

# **Campi Elettromagnetici**

**Corso di Laurea in Ingegneria Informatica,  
Biomedica e delle Telecomunicazioni**

**a.a. 2023-2024 - Laurea “Triennale” – Secondo semestre - Secondo anno**

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# Color legend

New formulas, important considerations,  
important formulas, important concepts

Very important for the discussion

Memo

Mathematical tools to be exploited

Mathematics

# Mathematical tools that we will exploit today

$$\nabla \cdot (\nabla \times \vec{\mathbf{A}}(\vec{\mathbf{r}})) = 0$$

# Maxwell equations

Time domain - Differential form

$$\begin{cases} \nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t) = -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \\ \nabla \times \vec{\mathbf{h}}(\vec{\mathbf{r}}, t) = \frac{\partial \vec{\mathbf{d}}(\vec{\mathbf{r}}, t)}{\partial t} + \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{d}}(\vec{\mathbf{r}}, t) = \rho(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t) = 0 \end{cases}$$



**James Clerk Maxwell 1831-1879**

# Maxwell equations

Time domain - Differential form

$$\begin{cases} \nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t) = -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \\ \nabla \times \vec{\mathbf{h}}(\vec{\mathbf{r}}, t) = \frac{\partial \vec{\mathbf{d}}(\vec{\mathbf{r}}, t)}{\partial t} + \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) + \vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{d}}(\vec{\mathbf{r}}, t) = \rho(\vec{\mathbf{r}}, t) + \rho_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t) = 0 \end{cases}$$



**James Clerk Maxwell 1831-1879**

# Maxwell equations

$$\begin{cases} \nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t) = -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \\ \nabla \times \vec{\mathbf{h}}(\vec{\mathbf{r}}, t) = \frac{\partial \vec{\mathbf{d}}(\vec{\mathbf{r}}, t)}{\partial t} + \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) + \vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{d}}(\vec{\mathbf{r}}, t) = \rho(\vec{\mathbf{r}}, t) + \rho_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t) = 0 \end{cases}$$

$$\begin{cases} \vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t) \\ \rho_0(\vec{\mathbf{r}}, t) \end{cases} \quad \text{Prescribed sources}$$

$$\begin{cases} \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) \\ \rho(\vec{\mathbf{r}}, t) \end{cases} \quad \text{Induced sources}$$

## Complex scenario



# The independence of the Maxwell equations

Time domain - Differential form

$$\left\{ \begin{array}{l} \nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t) = -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \\ \nabla \times \vec{\mathbf{h}}(\vec{\mathbf{r}}, t) = \frac{\partial \vec{\mathbf{d}}(\vec{\mathbf{r}}, t)}{\partial t} + \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) + \vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{d}}(\vec{\mathbf{r}}, t) = \rho(\vec{\mathbf{r}}, t) + \rho_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t) = 0 \end{array} \right.$$

$$\nabla \cdot (\nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t)) = \nabla \cdot \left( -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \right)$$

↓

$$0 = -\frac{\partial (\nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t))}{\partial t}$$

↓

$$\nabla \cdot (\nabla \times \vec{\mathbf{A}}(\vec{\mathbf{r}})) = 0$$

# The independence of the Maxwell equations

Time domain - Differential form

$$\left\{ \begin{array}{l} \nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t) = -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \\ \nabla \times \vec{\mathbf{h}}(\vec{\mathbf{r}}, t) = \frac{\partial \vec{\mathbf{d}}(\vec{\mathbf{r}}, t)}{\partial t} + \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) + \vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{d}}(\vec{\mathbf{r}}, t) = \rho(\vec{\mathbf{r}}, t) + \rho_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t) = 0 \end{array} \right.$$

$$\nabla \cdot (\nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t)) = \nabla \cdot \left( -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \right)$$

↓   ↓

$$0 = -\frac{\partial (\nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t))}{\partial t}$$

$\nabla \cdot \vec{\mathbf{b}}$  is independent of time. If the fields are equal to zero before a given time, then  $\nabla \cdot \vec{\mathbf{b}} = 0$  for all times, thus recovering the last Maxwell equation

# The independence of the Maxwell equations

Time domain - Differential form

$$\left\{ \begin{array}{l} \nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t) = -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \\ \nabla \times \vec{\mathbf{h}}(\vec{\mathbf{r}}, t) = \frac{\partial \vec{\mathbf{d}}(\vec{\mathbf{r}}, t)}{\partial t} + \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) + \vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{d}}(\vec{\mathbf{r}}, t) = \rho(\vec{\mathbf{r}}, t) + \rho_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t) = 0 \end{array} \right.$$

Number of independent scalar equations:  
 $3+3+1=7$

Let us assume knowledge of the impressed sources  $\vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t); \rho_0(\vec{\mathbf{r}}, t)$

Number of unknown scalar quantities:  
 $\vec{\mathbf{e}}(\vec{\mathbf{r}}, t); \vec{\mathbf{d}}(\vec{\mathbf{r}}, t); \vec{\mathbf{h}}(\vec{\mathbf{r}}, t); \vec{\mathbf{b}}(\vec{\mathbf{r}}, t); \vec{\mathbf{j}}(\vec{\mathbf{r}}, t); \rho(\vec{\mathbf{r}}, t)$

3	+3	+3	+3	+3	+1
				16	

# The independence of the Maxwell equations

Time domain - Differential form

$$\begin{cases} \nabla \times \vec{e}(\vec{r}, t) = -\frac{\partial \vec{b}(\vec{r}, t)}{\partial t} \\ \nabla \times \vec{h}(\vec{r}, t) = \frac{\partial \vec{d}(\vec{r}, t)}{\partial t} + \vec{j}(\vec{r}, t) + \vec{j}_0(\vec{r}, t) \\ \nabla \cdot \vec{d}(\vec{r}, t) = \rho(\vec{r}, t) + \rho_0(\vec{r}, t) \\ \nabla \cdot \vec{b}(\vec{r}, t) = 0 \end{cases}$$

Even assuming knowledge of the impressed sources  $\vec{j}_0(\vec{r}, t); \rho_0(\vec{r}, t)$

Number of independent scalar equations: 7

Number of unknown scalar quantities: 16

Maxwell equations involve a number of unknowns larger than the number of equations!



# Constitutive relationships

## Time domain - Differential form

$$\begin{cases} \nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t) = -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \\ \nabla \times \vec{\mathbf{h}}(\vec{\mathbf{r}}, t) = \frac{\partial \vec{\mathbf{d}}(\vec{\mathbf{r}}, t)}{\partial t} + \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) + \vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{d}}(\vec{\mathbf{r}}, t) = \rho(\vec{\mathbf{r}}, t) + \rho_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t) = 0 \end{cases}$$

Even assuming knowledge of the impressed sources  $\vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t); \rho_0(\vec{\mathbf{r}}, t)$

Number of independent scalar equations: 7

Number of unknown scalar quantities: 16

Maxwell equations involve a number of unknowns larger than the number of equations!

The additional missing equations are provided by the **constitutive relationships**, which describe interaction of fields and matter from a macroscopic point of view

# Memo

## Vacuum

$$\vec{d}(\vec{r}, t) = \epsilon_0 \vec{e}(\vec{r}, t)$$

$\mu_0 = 4\pi \times 10^{-7}$  Permeability [Henry / m]

$$\vec{b}(\vec{r}, t) = \mu_0 \vec{h}(\vec{r}, t)$$

$\epsilon_0 = 8.8 \times 10^{-12}$  Permittivity [Farad / m]

$$\vec{j}(\vec{r}, t) = \sigma \vec{e}(\vec{r}, t)$$

$\sigma = 0$  Conductivity [Siemens / m]

These relationships depend on the particular medium that we have considered, that is, the vacuum.

More generally, similar relationships, can be found also in other media. They depend upon the characteristics of the medium in which the electromagnetic field is considered, and are called **CONSTITUTIVE RELATIONSHIPS**

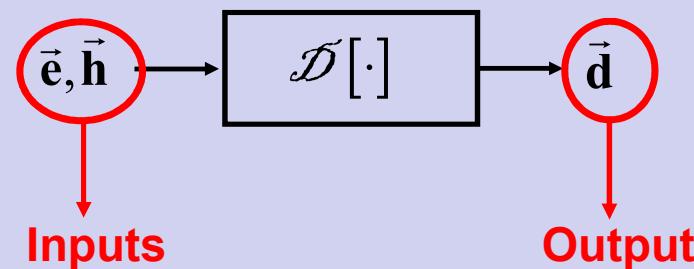
# Constitutive relationships

Inductions and currents must be represented in terms of fields

$$\vec{d} = \mathcal{D}[\vec{e}, \vec{h}]$$

$$\vec{b} = \mathcal{B}[\vec{e}, \vec{h}]$$

$$\vec{j} = \mathcal{J}[\vec{e}, \vec{h}]$$



$\mathcal{D}[\cdot]$ ,  $\mathcal{B}[\cdot]$  and  $\mathcal{J}[\cdot]$  are functionals which depend upon the medium in which the electromagnetic field is considered and upon the fields themselves.

# Constitutive relationships

## Linear media

$$\vec{d}_1 = \mathcal{D}[\vec{e}_1, \vec{h}_1] \quad \rightarrow \quad \vec{d}_1 + \vec{d}_2 = \mathcal{D}[\vec{e}_1 + \vec{e}_2, \vec{h}_1 + \vec{h}_2]$$
$$\vec{d}_2 = \mathcal{D}[\vec{e}_2, \vec{h}_2]$$

# Constitutive relationships

In the following we will consider linear media

# Constitutive relationships

## Linear media

### Example 1

$$\vec{d}(\vec{r}, t) = \boldsymbol{\epsilon} \cdot \vec{e}(\vec{r}, t) + \chi \cdot \vec{h}(\vec{r}, t)$$

where

$\vec{d}(\vec{r}, t)$ ;  $\vec{e}(\vec{r}, t)$ ;  $\vec{h}(\vec{r}, t)$ : 3x1 column vectors

$\boldsymbol{\epsilon}$  and  $\chi$ : 3x3 matrices

■ Local (nondispersive) media

■ Bianisotropic media

# Constitutive relationships

## Linear media

### Example 2

$$\vec{d}(\vec{r}, t) = \boldsymbol{\epsilon} \cdot \vec{e}(\vec{r}, t)$$

where

$\vec{d}(\vec{r}, t)$ ;  $\vec{e}(\vec{r}, t)$ ;  $\vec{h}(\vec{r}, t)$ : 3x1 column vectors

$\boldsymbol{\epsilon}$ : 3x3 matrix

- Local (nondispersive) media
- Anisotropic media

# Constitutive relationships

## Linear media

### Example 2

$$\vec{d}(\vec{r}, t) = \epsilon \cdot \vec{e}(\vec{r}, t)$$

$$\vec{b}(\vec{r}, t) = \mu \cdot \vec{h}(\vec{r}, t)$$

$$\vec{j}(\vec{r}, t) = \sigma \cdot \vec{e}(\vec{r}, t)$$

where

$\vec{d}(\vec{r}, t); \vec{b}(\vec{r}, t); \vec{j}(\vec{r}, t); \vec{e}(\vec{r}, t); \vec{h}(\vec{r}, t)$ : 3x1 column vectors

$\epsilon; \mu; \sigma$ : 3x3 matrices

■ Local (nondispersive) media

■ Anisotropic media

$\epsilon$ : permittivity [Farad/m]

$\mu$ : permeability [Henry/m]

$\sigma$ : conductivity [Siemens/m]

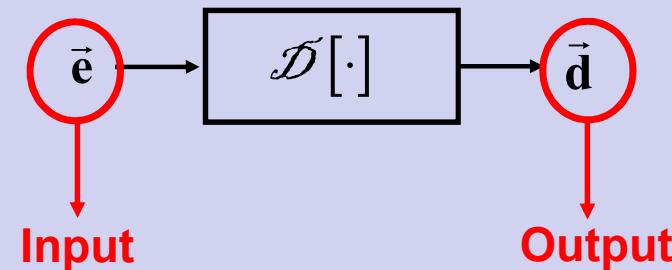
# Constitutive relationships

## Linear & Anisotropic media

$$\vec{d} = \mathcal{D}[\vec{e}]$$

$$\vec{b} = \mathcal{B}[\vec{h}]$$

$$\vec{j} = \mathcal{J}[\vec{e}]$$



# Constitutive relationships

## Linear media

$$\vec{d}(\vec{r}, t) = \epsilon \cdot \vec{e}(\vec{r}, t)$$

$\epsilon$ : 3x3 matrix

■ Local (nondispersive) media

■ Anisotropic media

### Class

Isotropic

### Property

A **rotation** of the input implies  
the **same rotation** of the  
output

### Effect on the I-O relation

$\epsilon$  becomes scalar. It is not a  
matrix anymore!

# Constitutive relationships

## Linear media

- Local (non-dispersive) media
- Anisotropic media

$$\vec{d}(\vec{r}, t) = \epsilon \cdot \vec{e}(\vec{r}, t)$$

$$\vec{b}(\vec{r}, t) = \mu \cdot \vec{h}(\vec{r}, t)$$

$$\vec{j}(\vec{r}, t) = \sigma \cdot \vec{e}(\vec{r}, t)$$

$\epsilon; \mu; \sigma$ : 3x3 matrices

- Local (non-dispersive) media
- Isotropic media

$$\vec{d}(\vec{r}, t) = \epsilon \vec{e}(\vec{r}, t)$$

$$\vec{b}(\vec{r}, t) = \mu \vec{h}(\vec{r}, t)$$

$$\vec{j}(\vec{r}, t) = \sigma \vec{e}(\vec{r}, t)$$

$\epsilon; \mu; \sigma$ : scalars

# Constitutive relationships

## Linear media

**Class**

**Property**

**Effect on the I-O relation**

# Constitutive relationships

## Linear media

### Class

Time-dispersive

Time-nondispersive

Time-variant

Time-invariant

### Property

### Effect on the I-O relation

# Constitutive relationships

## Linear media

### Class

**Space-dispersive**

**Space-nondispersive**

**Space-variant**

**Space-invariant**

### Property

### Effect on the I-O relation

# Constitutive relationships

## Linear media

### Class

Time-dispersive

Time-nondispersive

Time-variant

Time-invariant

### Property

The output **at time  $t$  depends only** on the value of the input **at the same time  $t$**

# Constitutive relationships

## Linear media

### Class

Time-dispersive

Time-nondispersive

Time-variant

Time-invariant

### Property

The output at time  $t$  depends on the values of the input throughout a time-interval.

# Constitutive relationships

## Linear media

### Class

Time-dispersive

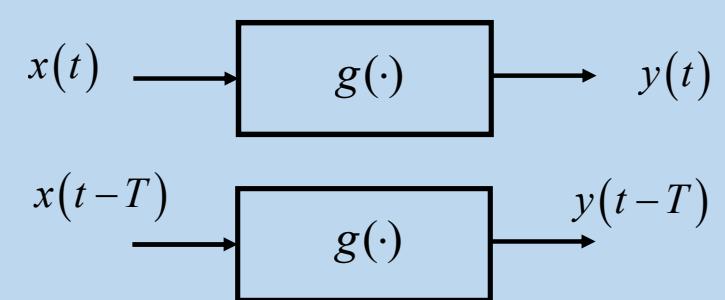
Time-nondispersive

Time-variant

Time-invariant

### Property

A time translation of the input implies the same translation of the output



# Constitutive relationships

## Linear media

### Class

Time-dispersive

Time-nondispersive

Time-variant

Time-invariant

### Property

A time translation of the input does not imply the same translation of the output

# Constitutive relationships

## Linear media

### Class

Space-dispersive

Space-nondispersive

Space-variant

Space-invariant

### Property

The output at space  $\vec{r}$   
depends only on the value of  
the input at the same space  $\vec{r}$

# Constitutive relationships

## Linear media

### Class

**Space-dispersive**

**Space-nondispersive**

**Space-variant**

**Space-invariant**

### Property

The output **at space  $\vec{r}$**   
**depends** on the values of  
the input **throughout a**  
**space-interval**

# Constitutive relationships

## Linear media

### Class

**Space-dispersive**

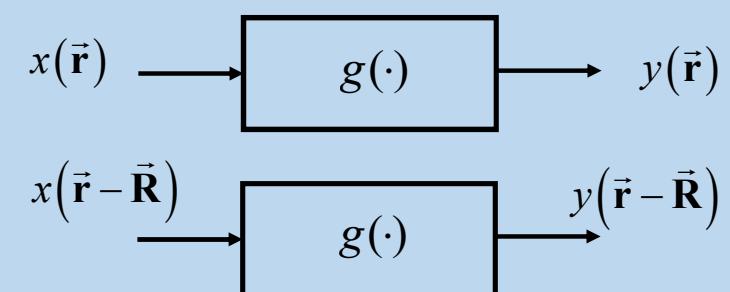
**Space-nondispersive**

**Space-variant**

**Space-invariant**

### Property

A **space translation** of the input implies **the same translation** of the output



# Constitutive relationships

## Linear media

### Class

**Space-dispersive**

**Space-nondispersive**

**Space-variant**

**Space-invariant**

### Property

A **space translation** of the input **does not imply the same translation** of the output

# Constitutive relationships

## Linear media

**Class**

**Property**

**Effect on the I-O relation**

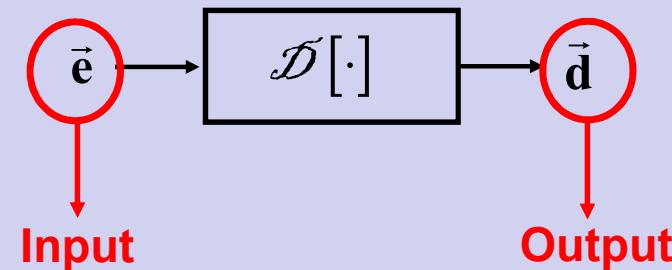
# Constitutive relationships

## Linear & Anisotropic media

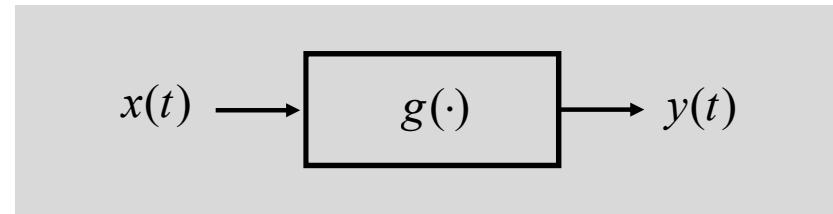
$$\vec{d} = \mathcal{D}[\vec{e}]$$

$$\vec{b} = \mathcal{B}[\vec{h}]$$

$$\vec{j} = \mathcal{J}[\vec{e}]$$



# Linear systems: effect on the I-O relation



$x(t)$  and  $y(t)$  are scalar

	Time-dispersive (TD)	Time-nondispersive (TND)
Time-variant (TV)	$y(t) = \int dt' g(t, t') x(t')$	$y(t) = \tilde{g}(t) x(t)$
Time-invariant (TI)	$y(t) = \int dt' g(t - t') x(t')$	$y(t) = \tilde{g} x(t)$