

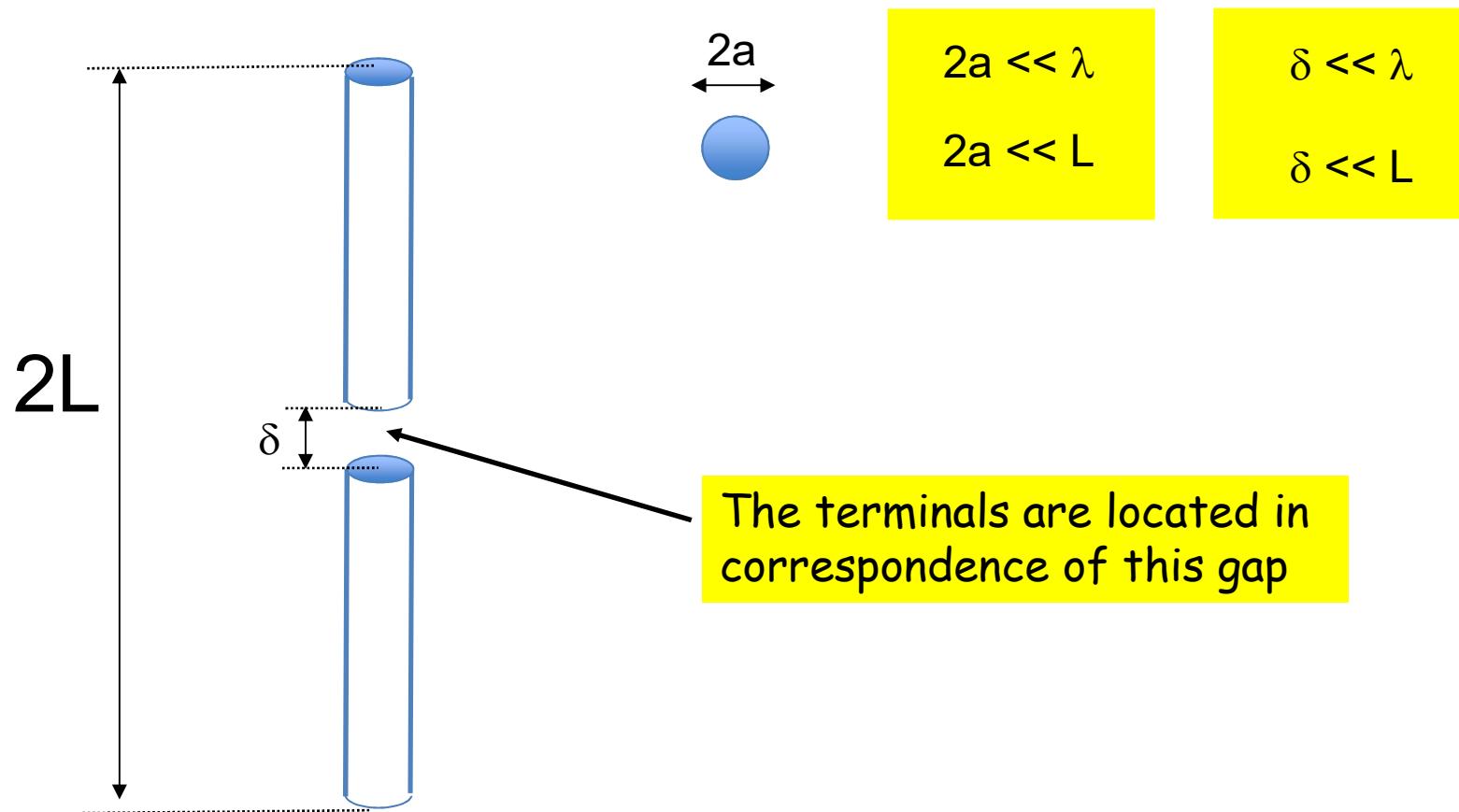
Ing. Stefano Perna

Wire antennas

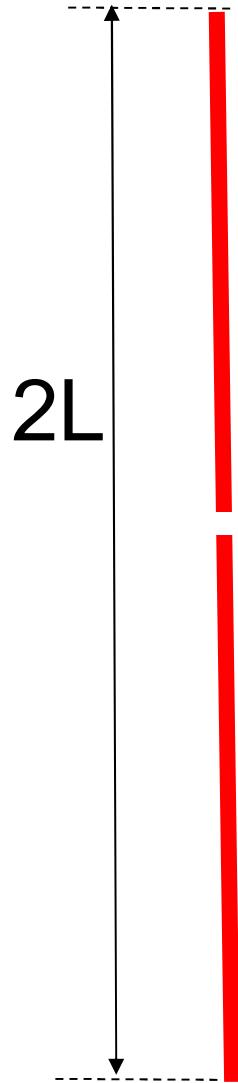
Wire antennas



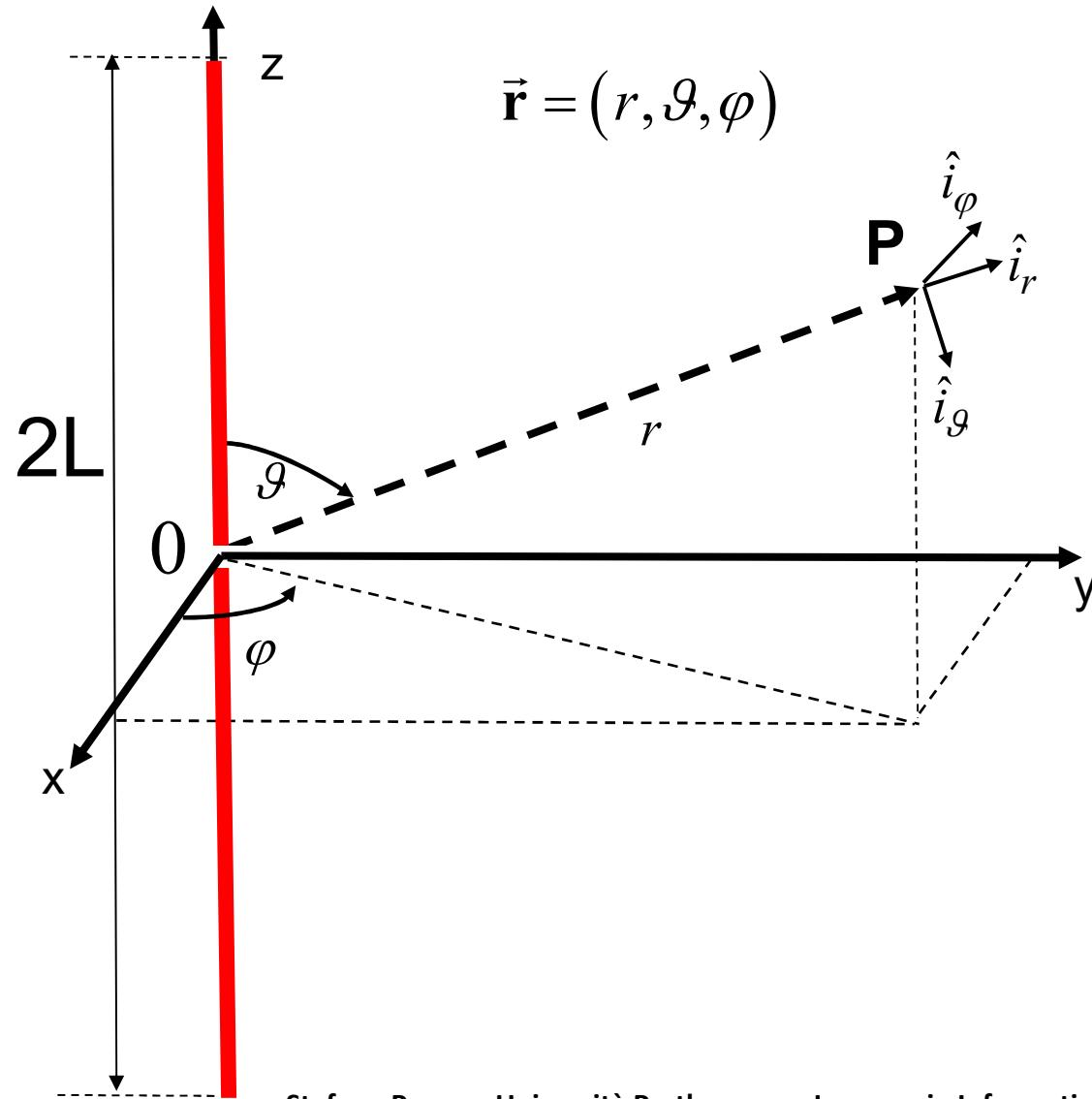
Wire antennas



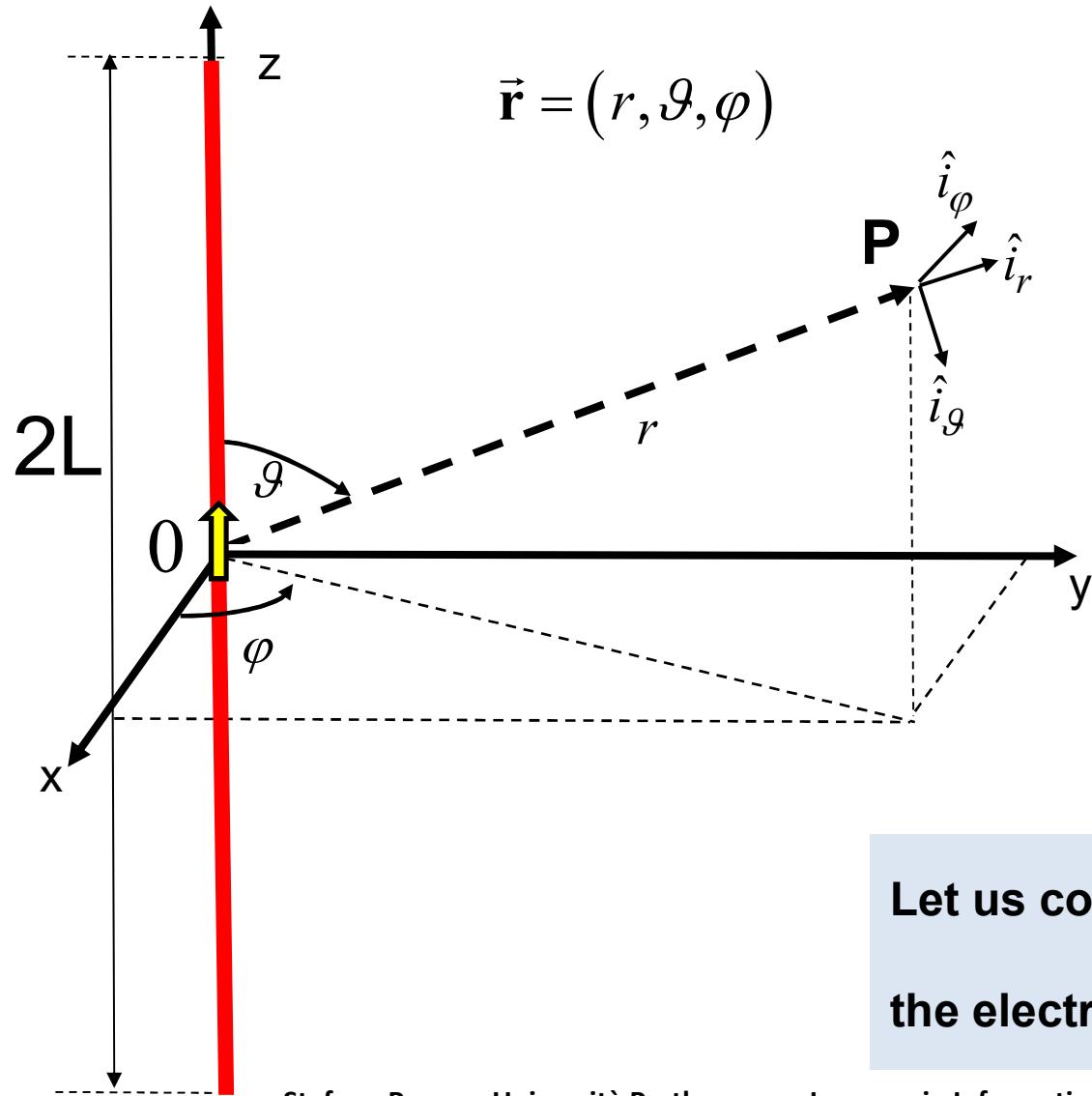
Wire antennas



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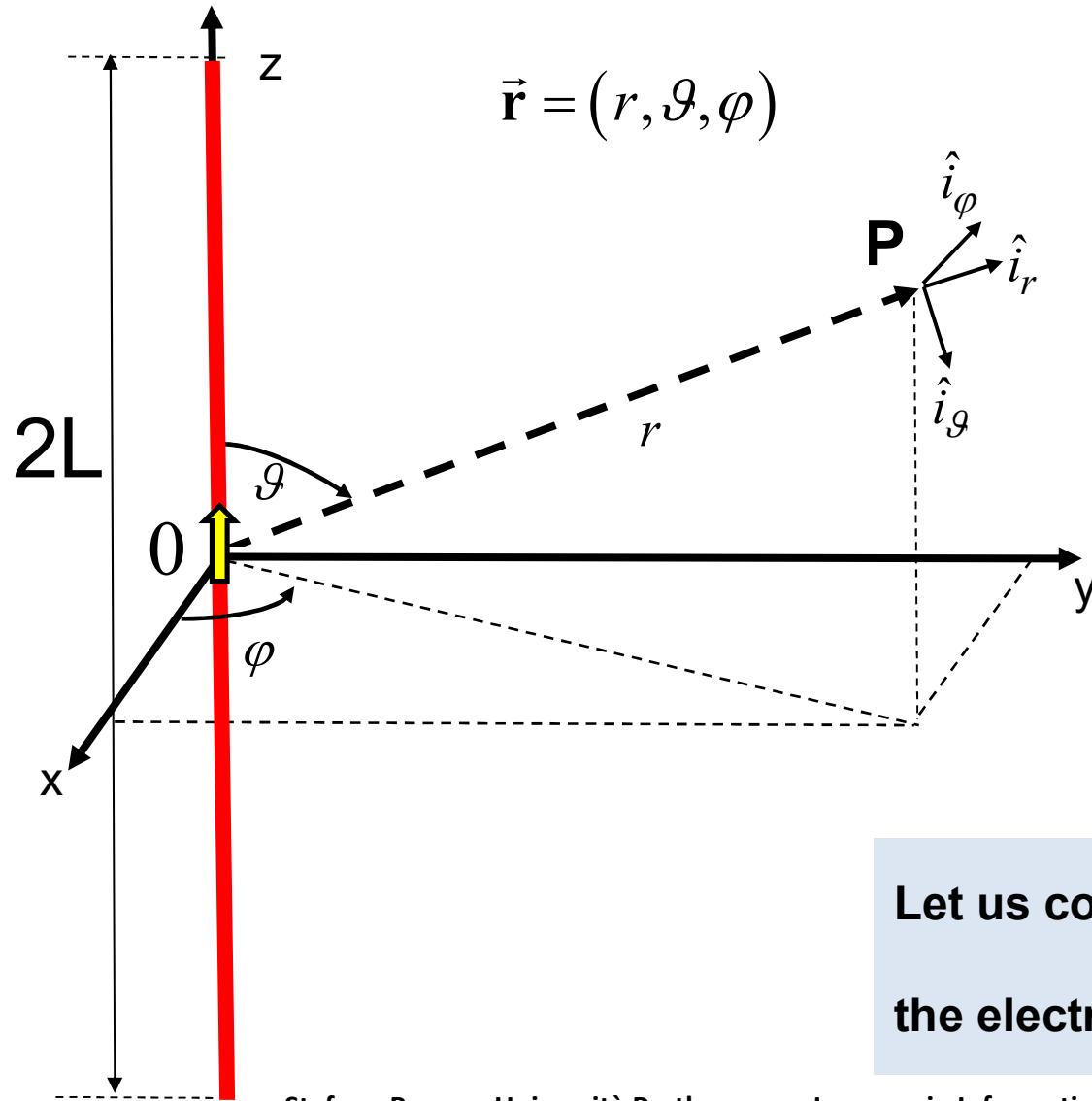


Wire antennas



Let us consider the elementary source:
the electric dipole  located in the origin

Wire antennas



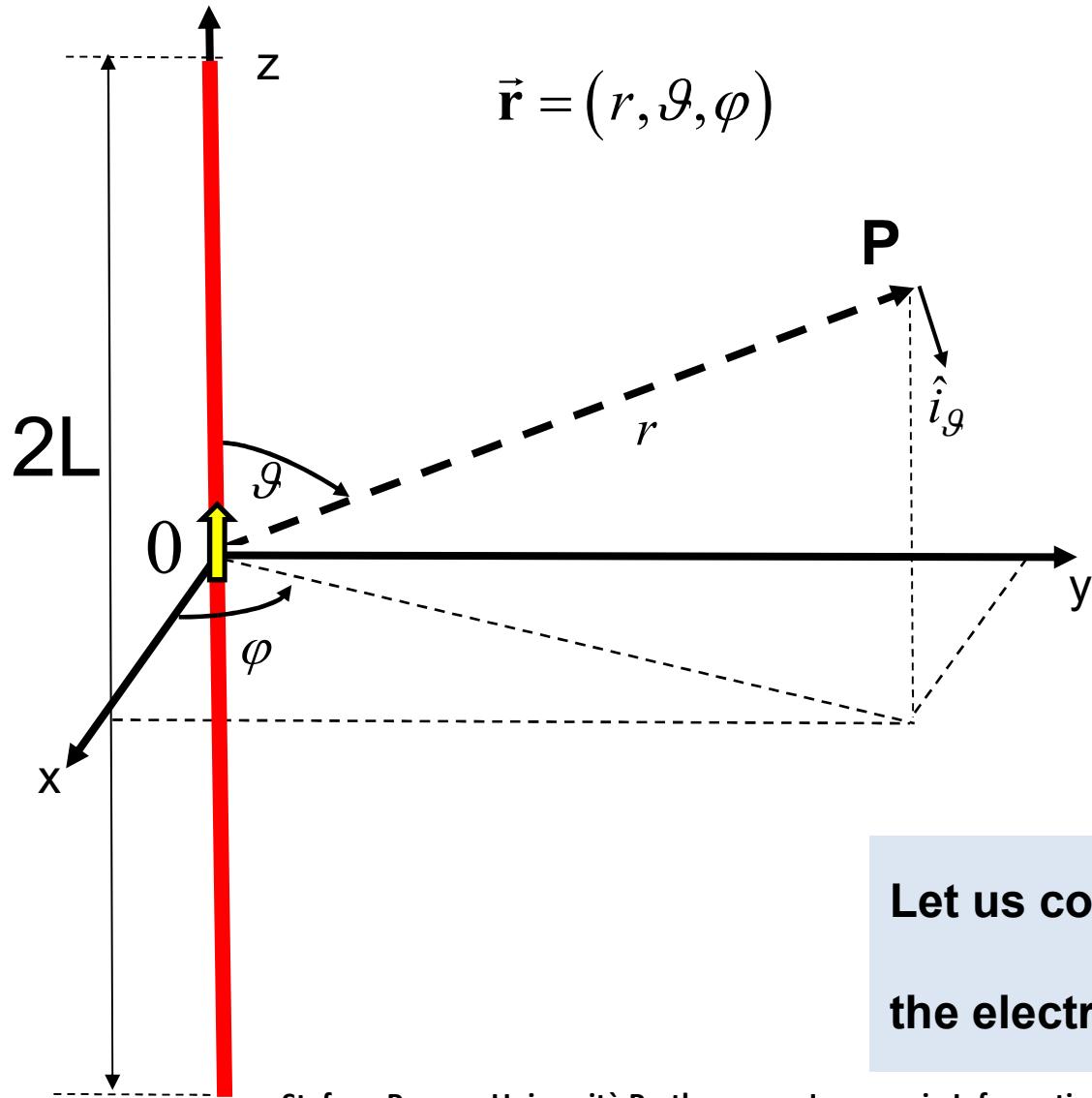
$$\vec{r} = (r, \vartheta, \varphi)$$

$$\vec{E}_0 = j \frac{\zeta I_0}{2\lambda} \frac{\exp(-j\beta R_0)}{R_0} \Delta z \sin \vartheta_0 \hat{i}_{\vartheta_0}$$

$$\vartheta_0 = \vartheta; \quad R_0 = r; \quad \hat{i}_{\vartheta_0} = \hat{i}_\vartheta$$

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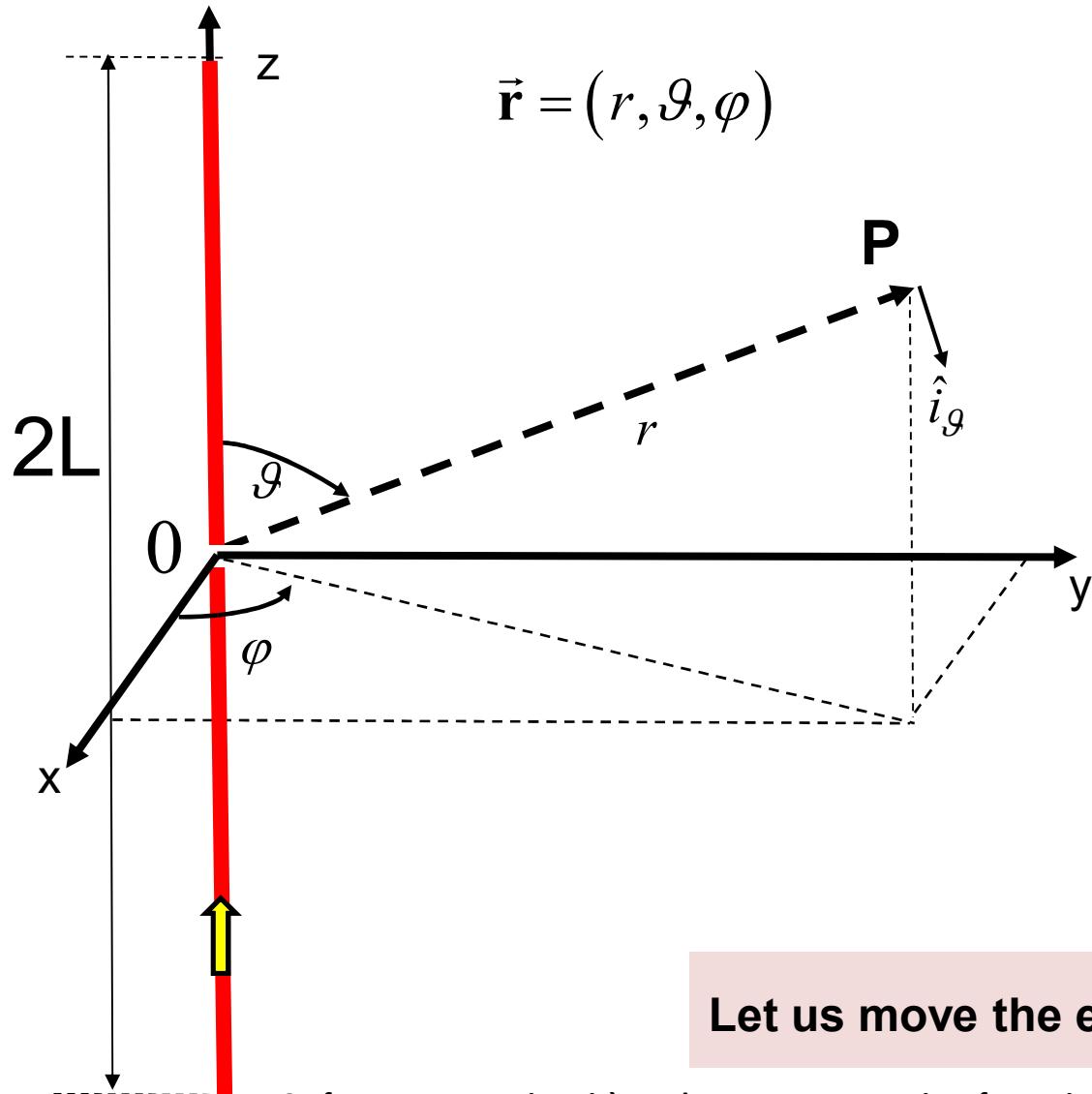
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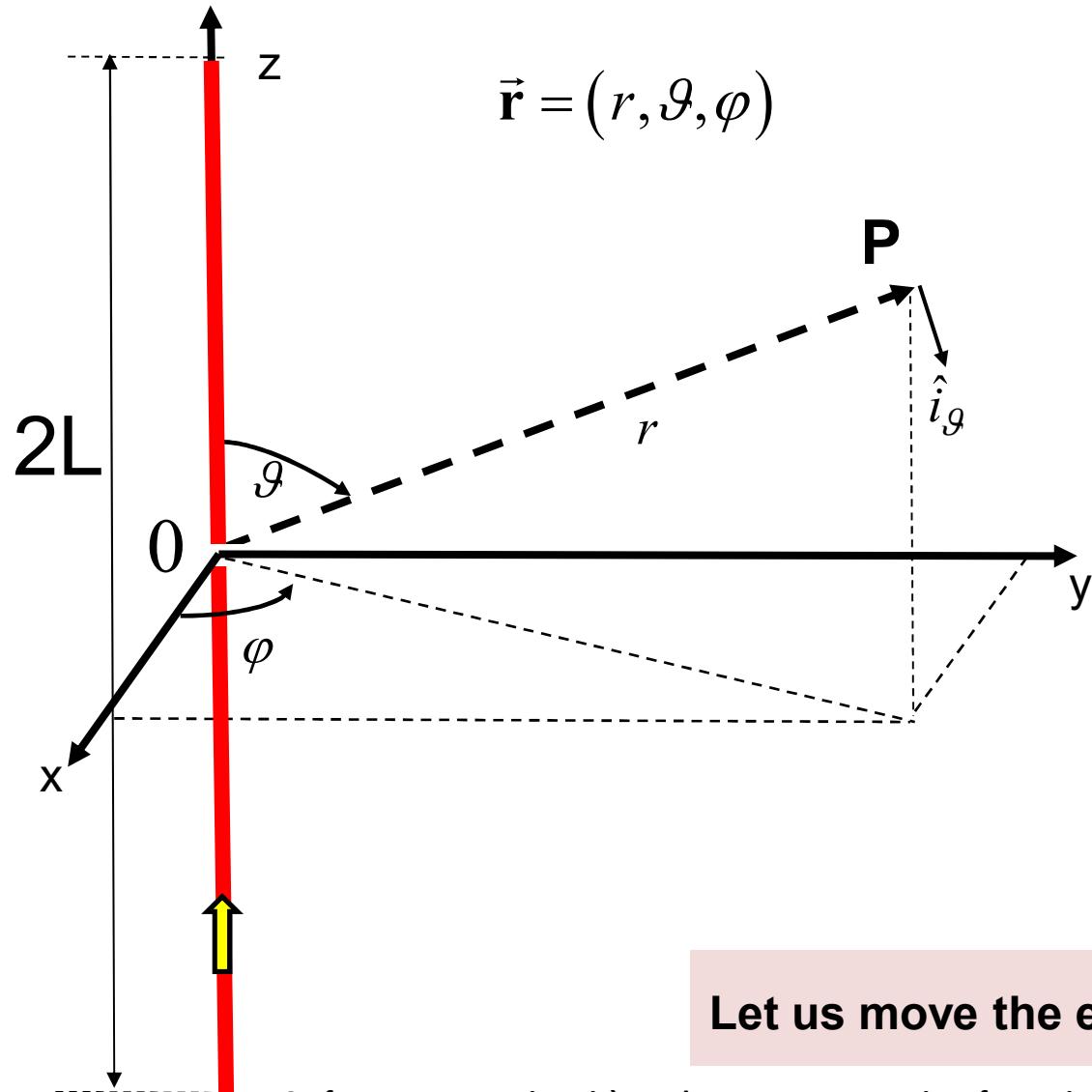
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Let us move the electric dipole along the z -axis

Wire antennas



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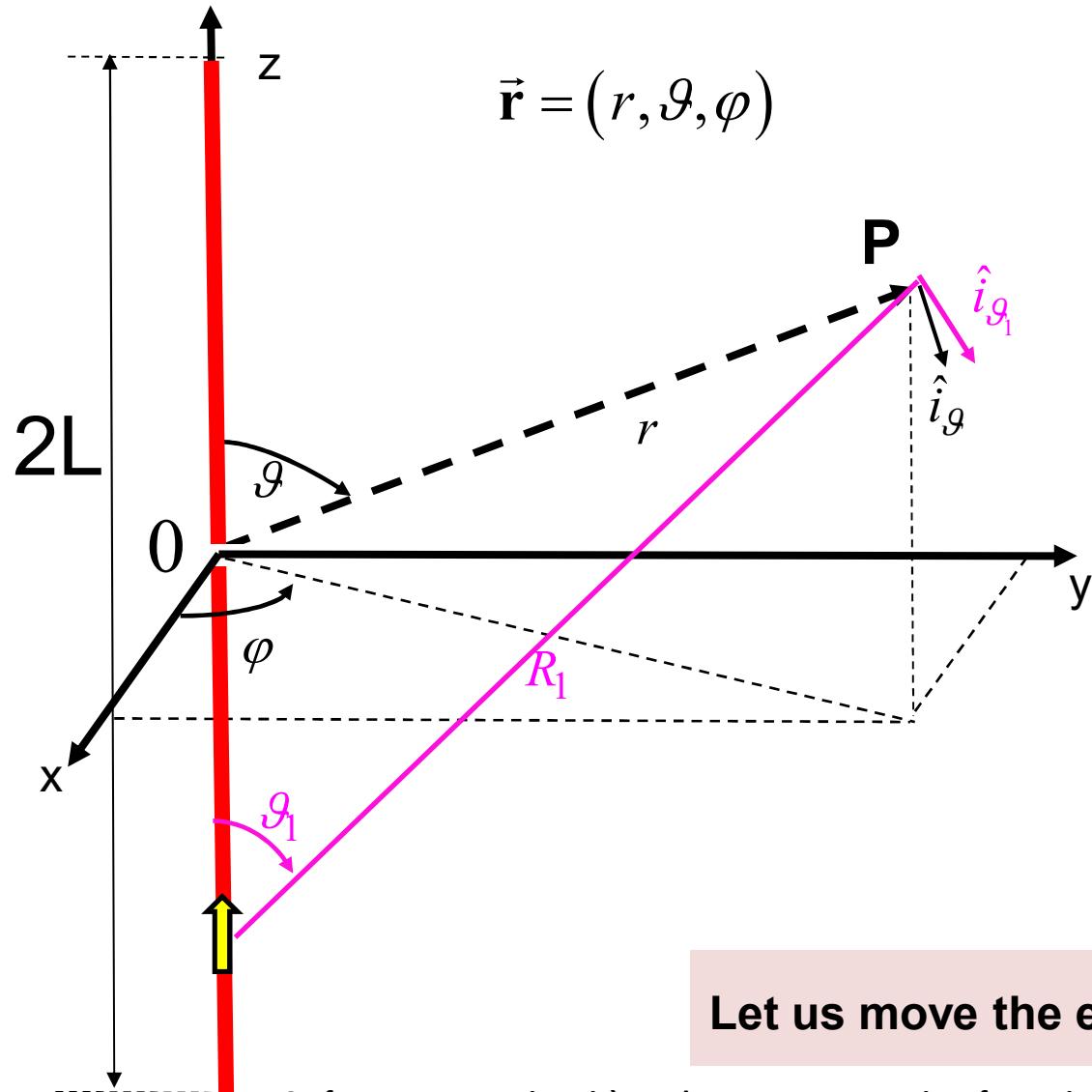
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Wire antennas



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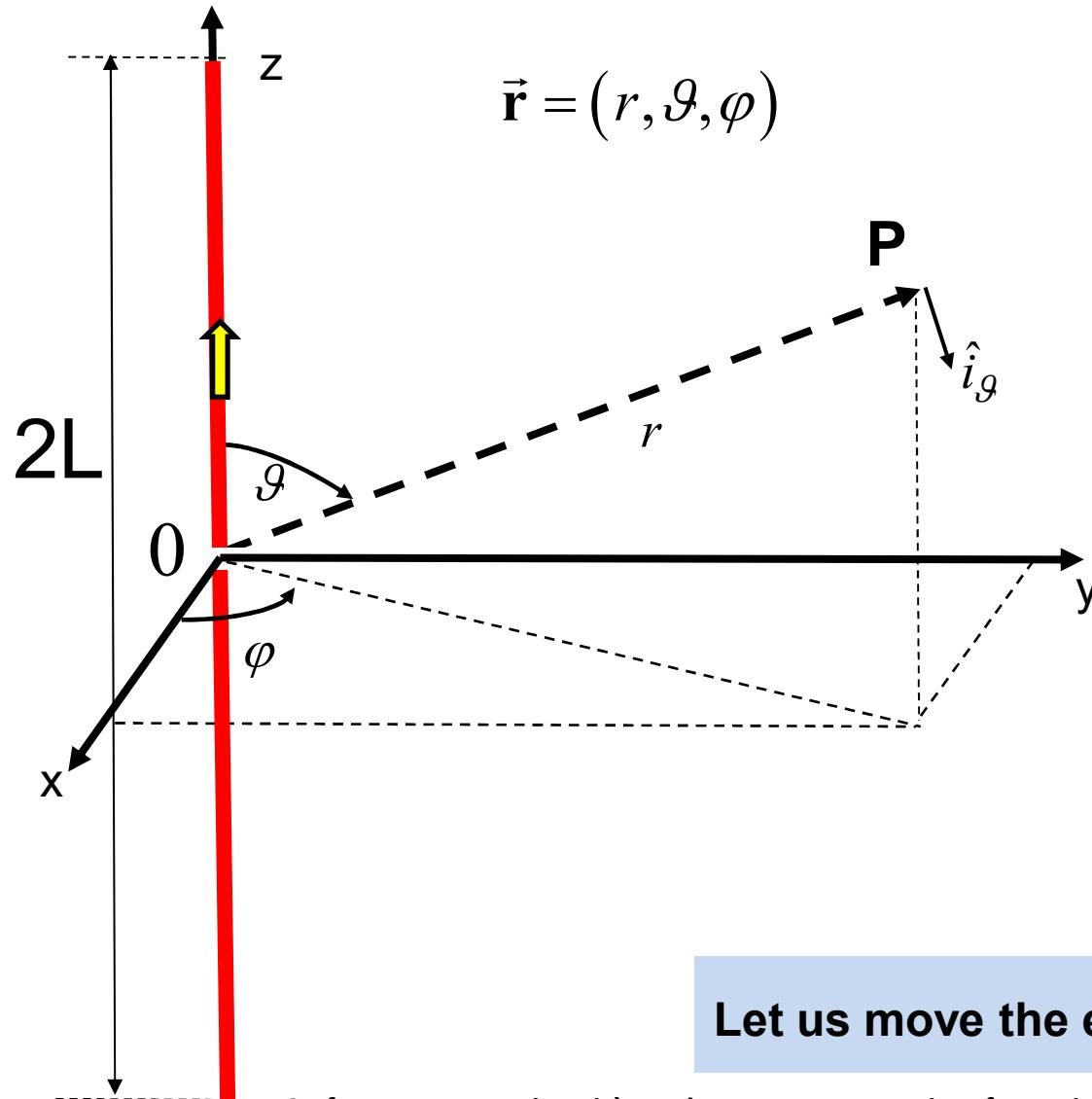
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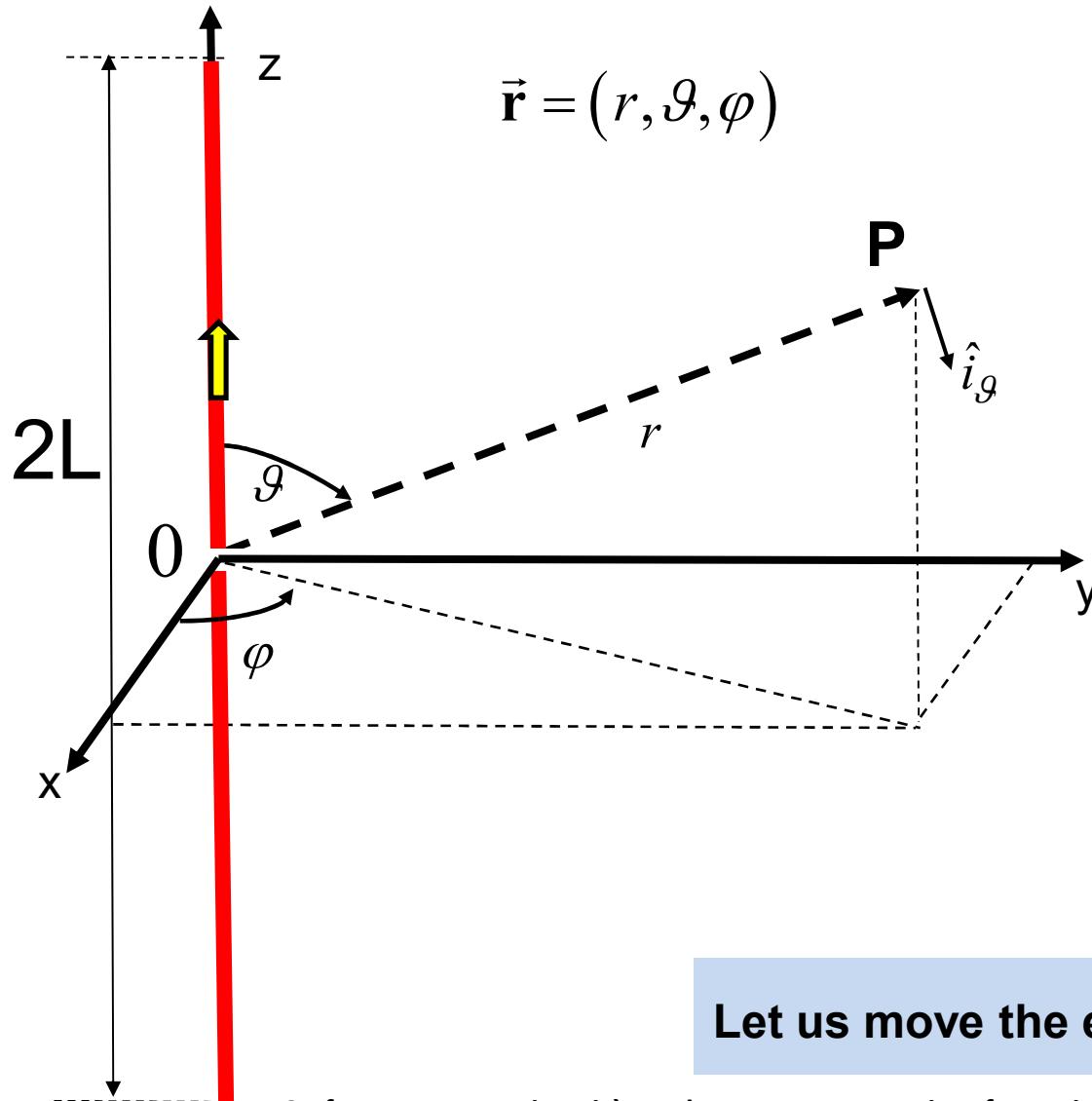
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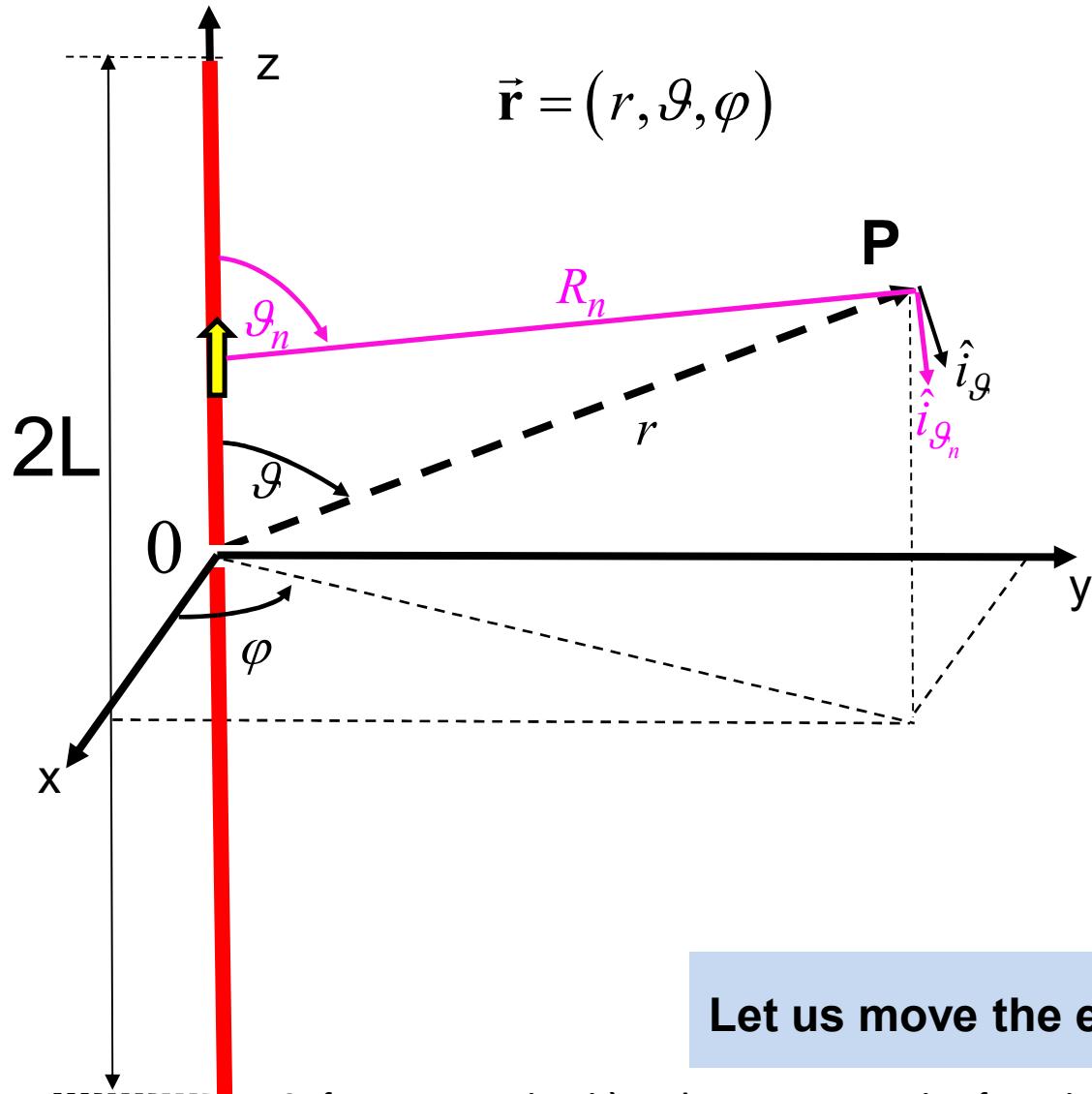
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Wire antennas



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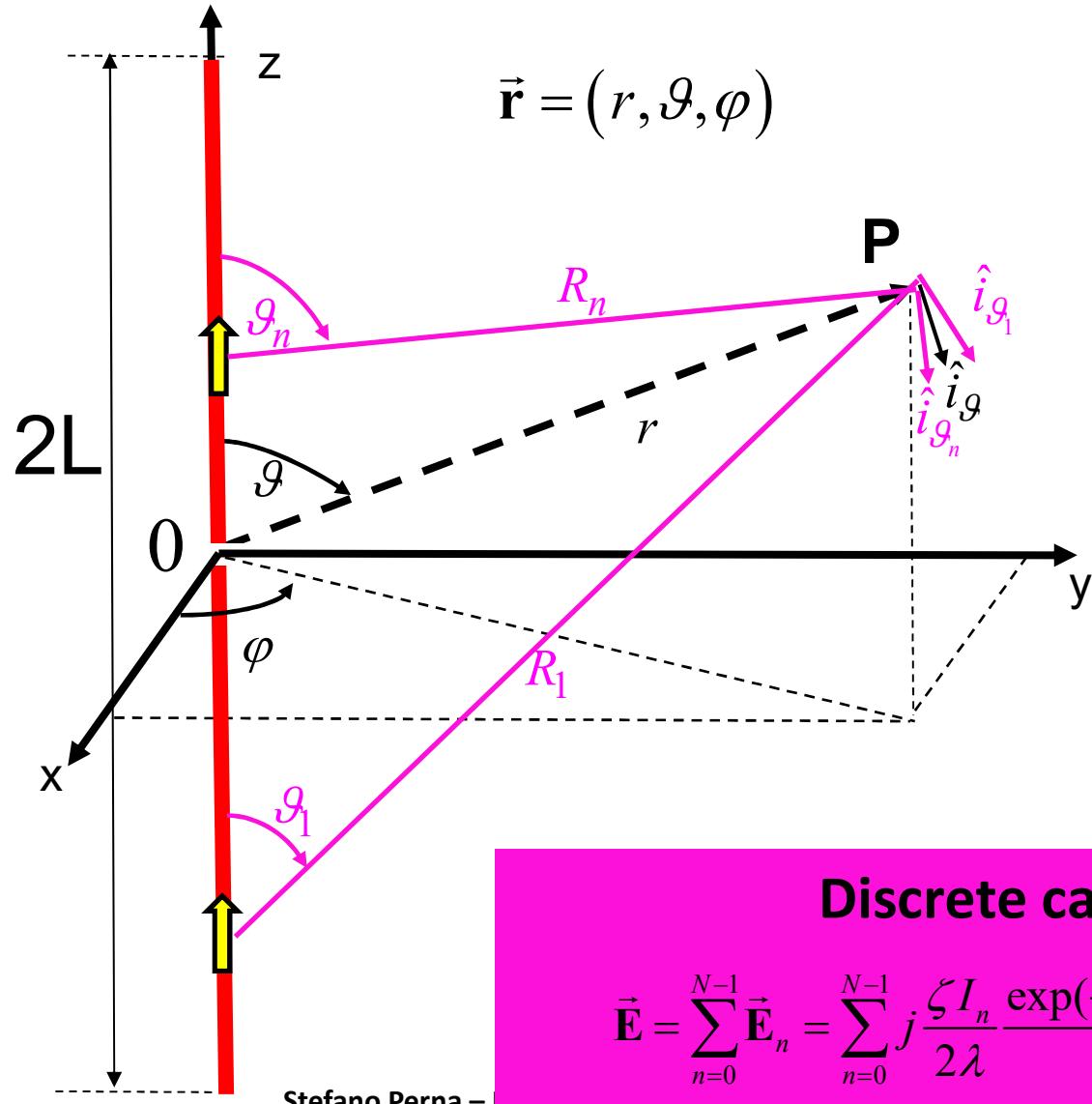
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Wire antennas



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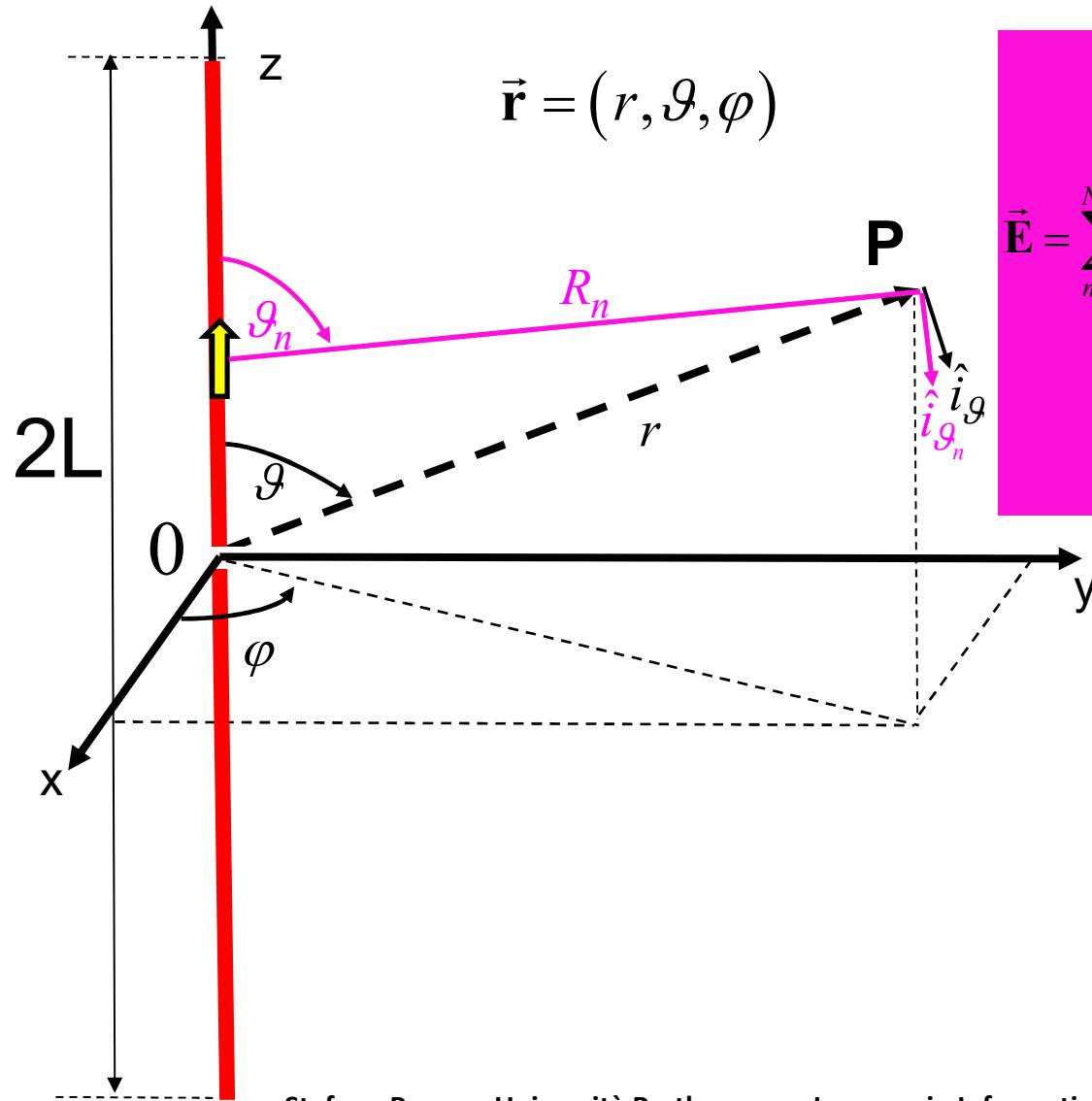
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Discrete case

$$\vec{E} = \sum_{n=0}^{N-1} \vec{E}_n = \sum_{n=0}^{N-1} j \frac{\zeta I_n}{2\lambda} \frac{\exp(-j\beta R_n)}{R_n} \Delta z \sin \vartheta_n \hat{i}_{\vartheta_n}$$

di "Antenne"

Wire antennas

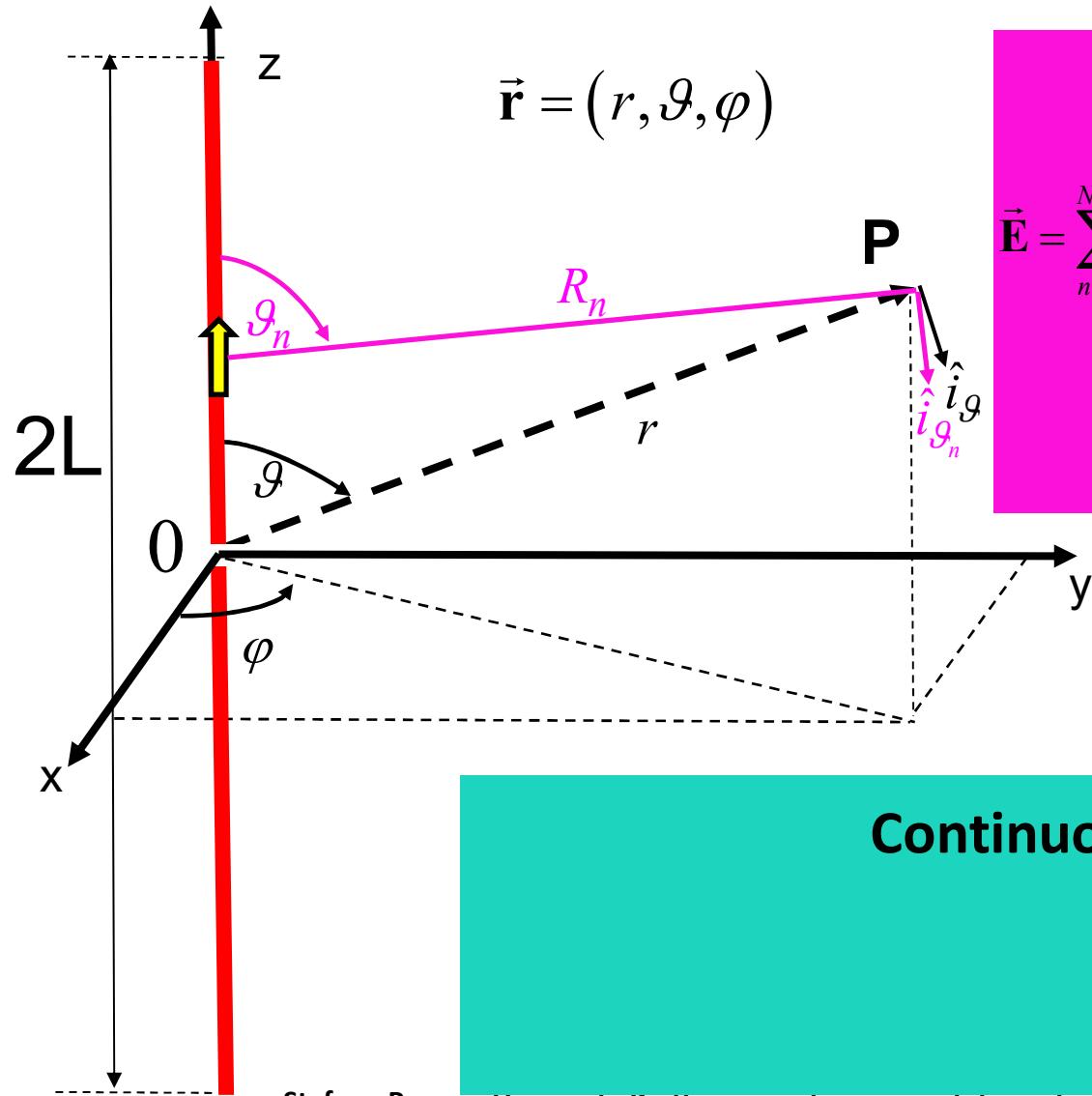


Discrete case

$$\vec{\mathbf{E}} = \sum_{n=0}^{N-1} \vec{\mathbf{E}}_n = \sum_{n=0}^{N-1} j \frac{\zeta I_n}{2\lambda} \frac{\exp(-j\beta R_n)}{R_n} \Delta z \sin \vartheta_n \hat{i}_{\vartheta_n}$$

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Wire antennas



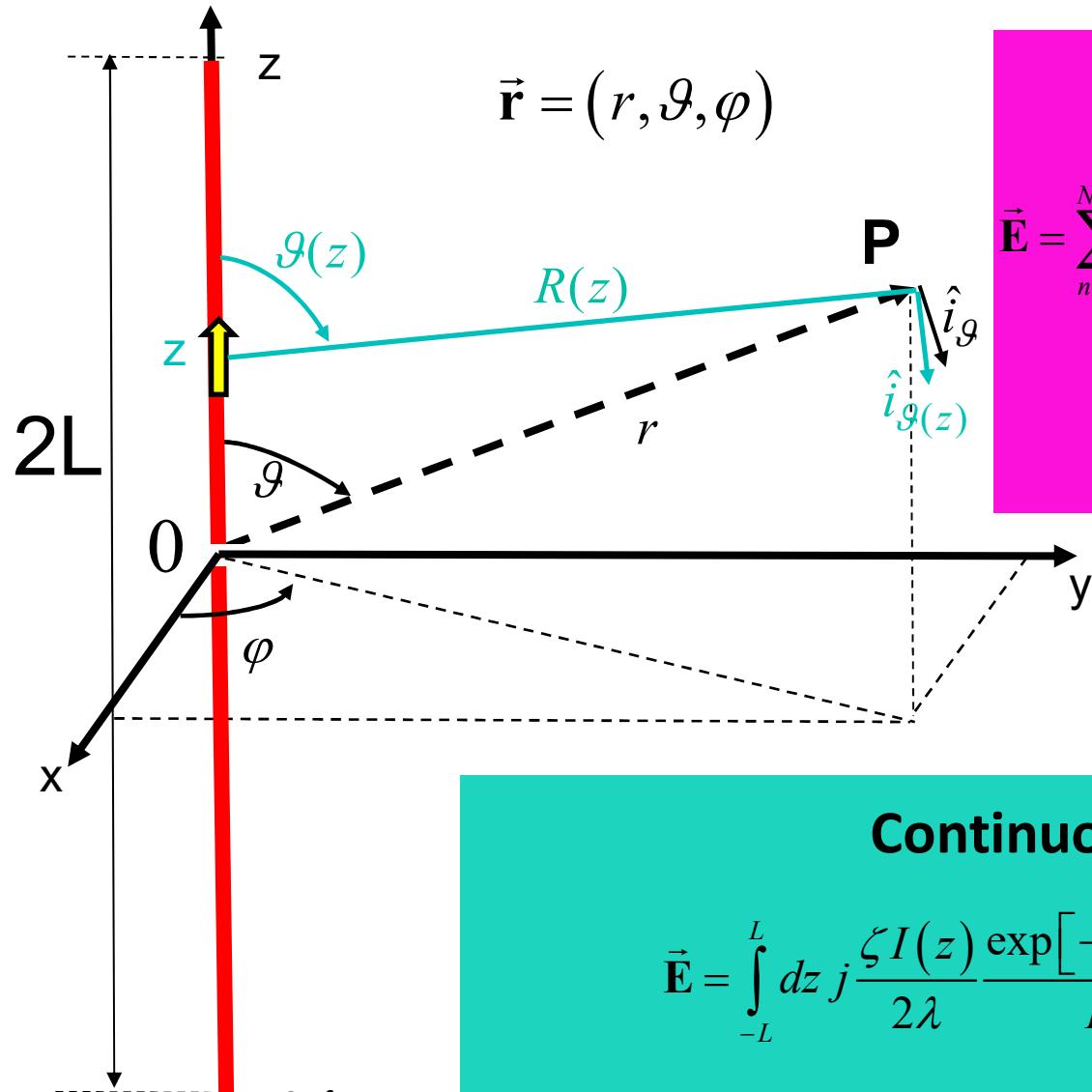
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Continuous case

Wire antennas



Discrete case

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Continuous case

$$\vec{\mathbf{E}} = \int_{-L}^L dz j \frac{\zeta I(z)}{2\lambda} \frac{\exp[-j\beta R(z)]}{R(z)} \sin \vartheta(z) \hat{i}_{\vartheta(z)}$$

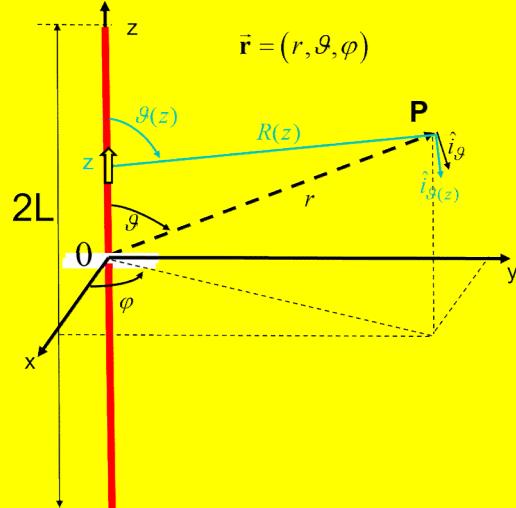
Wire antennas

$$\vec{E} = \int_{-L}^L dz j \frac{\zeta I(z) \exp[-j\beta R(z)]}{2\lambda} \sin \vartheta(z) \hat{i}_{\vartheta(z)} = j \frac{\zeta}{2\lambda} \int_{-L}^L dz I(z) \frac{\exp[-j\beta R(z)]}{R(z)} \sin \vartheta(z) \hat{i}_{\vartheta(z)}$$

Wire antennas

$$\begin{aligned}\vec{\mathbf{E}} &= \int_{-L}^L dz j \frac{\zeta I(z) \exp[-j\beta R(z)]}{2\lambda} \sin \vartheta(z) \hat{i}_{\vartheta(z)} = j \frac{\zeta}{2\lambda} \int_{-L}^L dz I(z) \frac{\exp[-j\beta R(z)]}{R(z)} \sin \vartheta(z) \hat{i}_{\vartheta(z)} \\ &= j \frac{\zeta}{2\lambda} \int_{-L}^L dz I(z) \frac{\exp[-j\beta r] \exp[j\beta z \cos \vartheta]}{r} \sin \vartheta \hat{i}_\vartheta\end{aligned}$$

Let us suppose that P is located in the **Fraunhofer Region** relevant to the considered wire antenna



■ $R(z) \approx r - \vec{\mathbf{r}}' \cdot \hat{i}_r = r - z \hat{i}_z \cdot \hat{i}_r = r - z \cos \vartheta$

$$\Rightarrow \frac{\exp[-j\beta R(z)]}{R(z)} \approx \frac{\exp[-j\beta r] \exp[j\beta z \cos \vartheta]}{r}$$

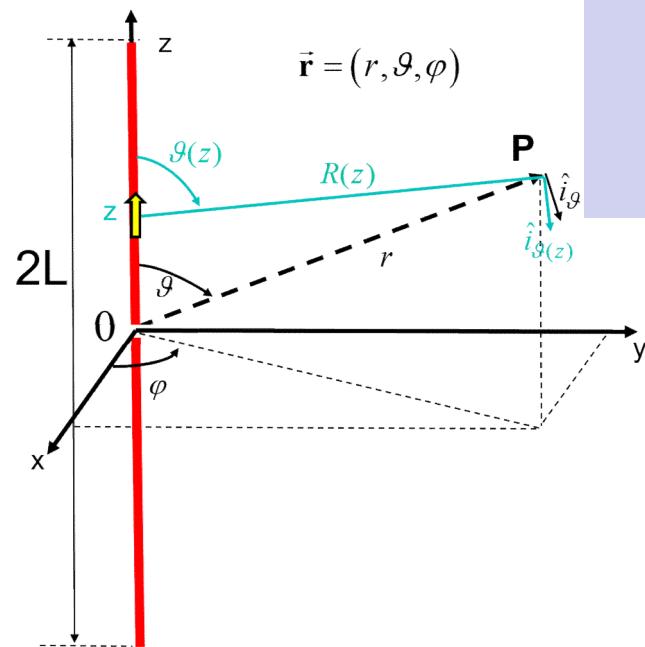
■ $\vartheta(z) \approx \vartheta$

■ $\hat{i}_{\vartheta(z)} \approx \hat{i}_\vartheta$

Wire antennas

$$\begin{aligned}\vec{\mathbf{E}} &= \int_{-L}^L dz j \frac{\zeta I(z) \exp[-j\beta R(z)]}{2\lambda R(z)} \sin \vartheta(z) \hat{i}_{\vartheta(z)} = j \frac{\zeta}{2\lambda} \int_{-L}^L dz I(z) \frac{\exp[-j\beta R(z)]}{R(z)} \sin \vartheta(z) \hat{i}_{\vartheta(z)} \\ &= j \frac{\zeta}{2\lambda} \int_{-L}^L dz I(z) \frac{\exp[-j\beta r] \exp(j\beta z \cos \vartheta)}{r} \sin \vartheta \hat{i}_\vartheta \\ &= j \frac{\zeta}{2\lambda} \frac{\exp[-j\beta r]}{r} \sin \vartheta \hat{i}_\vartheta \left[\int_{-L}^L dz I(z) \exp(j\beta z \cos \vartheta) \right] \\ &= j \frac{\zeta}{2\lambda} I_0 \frac{\exp[-j\beta r]}{r} \sin \vartheta \hat{i}_\vartheta \left[\int_{-L}^L dz \frac{I(z)}{I_0} \exp(j\beta z \cos \vartheta) \right]\end{aligned}$$

Wire antennas



In the Fraunhofer Region the expression of the radiated field simplifies as

$$\vec{E} = j \frac{\zeta}{2\lambda} I_0 \frac{\exp[-j\beta r]}{r} \sin \vartheta \hat{i}_\vartheta \int_{-l}^l dz \frac{I(z)}{I_0} \exp(j\beta z \cos \vartheta)$$

Color legend

New formulas, important considerations,
important formulas, important concepts

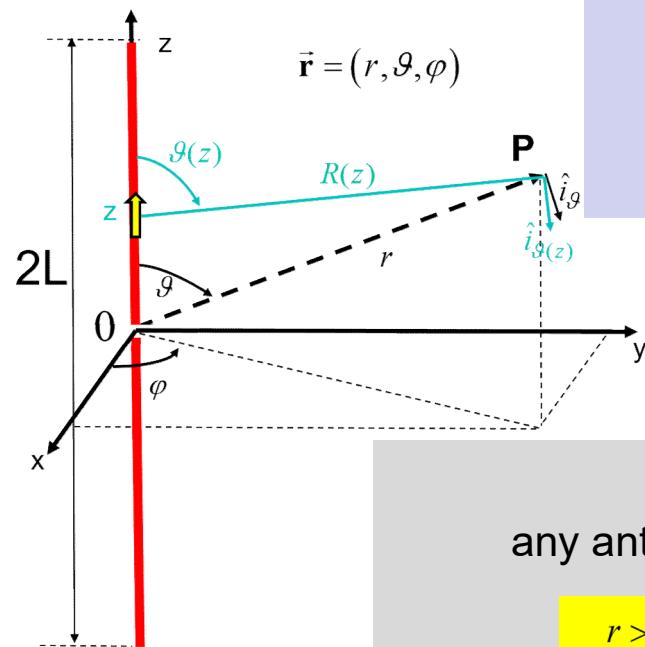
Very important for the discussion

Memo

Mathematical tools to be exploited

Mathematics

Wire antennas



In the Fraunhofer Region the expression of the radiated field simplifies as

$$\vec{E} = j \frac{\zeta}{2\lambda} I_0 \frac{\exp[-j\beta r]}{r} \sin \theta \left[\int_{-L}^L dz \frac{I(z)}{I_0} \exp(j\beta z \cos \theta) \right] \hat{i}_g$$

Effective length of the wire antenna

.... Memo

any antenna, in the Fraunhofer region, behaves as follows

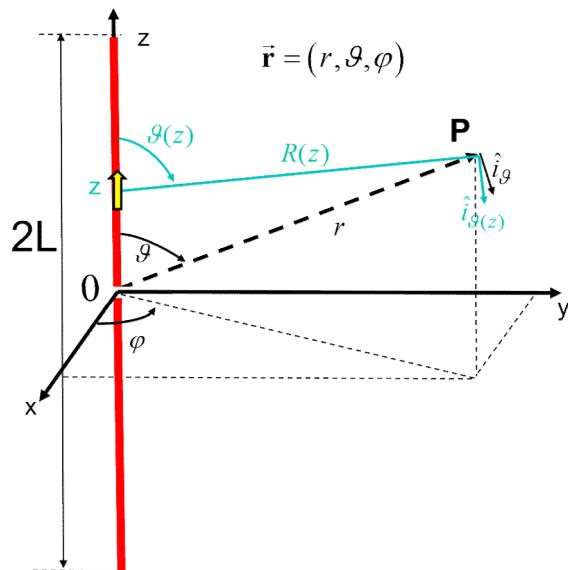
$$\begin{aligned} r &> D \\ r &> \frac{2D^2}{\lambda} \\ r &> \lambda \end{aligned}$$

$$\left\{ \begin{array}{l} \mathbf{E}(\vec{r}) = \mathbf{E}(r, \theta, \varphi) = \frac{j\zeta I}{2\lambda} \frac{e^{-j\beta r}}{r} \mathbf{l}(\theta, \varphi) \\ \zeta \mathbf{H} = \hat{i}_r \times \mathbf{E} \end{array} \right.$$

$$\mathbf{l}(\theta, \varphi) = l_g(\theta, \varphi) \hat{i}_g + l_\varphi(\theta, \varphi) \hat{i}_\varphi \quad \text{Effective length}$$

Wire antennas: effective length

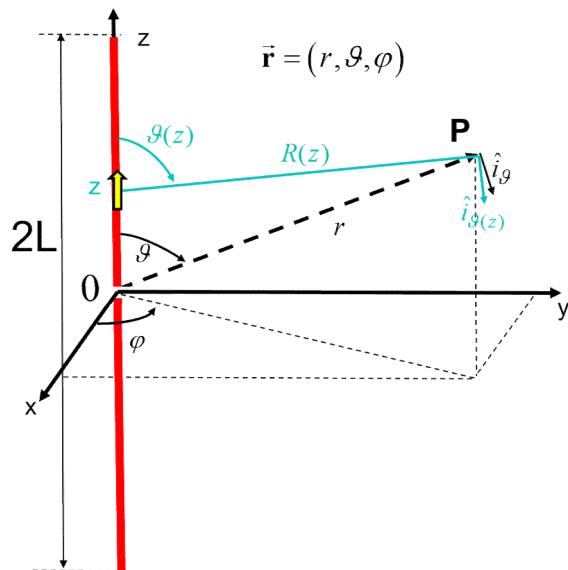
$$\vec{I} = \sin \vartheta \left[\int_{-L}^L dz \frac{I(z)}{I_0} \exp(j\beta z \cos \vartheta) \right] \hat{i}_\vartheta$$



 **The effective length is independent of ϕ**
... absolutely not surprising

Wire antennas: effective length

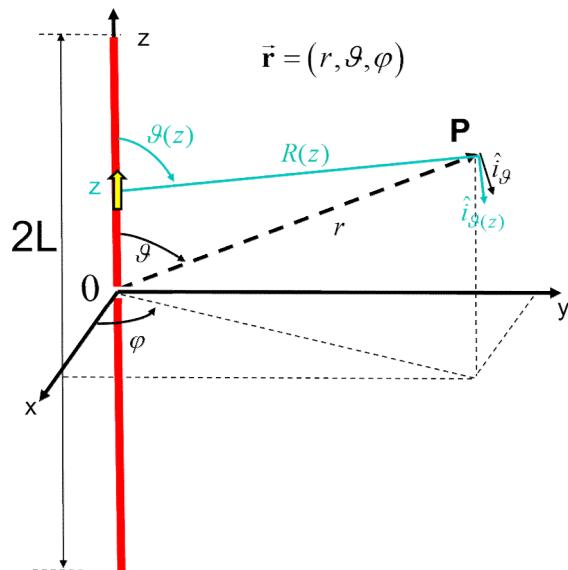
$$\vec{I}(\vartheta) = l_\vartheta(\vartheta) \hat{i}_\vartheta = \sin \vartheta \left[\int_{-L}^L dz \frac{I(z)}{I_0} \exp(j\beta z \cos \vartheta) \right] \hat{i}_\vartheta$$



- The effective length is independent of ϕ
... absolutely not surprising
- The effective length depends by the current distribution $I(z)$
... absolutely not surprising
- The effective length depends on L
... absolutely not surprising

Wire antennas: effective length

$$\vec{I}(\vartheta) = l_\vartheta(\vartheta) \hat{i}_\vartheta = \sin \vartheta \left[\int_{-L}^L dz \frac{I(z)}{I_0} \exp(j\beta z \cos \vartheta) \right] \hat{i}_\vartheta$$



$$u = -\beta \cos \vartheta \quad \tilde{I}(z) = \frac{I(z)}{I_0}$$
$$F(u) = \int_{-L}^L dz \tilde{I}(z) e^{-juz}$$

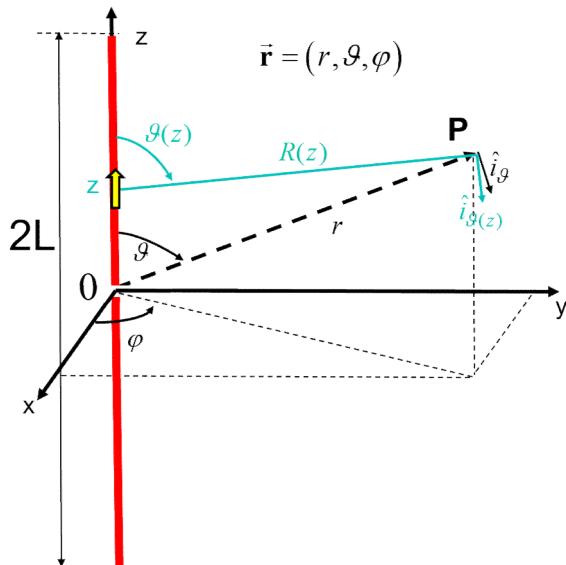
For the wire antennas the source impressed on the antenna is related to the radiated field through the Fourier Transformation rule

Wire antennas: effective length

$$\vec{E} = j \frac{\zeta}{2\lambda} I_0 \frac{\exp[-j\beta r]}{r} \left[\sin \vartheta F(\vartheta) \hat{i}_\vartheta \right]$$

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$$F(\vartheta) = F(u) \Big|_{u = -\beta \cos \vartheta}$$
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The properties of the Fourier Transformation suggest some interesting considerations

Wire antennas: effective length

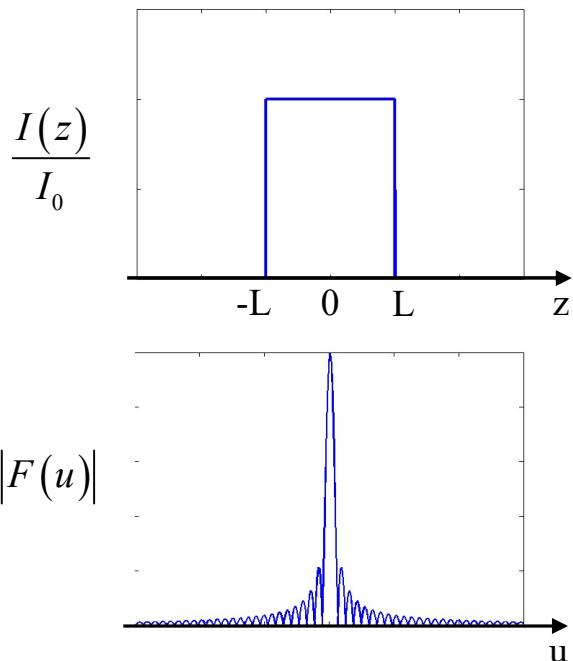
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The properties of the Fourier Transformation suggest some interesting considerations

- Antenna's size and beamwidth

Wire antennas: effective length

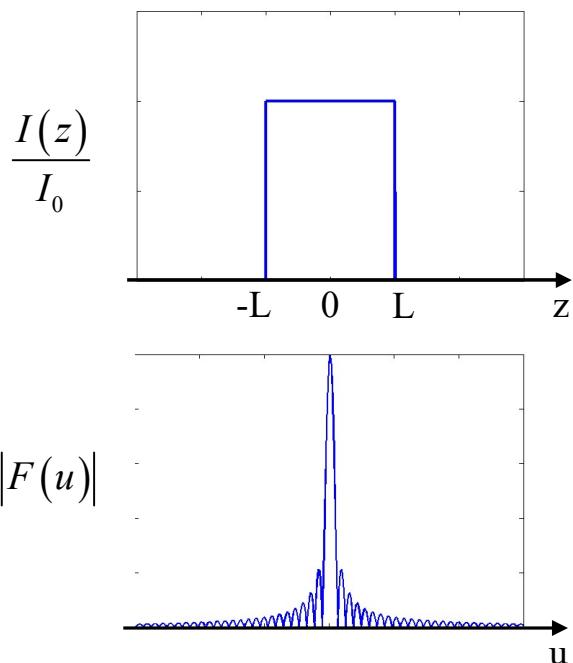
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The properties of the Fourier Transformation suggest some interesting considerations

- Antenna's size and beamwidth
- Scanning of the pattern
- Synthesis of the pattern

Wire antennas: visible region

$$\vec{E} = j \frac{\zeta}{2\lambda} I_0 \frac{\exp[-j\beta r]}{r} \left[\sin \vartheta F(\vartheta) \hat{i}_\vartheta \right]$$

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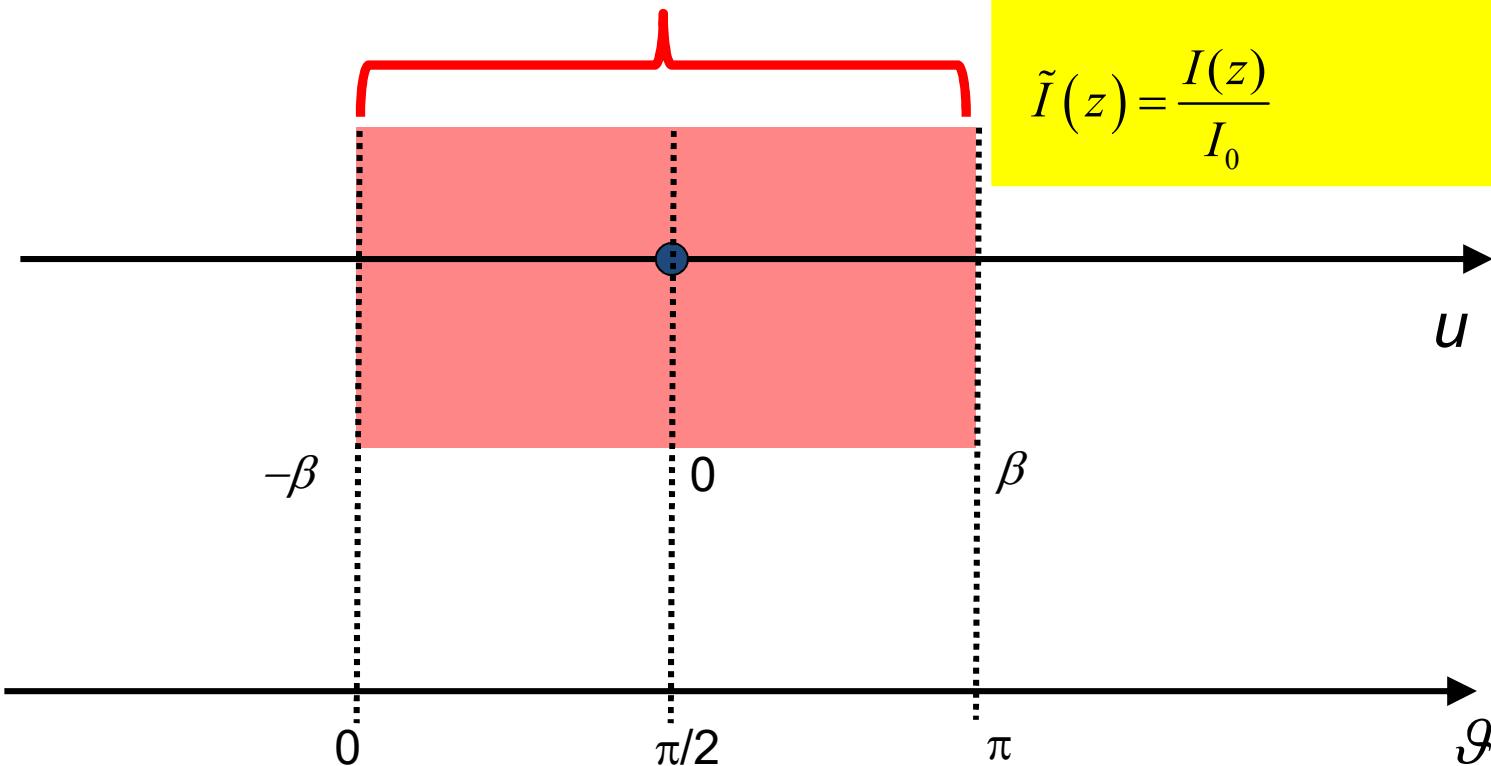
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Wire antennas: visible region

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Visible region of the spectrum



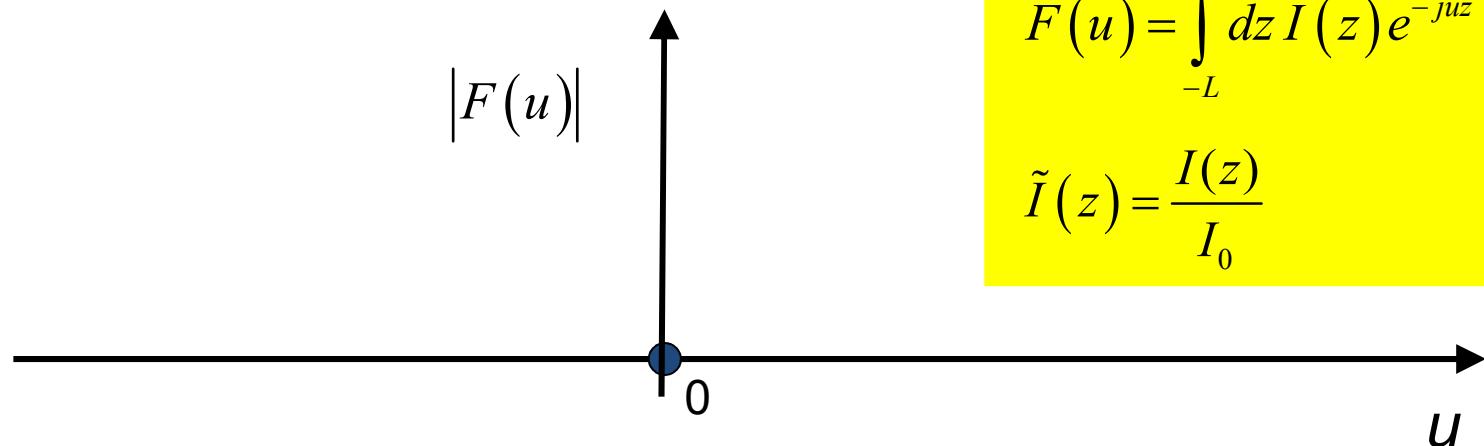
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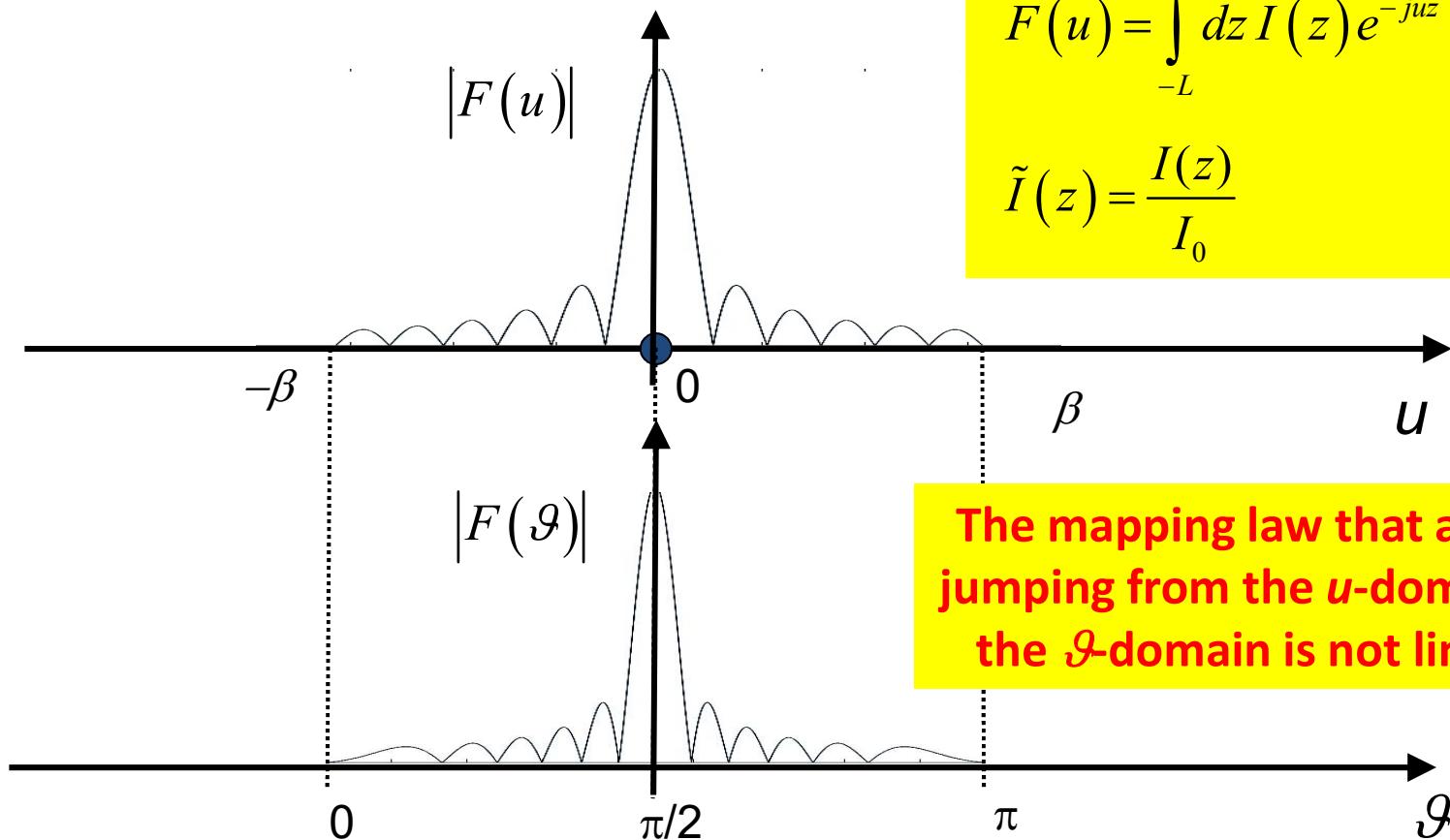
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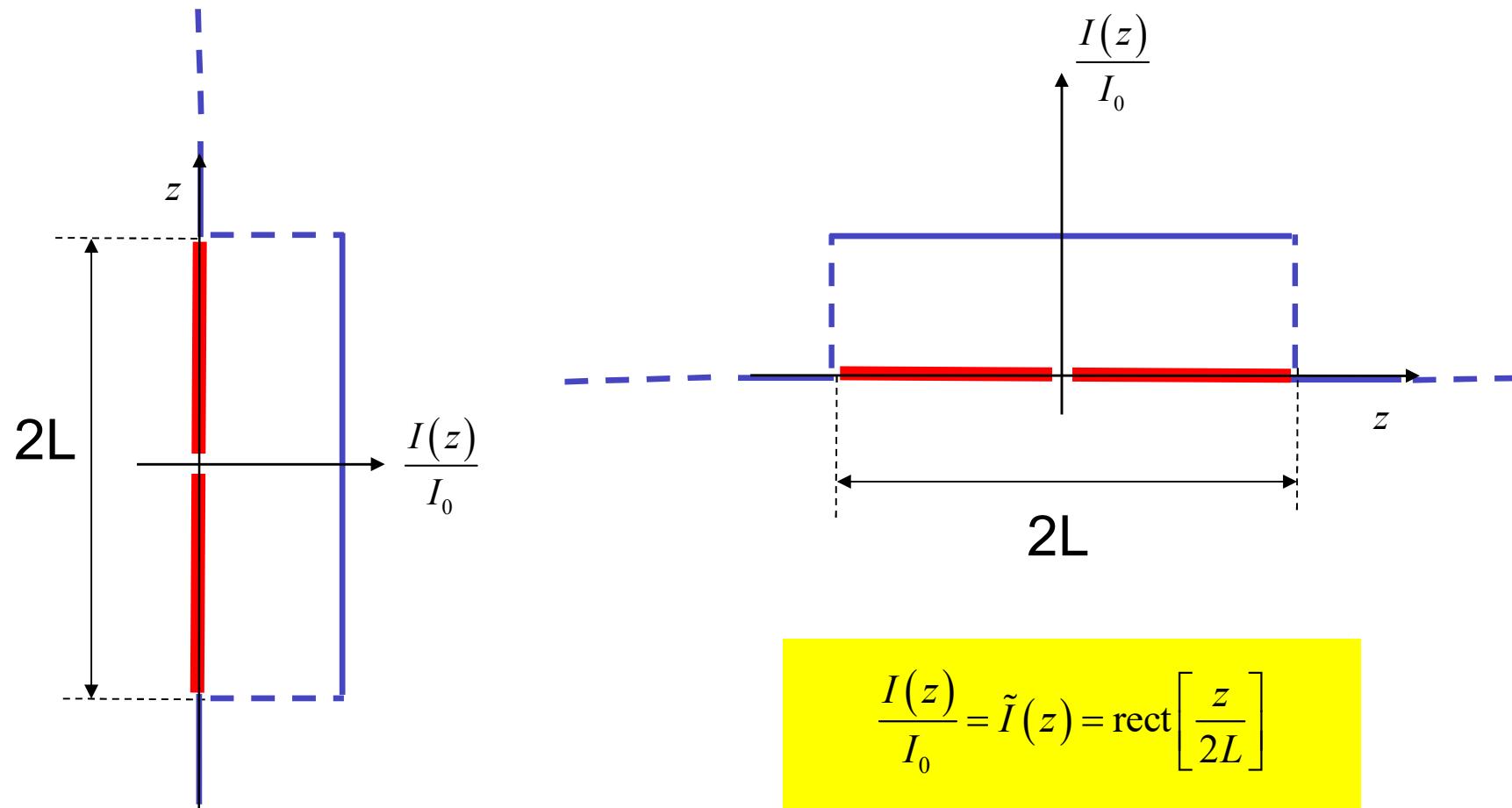
Memo

Mathematical tools to be exploited

Mathematics

Wire antennas: an ideal case

Uniform current distribution



$$\frac{I(z)}{I_0} = \tilde{I}(z) = \text{rect}\left[\frac{z}{2L}\right]$$