Campi Elettromagnetici

Corso di Laurea in Ingegneria Informatica, Biomedica e delle Telecomunicazioni

a.a. 2023-2024 - Laurea "Triennale" – Secondo semestre - Secondo anno

Università degli Studi di Napoli "Parthenope"

Stefano Perna

Color legend

New formulas, important considerations,
important formulas, important concepts

Very important for the discussion

Memo

Mathematical tools to be exploited

Mathematics

THEOREMS

Poynting

Time domain – Phasor domain

Uniqueness (Interior problem – Exterior problem)

Time domain – Phasor domain

Equivalence

Phasor domain

Image Theory

Reciprocity

Phasor domain

Maxwell Equations (Spectral Domains)

$$\begin{cases} \nabla \times \vec{\mathbf{E}} = -j\omega \vec{\mathbf{B}} \\ \nabla \times \vec{\mathbf{H}} = j\omega \vec{\mathbf{D}} + \vec{\mathbf{J}} \\ \nabla \cdot \vec{\mathbf{D}} = \rho \\ \nabla \cdot \vec{\mathbf{B}} = 0 \end{cases}$$

James Clerk Maxwell 1831-1879

Maxwell Equations (Spectral Domains)

Magnetic Sources

James Clerk Maxwell 1831-1879

$$\begin{cases} \nabla \times \vec{\mathbf{E}} = -j\omega\vec{\mathbf{B}} - \vec{\mathbf{J}}_{\mathrm{m}} \\ \nabla \times \vec{\mathbf{H}} = j\omega\vec{\mathbf{D}} + \vec{\mathbf{J}} \\ \nabla \cdot \vec{\mathbf{D}} = \rho \\ \nabla \cdot \vec{\mathbf{B}} = \rho_{\mathrm{m}} \end{cases}$$

$$\left[\vec{\mathbf{e}}(\vec{\mathbf{r}},t)\right] : \frac{Volt}{m} \qquad \left[\vec{\mathbf{b}}(\vec{\mathbf{r}},t)\right] : \frac{Weber}{m^2}$$

$$\left[\vec{\mathbf{j}}_m(\vec{\mathbf{r}},t)\right] : \frac{Volt}{m^2} \qquad \left[\rho_m(\vec{\mathbf{r}},t)\right] : \frac{Weber}{m^3}$$

Equivalence theorem

$\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0$

Consider a source distribution $\vec{\mathbf{J}}_0$ with its associated electromagnetic field $\left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$

Equivalence theorem

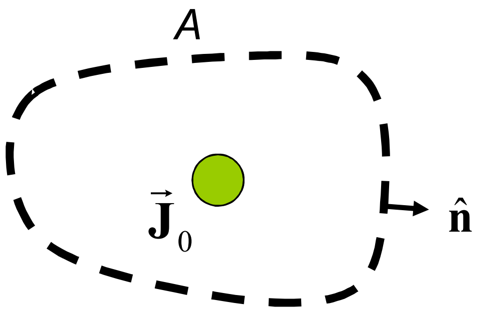
Consider a source distribution $\vec{\mathbf{J}}_0$ with its associated electromagnetic field $\left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$

Equivalence theorem

Consider a source distribution $\vec{\mathbf{J}}_0$ with its associated electromagnetic field $\left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$

$$\vec{\mathbf{J}}_0 \rightarrow \left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$$

Equivalence theorem



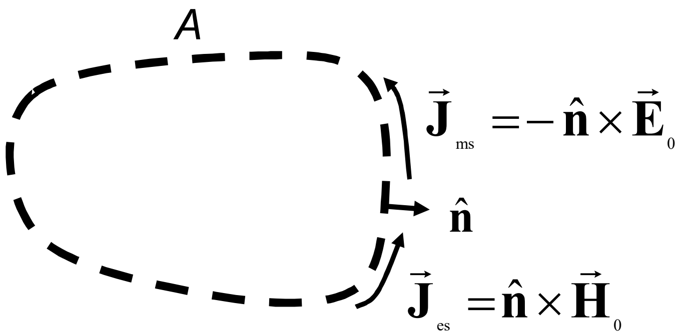
$\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0$

Consider a source distribution $\vec{\mathbf{J}}_0$ with its associated electromagnetic field $\left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$

Consider a (smooth) surface *A* with an everywhere defined unit normal $\hat{\mathbf{n}}$

$$\vec{\mathbf{J}}_0 \rightarrow \left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$$

Equivalence theorem



$\left[\vec{\mathbf{e}}(\vec{\mathbf{r}},t)\right] : \frac{Volt}{m}$ $\left[\vec{\mathbf{j}}_m(\vec{\mathbf{r}},t)\right] : \frac{Volt}{m^2}$

$\left[\vec{\mathbf{j}}_{ms}(\vec{\mathbf{r}},t)\right] : \frac{Volt}{m}$

Consider a source distribution $\vec{\mathbf{J}}_0$ with its associated electromagnetic field $\left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$

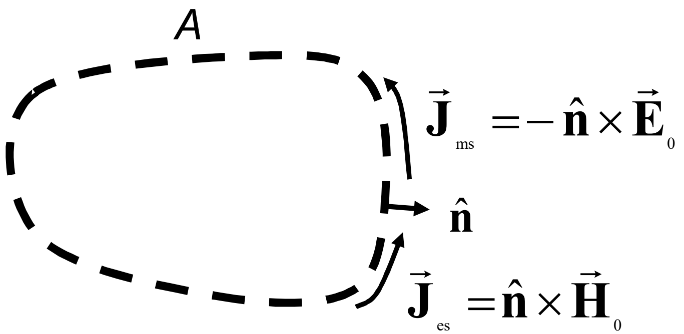
Consider a (smooth) surface *A* with an everywhere defined unit normal $\hat{\mathbf{n}}$

The original sources $\vec{\mathbf{J}}_0$ enclosed in *A* can be removed and substituted by equivalent sources, i.e., electric $\vec{\mathbf{J}}_{es} = \hat{\mathbf{n}} \times \vec{\mathbf{H}}_0$ and magnetic $\vec{\mathbf{J}}_{ms} = -\hat{\mathbf{n}} \times \vec{\mathbf{E}}_0$ current densities distributed over the surface *A*.

$$\vec{\mathbf{J}}_0 \rightarrow \left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$$

Equivalence theorem

$\left[\vec{h}(\vec{\mathbf{r}},t)\right]: \frac{Ampere}{m}$	$\left[\vec{\mathbf{j}}_e(\vec{\mathbf{r}},t)\right]: \frac{Ampere}{m^2}$
$\left[\vec{\mathbf{j}}_{ms}(\vec{\mathbf{r}},t)\right]: \frac{Volt}{m}$	$\left[\vec{\mathbf{j}}_{es}(\vec{\mathbf{r}},t)\right]: \frac{Ampere}{m}$



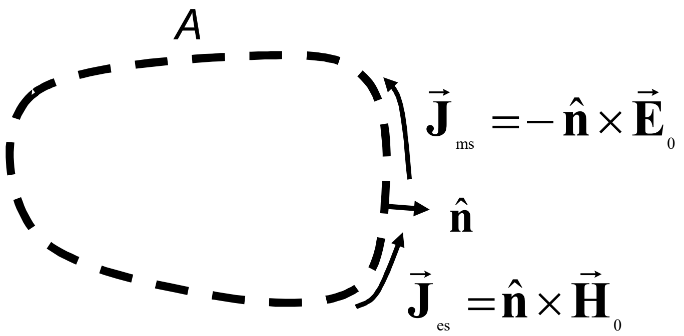
Consider a source distribution $\vec{\mathbf{J}}_0$ with its associated electromagnetic field $\left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$

Consider a (smooth) surface *A* with an everywhere defined unit normal $\hat{\mathbf{n}}$

The original sources $\vec{\mathbf{J}}_0$ enclosed in *A* can be removed and substituted by equivalent sources, i.e., electric $\vec{\mathbf{J}}_{es} = \hat{\mathbf{n}} \times \vec{\mathbf{H}}_0$ and magnetic $\vec{\mathbf{J}}_{ms} = -\hat{\mathbf{n}} \times \vec{\mathbf{E}}_0$ current densities distributed over the surface *A*.

$$\vec{\mathbf{J}}_0 \rightarrow \left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$$

Equivalence theorem



Consider a source distribution $\vec{\mathbf{J}}_0$ with its associated electromagnetic field $\left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$

Consider a (smooth) surface *A* with an everywhere defined unit normal $\hat{\mathbf{n}}$

The original sources $\vec{\mathbf{J}}_0$ enclosed in *A* can be removed and substituted by equivalent sources, i.e., electric $\vec{\mathbf{J}}_{es} = \hat{\mathbf{n}} \times \vec{\mathbf{H}}_0$ and magnetic $\vec{\mathbf{J}}_{ms} = -\hat{\mathbf{n}} \times \vec{\mathbf{E}}_0$ current densities distributed over the surface *A*.

$$\vec{\mathbf{J}}_0 \rightarrow \left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$$

Maxwell Equations (Spectral Domains)

Magnetic Sources

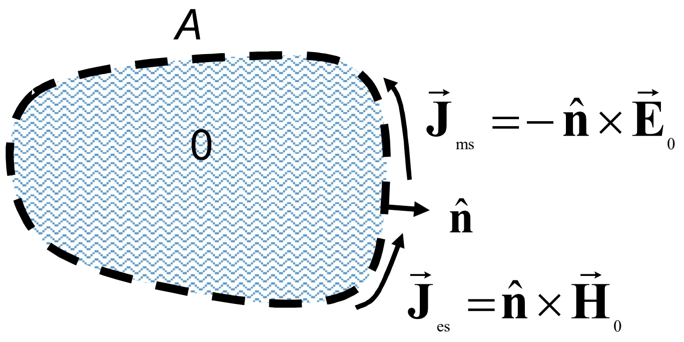
James Clerk Maxwell 1831-1879

$$\begin{cases} \nabla \times \vec{\mathbf{E}} = -j\omega\vec{\mathbf{B}} - \vec{\mathbf{J}}_m \\ \nabla \times \vec{\mathbf{H}} = j\omega\vec{\mathbf{D}} + \vec{\mathbf{J}} \\ \nabla \cdot \vec{\mathbf{D}} = \rho \\ \nabla \cdot \vec{\mathbf{B}} = \rho_m \end{cases}$$

$$\left[\vec{\mathbf{e}}(\vec{\mathbf{r}},t)\right]: \frac{Volt}{m} \qquad \left[\vec{\mathbf{b}}(\vec{\mathbf{r}},t)\right]: \frac{Weber}{m^2}$$

$$\left[\vec{\mathbf{j}}_m(\vec{\mathbf{r}},t)\right]: \frac{Volt}{m^2} \qquad \left[\rho_m(\vec{\mathbf{r}},t)\right]: \frac{Weber}{m^3}$$

Equivalence theorem

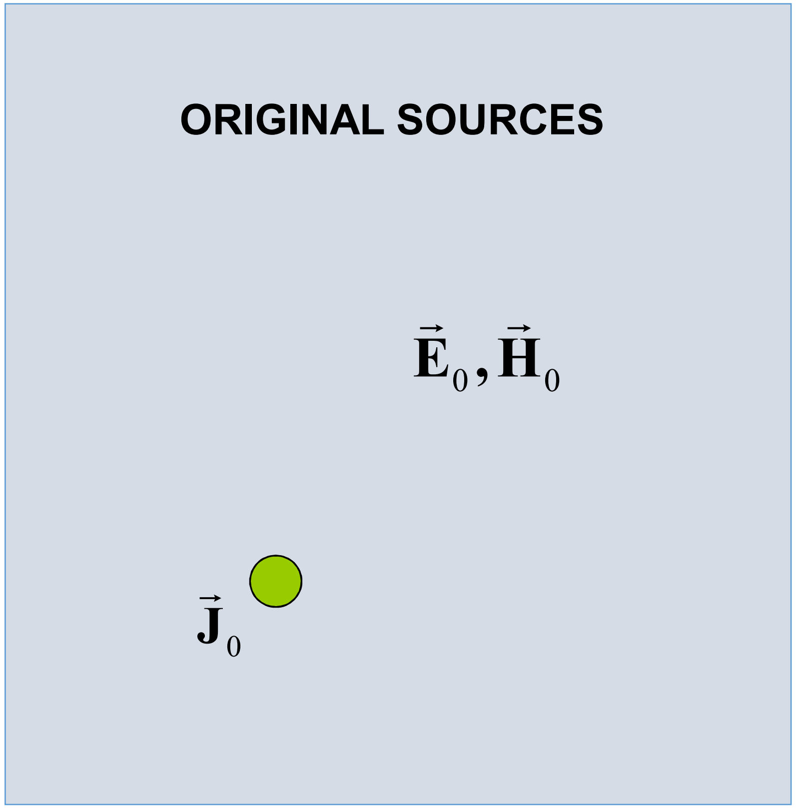


$$\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0$$

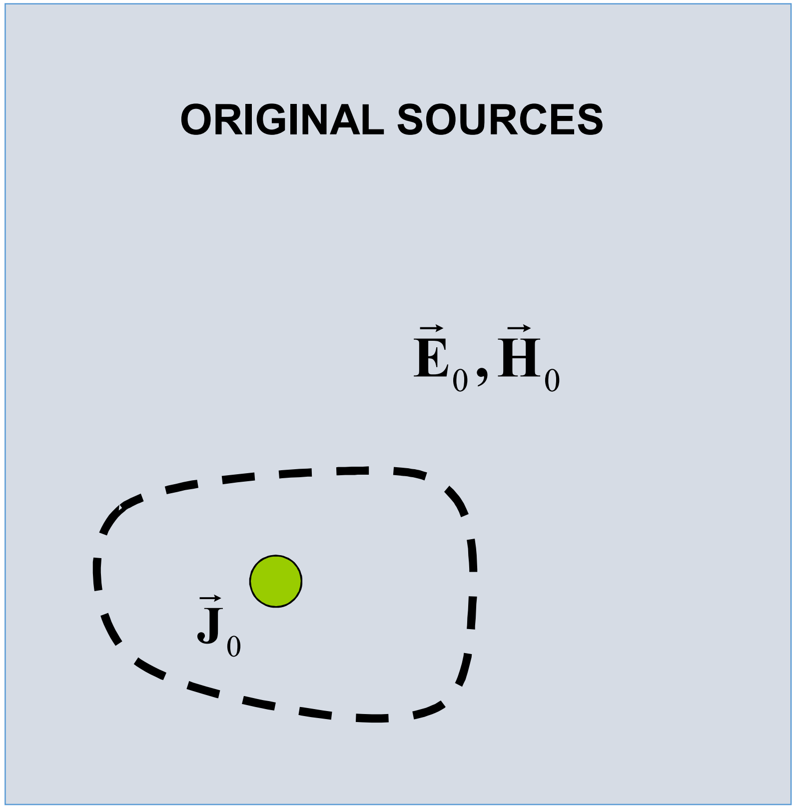
The Equivalence Theorem states that the equivalent sources $\vec{\mathbf{J}}_{es}$ and $\vec{\mathbf{J}}_{ms}$ generate a field $\left(\vec{\mathbf{E}}', \vec{\mathbf{H}}'\right)$ coincident with $\left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$ outside *A* and identically equal to zero inside

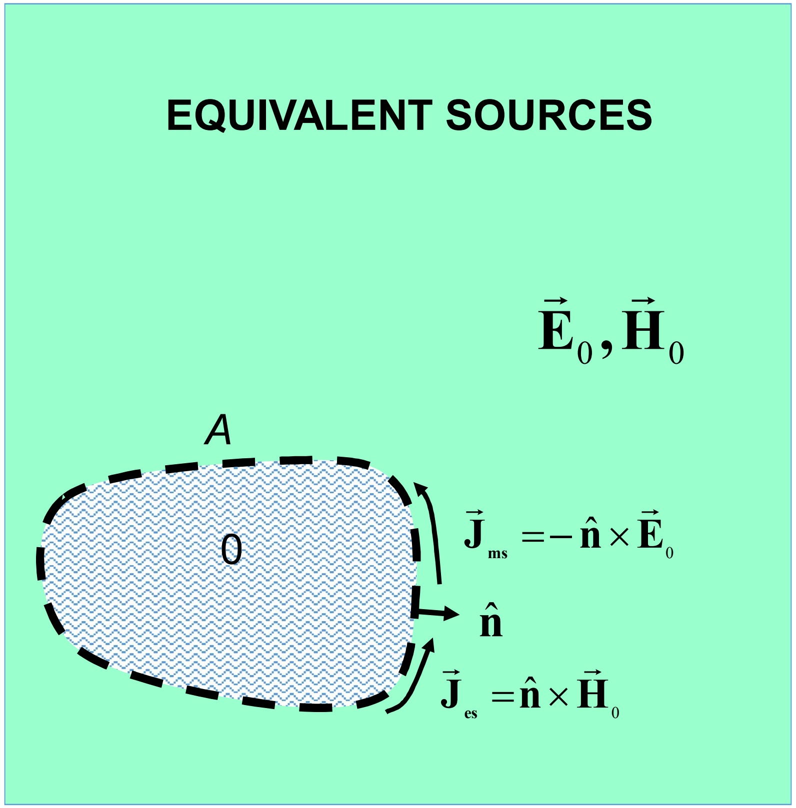
$$\vec{\mathbf{J}}_0 \rightarrow \left(\vec{\mathbf{E}}_0, \vec{\mathbf{H}}_0\right)$$

Equivalence theorem



Equivalence theorem

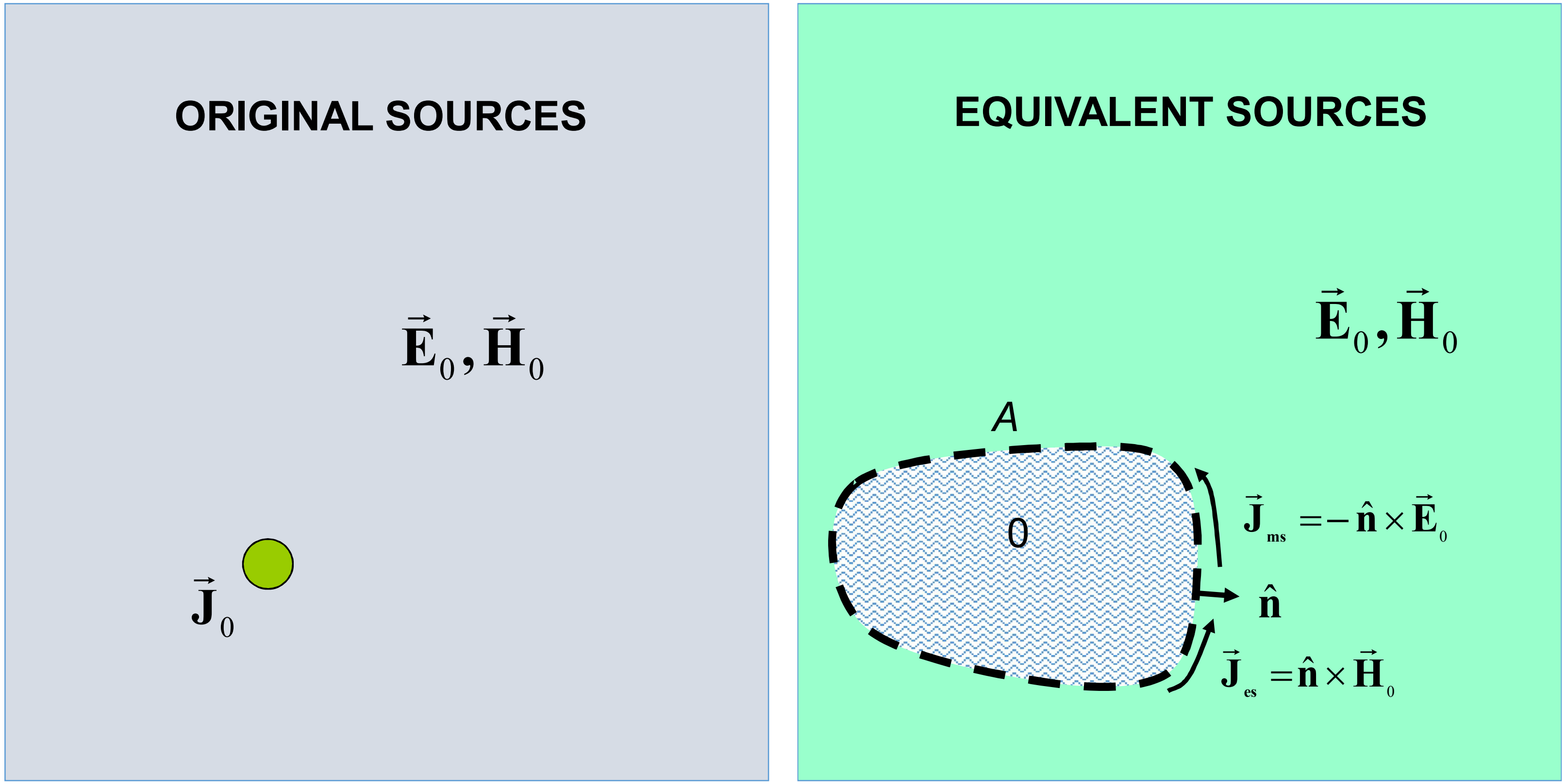




Equivalence theorem

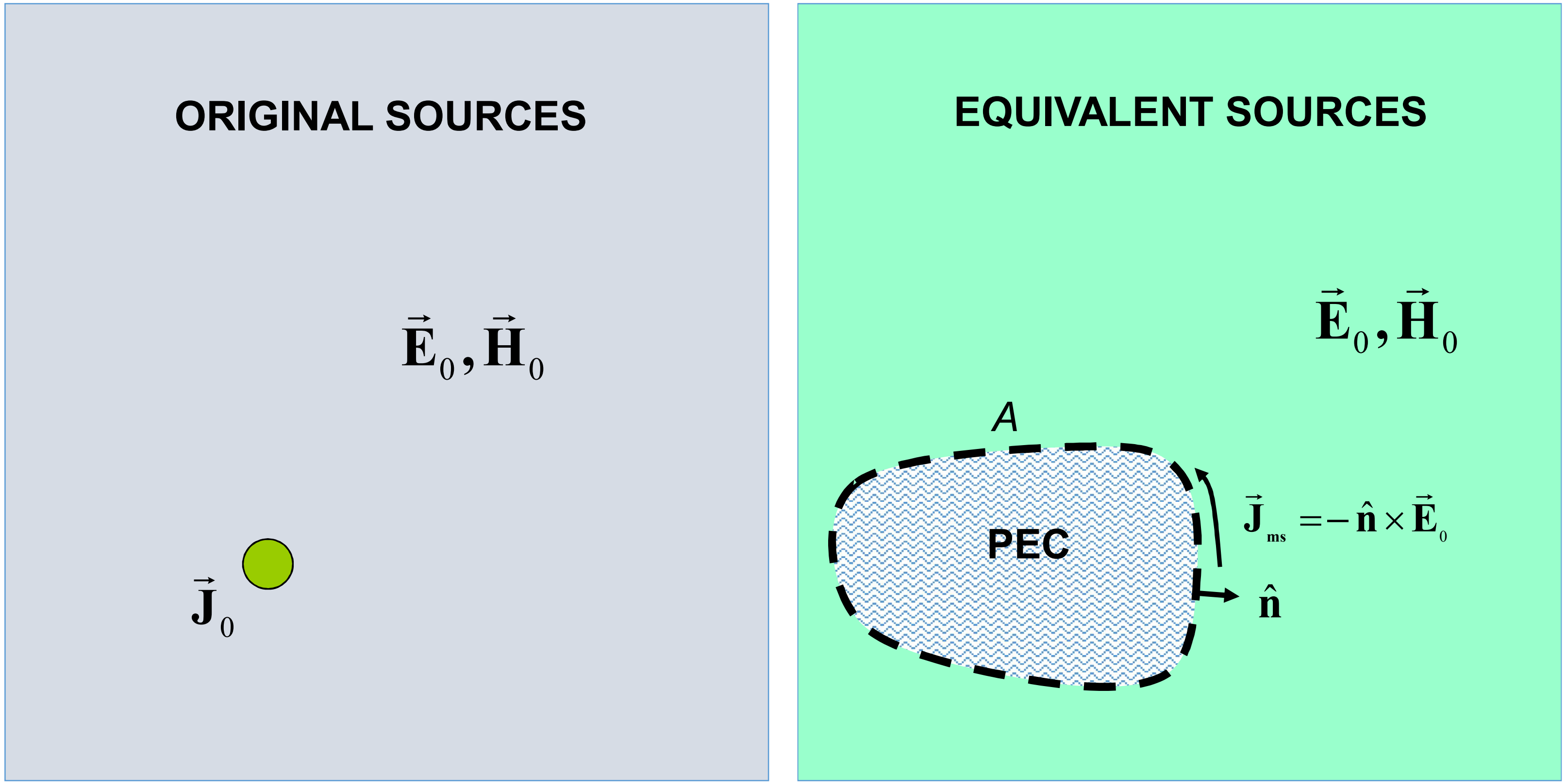
It's a powerful theorem that allows calculating the e.m. field in **all** the space, starting from the knowledge of its value **just on a surface**.

Equivalence theorem



Equivalence theorem

Alternative formulation



THEOREMS

Poynting

Time domain – Phasor domain

Uniqueness (Interior problem – Exterior problem)

Time domain – Phasor domain

Equivalence

Phasor domain

Image Theory

Reciprocity

Phasor domain

Image theory

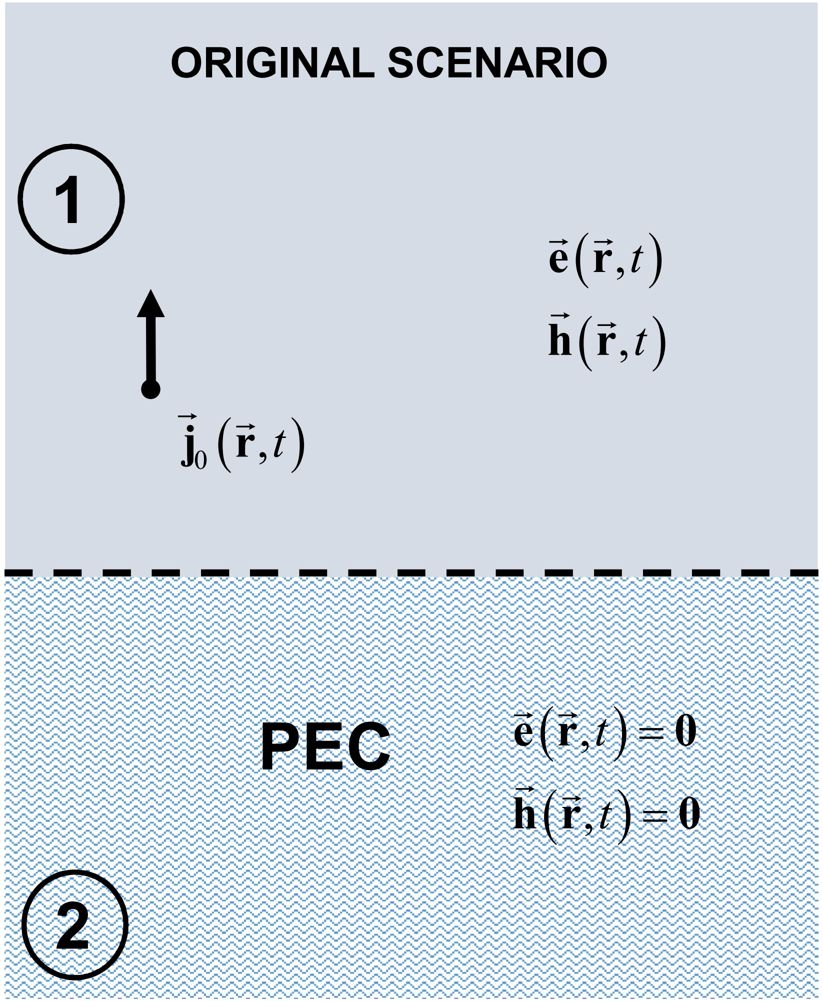
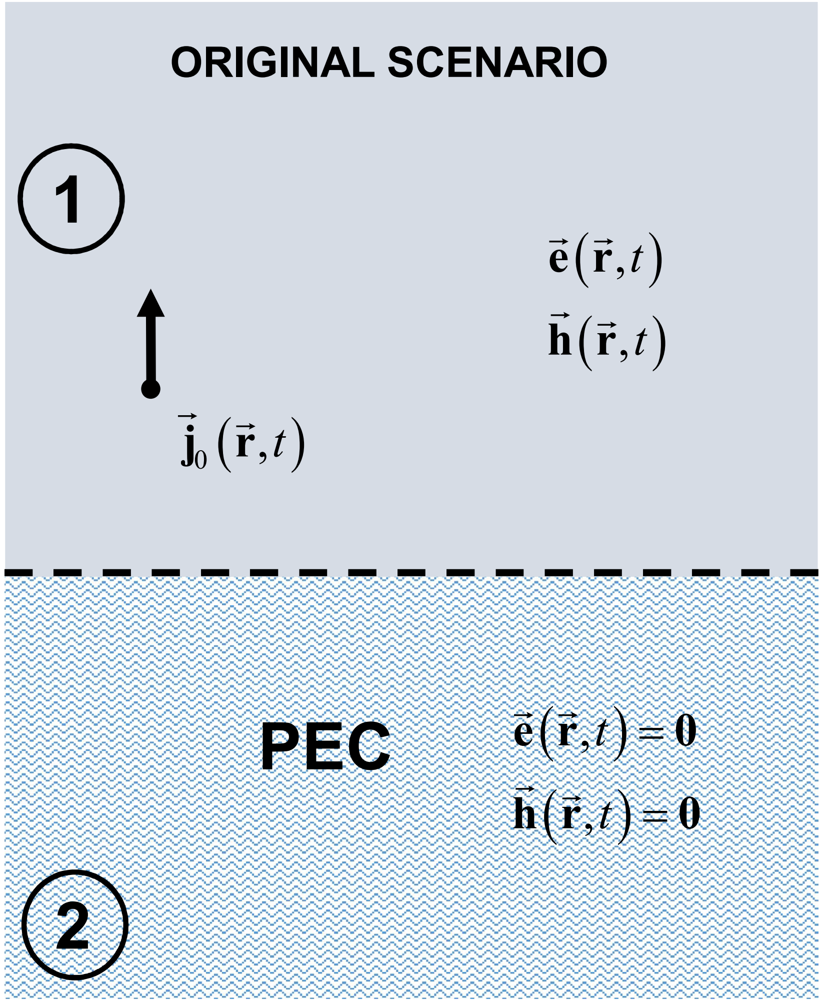


Image theory



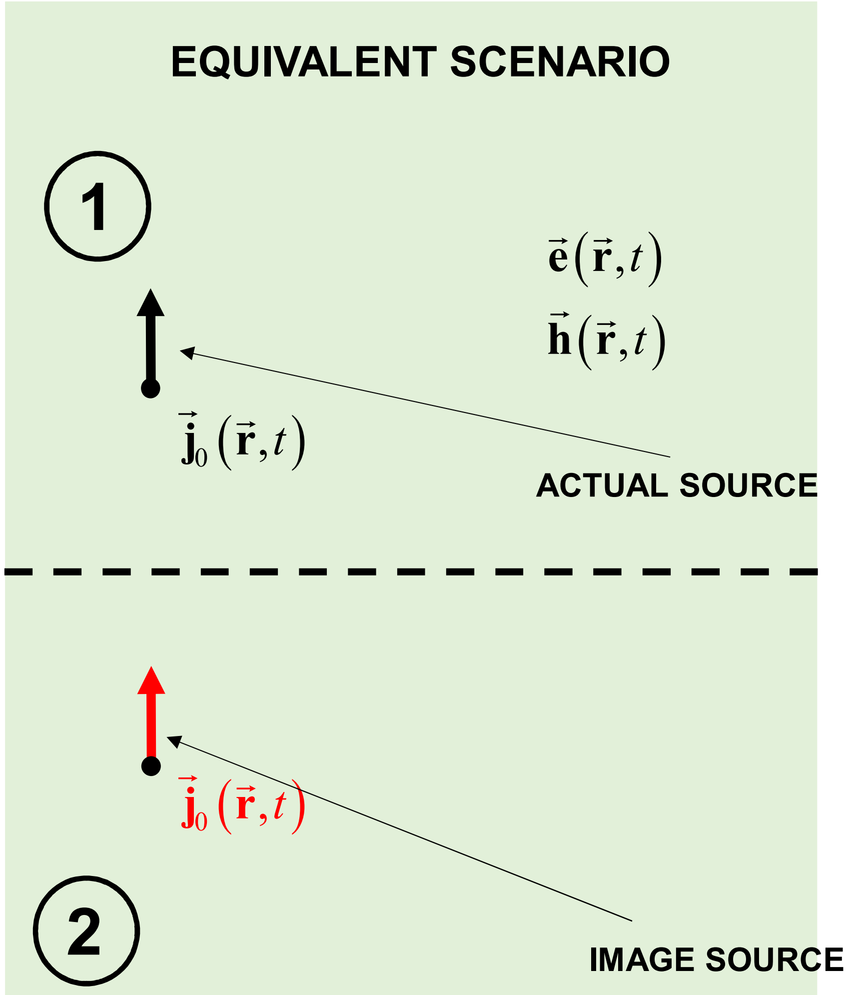


Image theory

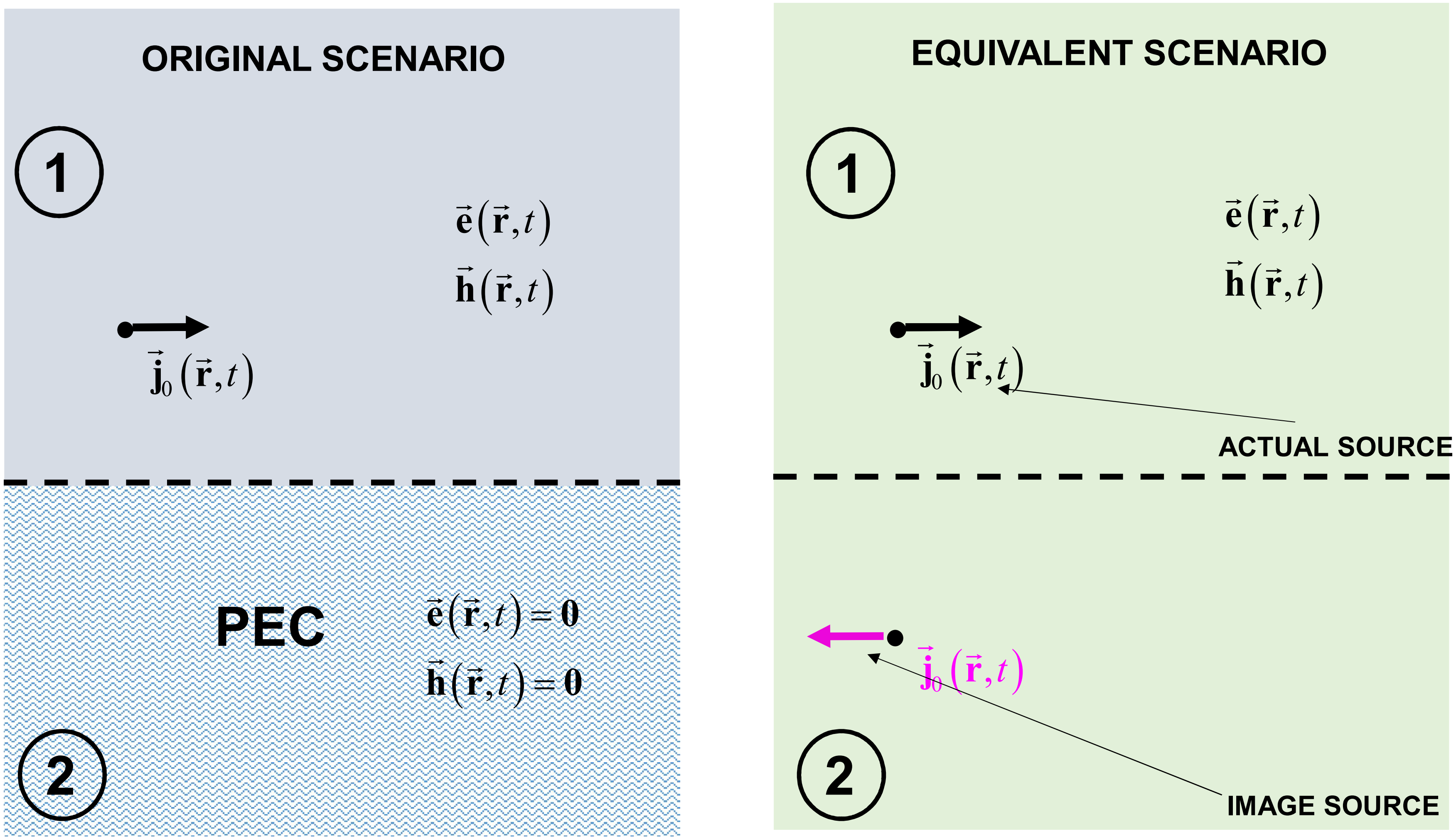
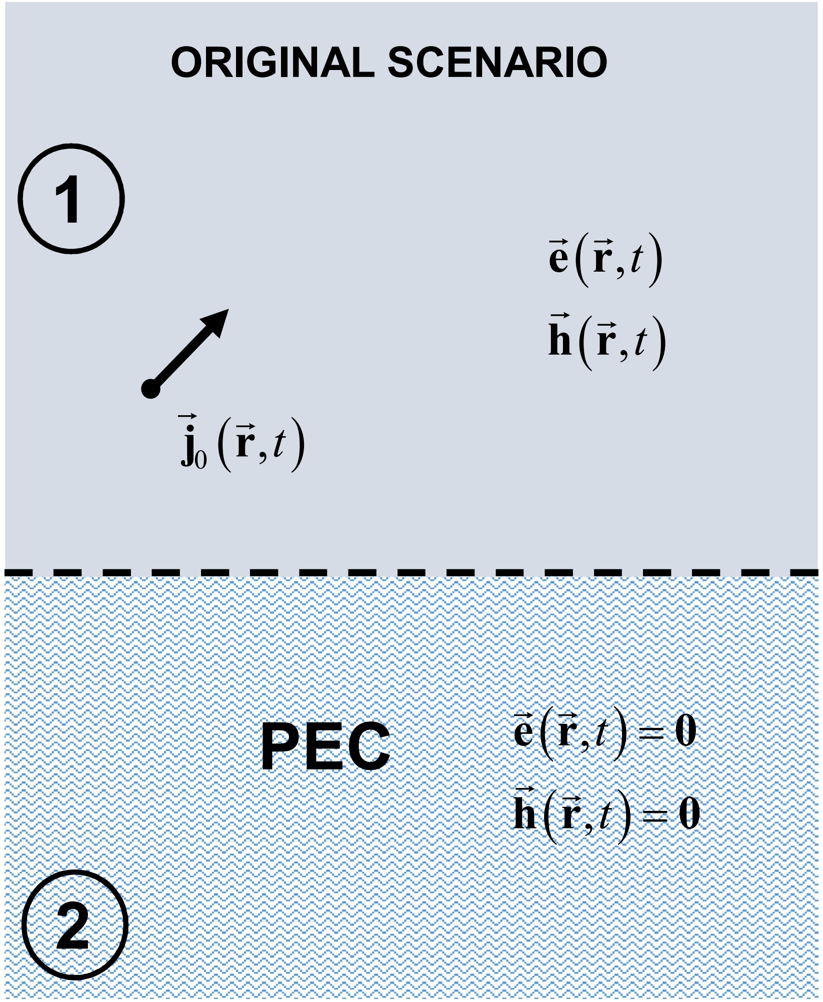


Image theory



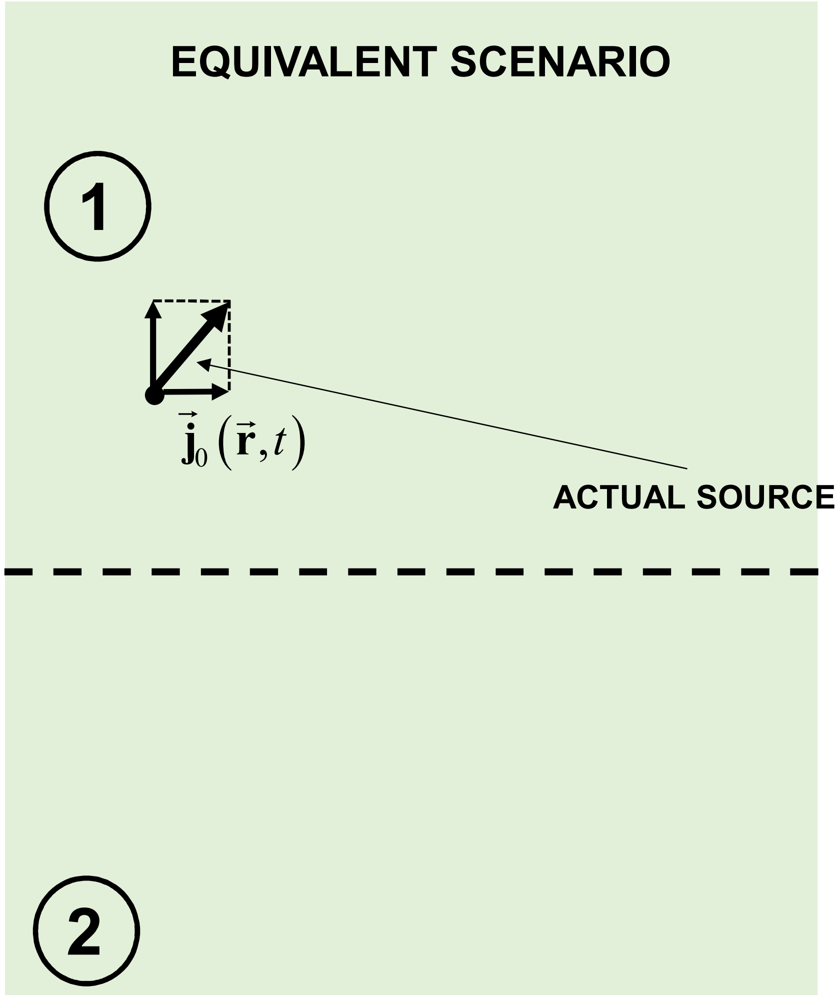


Image theory

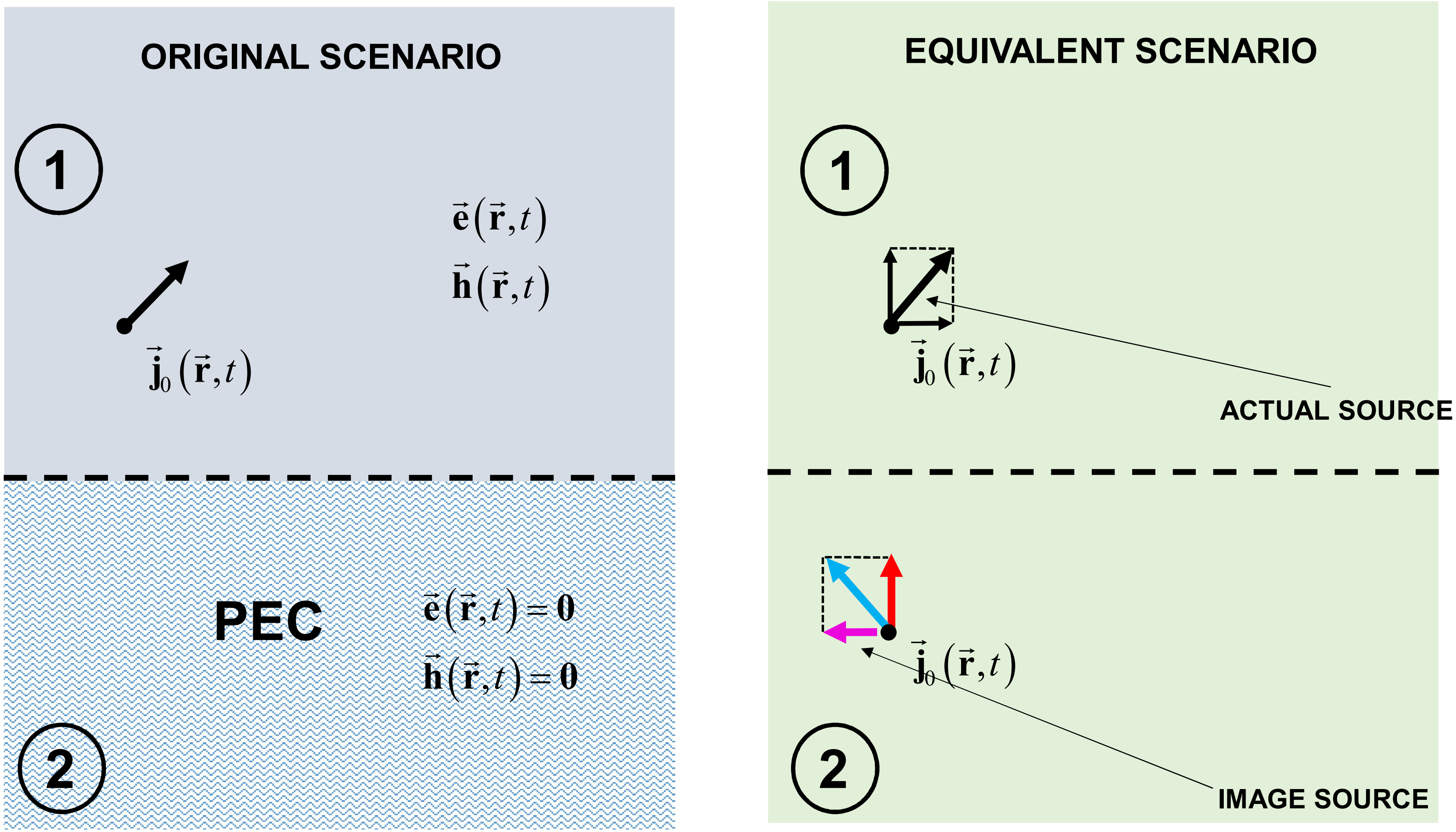


Image theory

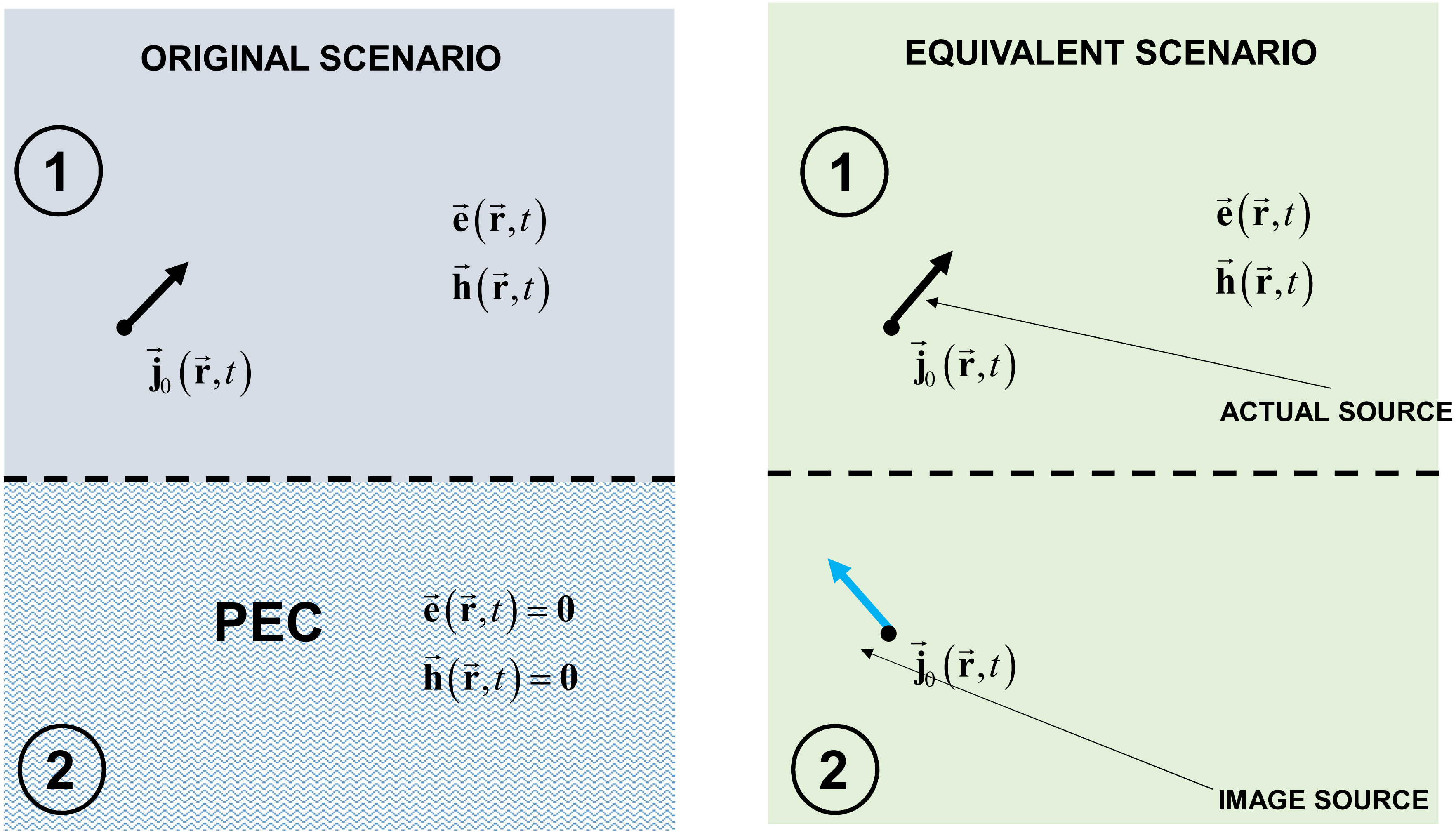
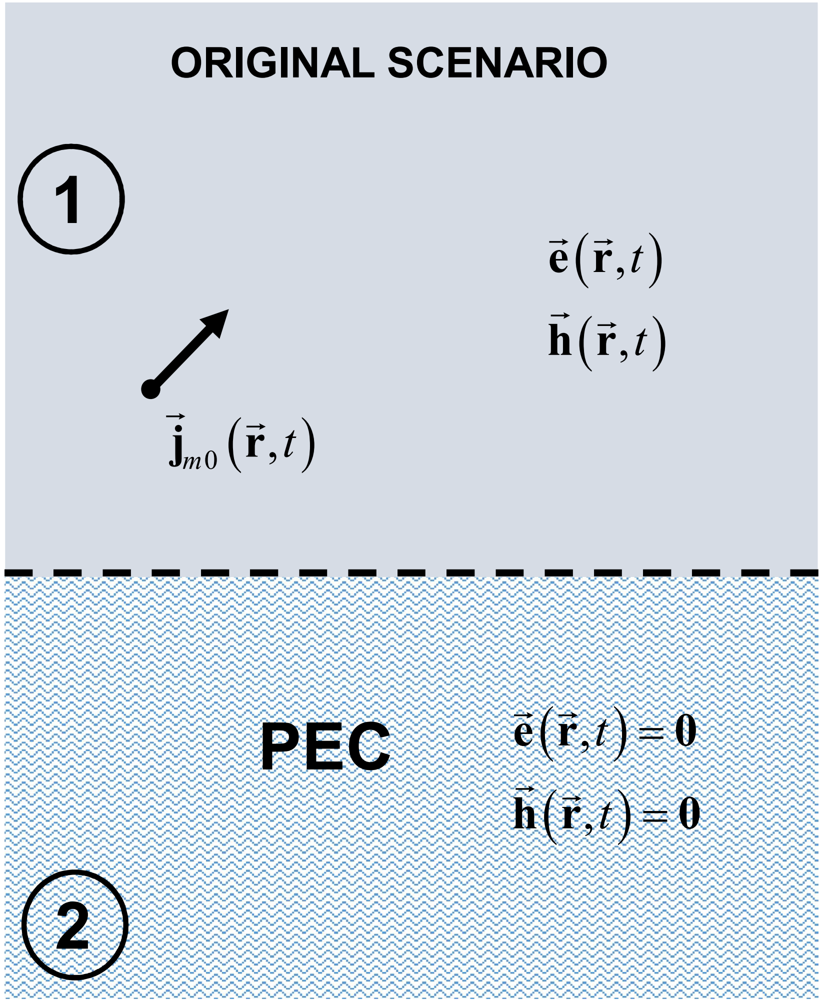
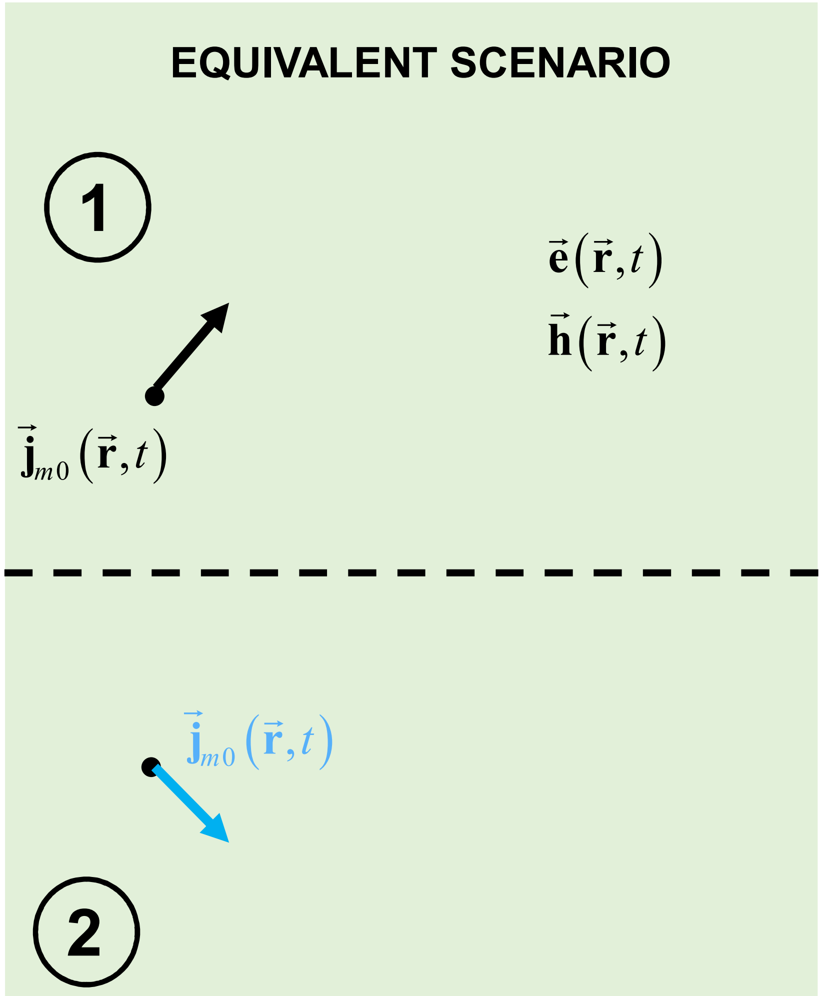


Image theory

Image theory (magnetic sources)





THEOREMS

Poynting

Time domain – Phasor domain

Uniqueness (Interior problem – Exterior problem)

Time domain – Phasor domain

Equivalence

Phasor domain

Image Theory

Reciprocity

Phasor domain

HOMEWORKS

Come si calcola il flusso di potenza irradiato da una sorgente?

Espressione del vettore di Poynting (TD e PD)

Espressione del vettore di Poynting all'infinito (TD e PD)

Il vettore di Poynting all'infinito nel PD è una quantità complessa?

PD: perchè le perdite sono legate a σ e a parte imaginaria di ε e μ?

Legame tra perdite e dispersione del mezzo

Il teorema di unicità (PD) è sempre valido?

Condizione di radiazione all'infinito