



Course of
"Automatic Control Systems"
2022/23

Nyquist stability criterion

Prof. Francesco Montefusco

Department of Economics, Law, Cybersecurity, and Sports Sciences

Università degli Studi di Napoli Parthenope

francesco.montefusco@uniparthenope.it

Team code: **uxbsz19**



Phase variation formula

- ✦ The previous lesson the concept of phase variation has been introduced.
- ✦ The phase variation is related to the number and sign of poles/zeros of the transfer function $F(s)|_{s=j\omega}$.

Given a transfer function $F(s)|_{s=j\omega}$, said:

- n the total number of poles
- m the total number of zeros
- n_p the number of poles with positive real part
- m_p the total number of zeros with positive real part

$$\Delta \angle F(j\omega) \Big|_{-\infty}^{\infty} = \pi(m - n) - 2\pi(m_p - n_p)$$



Phase variation formula

- ✦ The previous phase variation formula doesn't consider the case of poles and zeros on the imaginary axis.
- ✦ Indeed, in case of poles and zeros on the imaginary axis the phase variation can not be defined
- ✦ In the following slides we will consider these two critical cases and we will illustrate how to extend the definition of phase variation



Phase variation with poles on the imaginary axis

✧ Open loop poles on the imaginary axis (null real part), can be due to:

✧ One or more integrators $1/s^h$

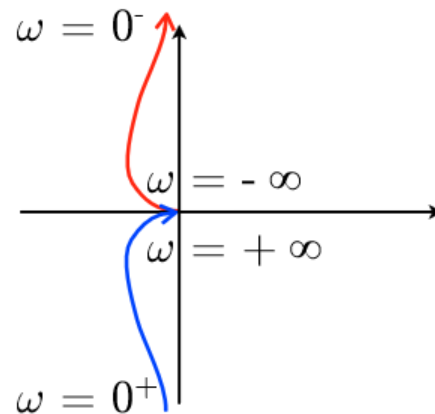
✧ Resonance $1/(1 + s^2/\omega_n^2)^h$

✧ In both the cases we have a *discontinuity in the phase margin*:

✧ **Integrator**: passing *from $\pi/2$ to $-\pi/2$* with infinite magnitude *at $\omega = 0$*

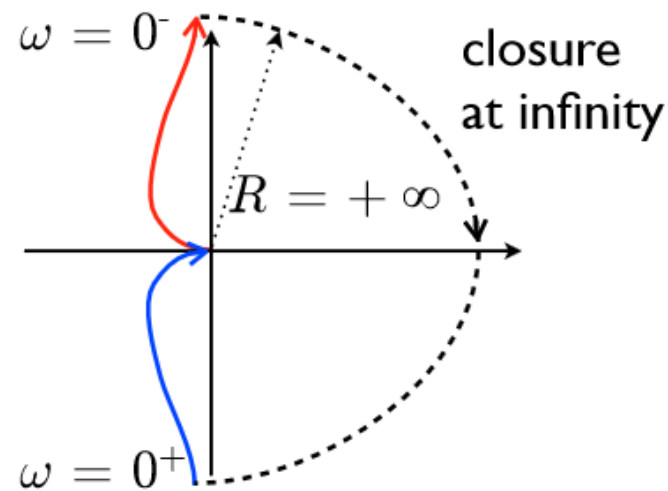
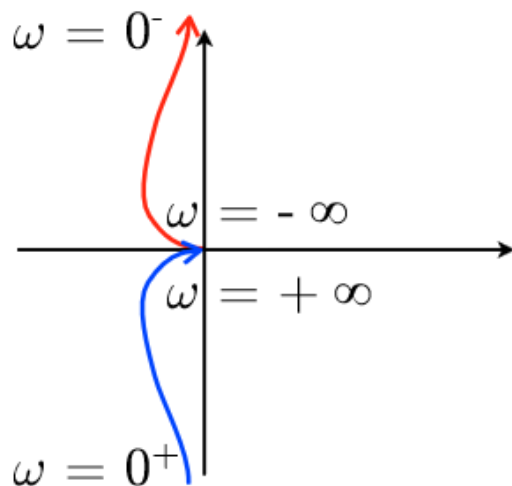
✧ **Resonance**: passing *from 0 to $-\pi$* with infinity magnitude *at $\omega = \omega_n$* and *from π to 0* with infinity magnitude *at $\omega = -\omega_n$*

$$F(s) = \frac{1}{s(s+1)}$$



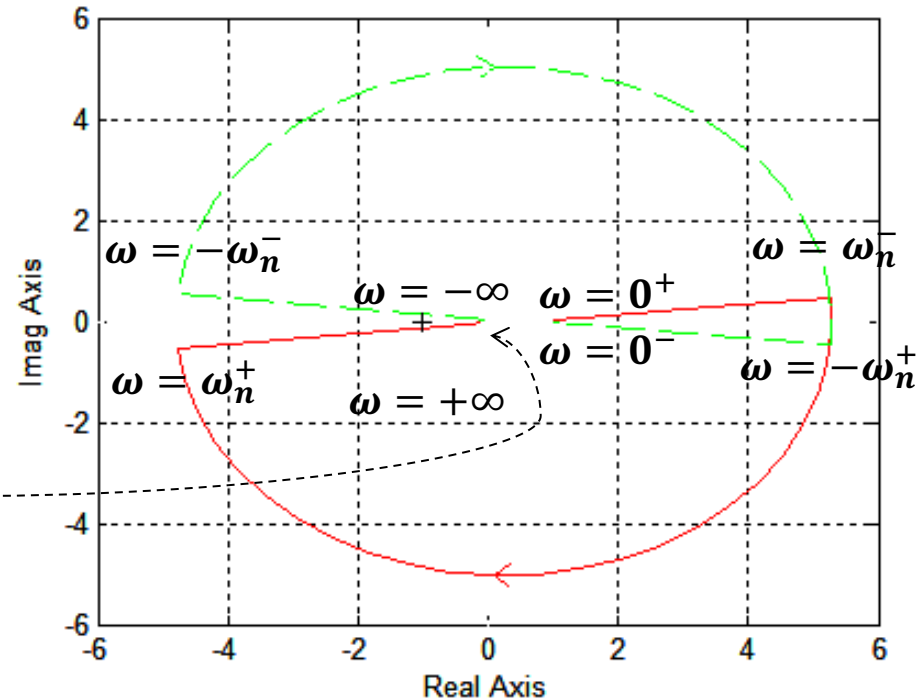
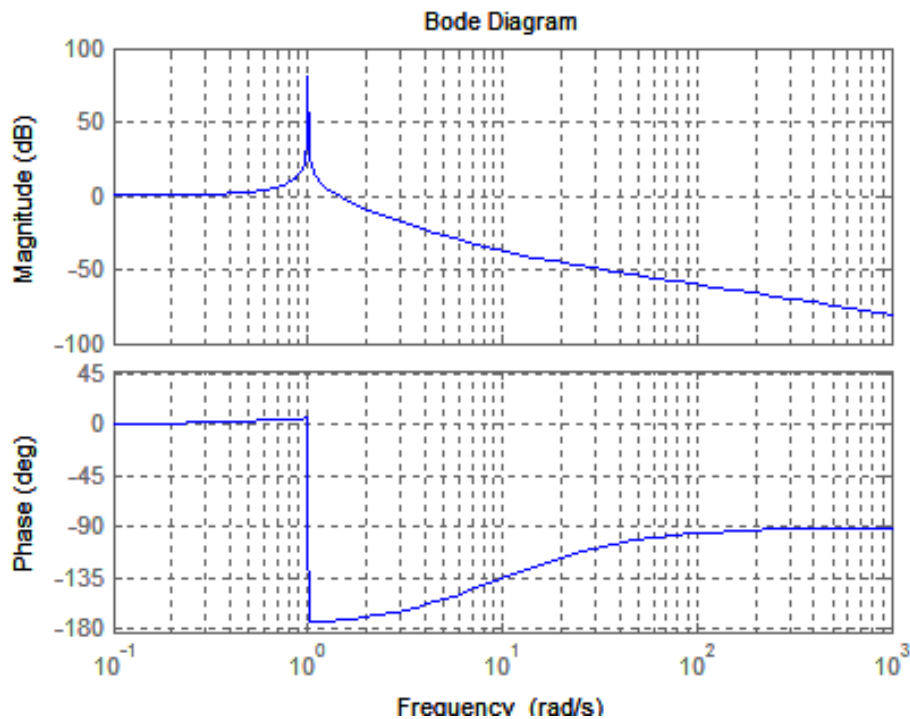
- ▶ In order to obtain a closed polar plot and to extend the definition of phase variation, we introduce *the closures at infinity*.
- ▶ The closures at infinity consists in rotating clockwise the Nyquist plot of the $F(j\omega)$ in the discontinuity frequency with an infinite radius.
- ▶ With this manipulation, the contribution to the phase variation of poles on the imaginary axis will be the same as the poles with negative real part.

$$F(s) = \frac{1}{s(s+1)}$$



Transfer function with resonance

$$F(s) = \frac{1 + 0.1s}{1 + s^2}$$

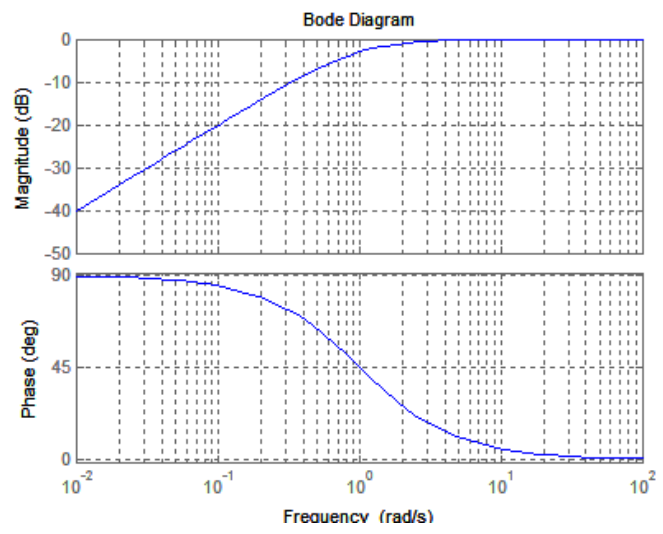


$$\Delta \angle F(j\omega) \Big|_{-\infty}^{\infty} = \pi(m - n) - 2\pi(m_p - n_p) = \pi(1 - 2) = -\pi$$

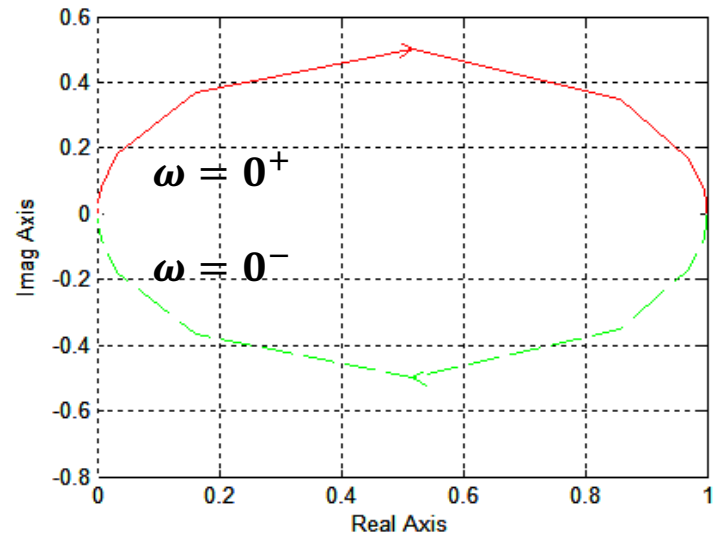


Phase variation with zeros on the imaginary axis

- Open loop zeros on the imaginary axis (null real part), can be due to:
 - One or more derivative s^h
 - Anti-resonance $(1 + s^2/\omega_n^2)^h$
- In both the cases we have a *discontinuity in the phase margin*:
 - Derivate**: passing *from $-\pi/2$ to $\pi/2$* with zero magnitude *at $\omega = 0$*
 - Anti-Resonance**: passing *from 0 to π* with zero magnitude *at $\omega = \omega_n$* and *from $-\pi$ to 0* with zero magnitude *at $\omega = -\omega_n$*



$$F(s) = \frac{s}{1+s}$$



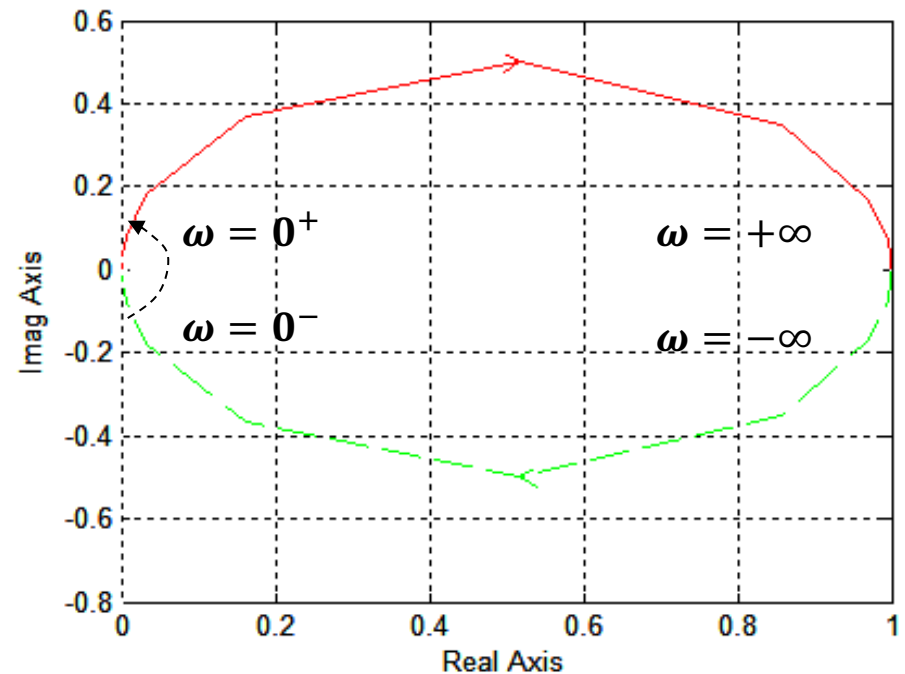


Phase variation with zeros on the imaginary axis

- ✦ In order to extend the definition of phase variation, we will assume that *in* $\omega = 0$ and $\omega = \omega_n$ the Nyquist plot of the frequency response $F(j\omega)$ will rotate *counterclockwise* with in infinitesimal magnitude.
- ✦ With this manipulation, the contribution to the phase variation of zeros on the imaginary axis will be the same as the zeros with negative real part.

$$F(s) = \frac{s}{1+s}$$

$$\Delta \angle F(j\omega) \Big|_{-\infty}^{\infty} = 0$$



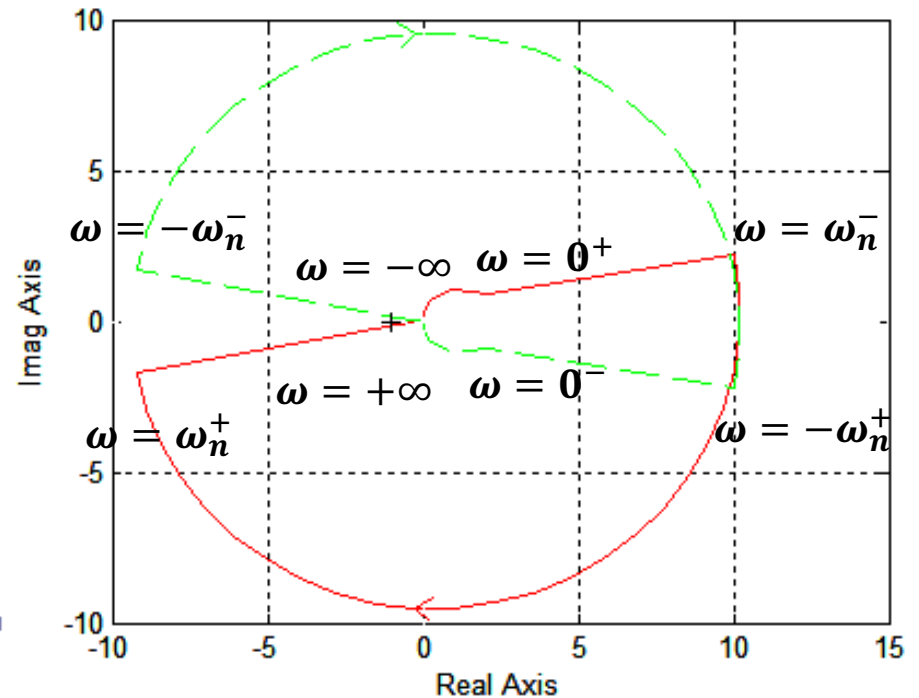
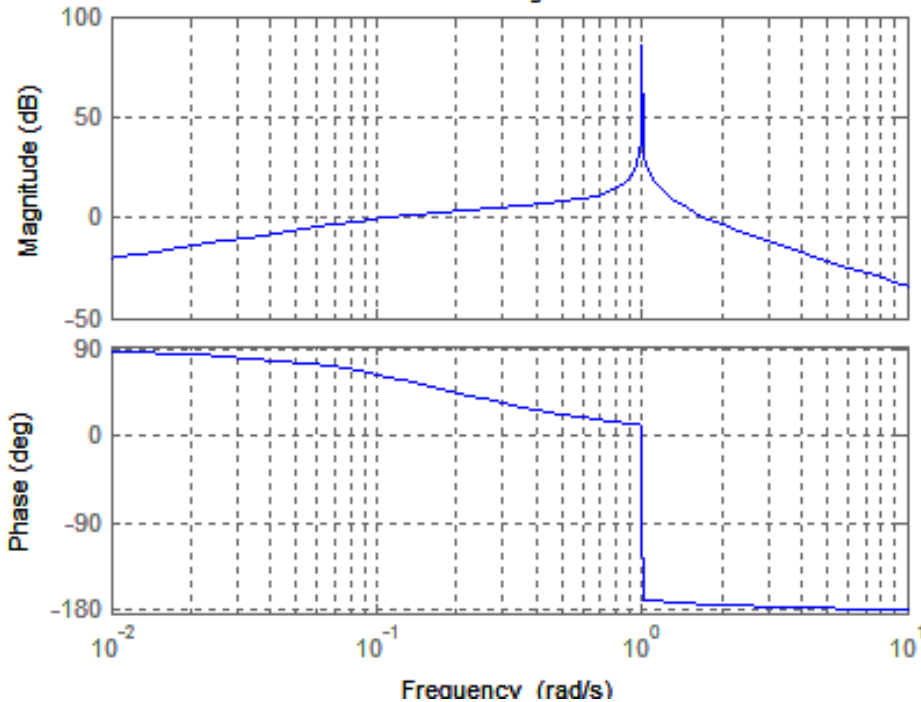


Phase variation with poles and zeros on the imaginary axis

Transfer function with resonance

$$F(s) = \frac{10s}{(1 + 5s)(1 + s^2)}$$

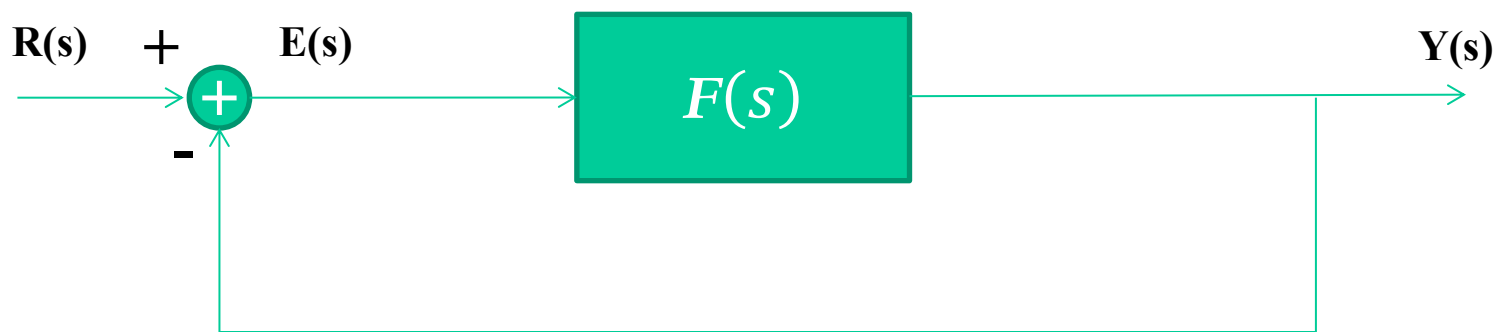
Bode Diagram



$$\Delta \angle F(j\omega)_{-\infty \rightarrow \infty} = \pi(m - n) - 2\pi(m_p - n_p) = \pi(1 - 3) = -2\pi$$

Stability of the closed loop system

- Let us consider the $R(s) \rightarrow Y(s)$ closed loop system



- Assume that the **hidden modes** of the open loop function $F(s) = K(s)G(s)$ are asymptotically stable
- The stability of the closed loop system depends on the poles of the transfer function

$$T(s) = \frac{F(s)}{1+F(s)}$$



Stability of the closed loop system

- ✦ Indicate with $N_F(s)$ and $D_F(s)$ the numerator and the denominator of the open loop function

$$F(s) = \frac{N_F(s)}{D_F(s)} = \frac{N_K(s)N_G(s)}{D_K(s)D_G(s)}$$

- ✦ The closed loop function can be written as

$$T(s) = \frac{\frac{N_F(s)}{D_F(s)}}{1 + \frac{N_F(s)}{D_F(s)}} = \frac{N_F(s)}{D_F(s) + N_F(s)}$$

- ✦ The denominator of $T(s)$ is given by the sum of $N_F(s)$ and $D_F(s)$; therefore the design problem of a controller able to guarantee the stability of the closed loop system is to be very demanding.
- ✦ By means of the Nyquist plots and the *Nyquist criteria*, we are going *to determine the stability of the closed loop* system *from the open loop* system features



Nyquist Stability Criterion

Let us consider a strictly proper open-loop function $F(s)$ and assume that the Nyquist diagram of $F(s)$ doesn't intersect the critical point $-1 + j0$.

Said

- ✦ $\overleftarrow{\mathcal{N}}$ the number of counter-clockwise encirclements of the critical point $-1 + j0$ of the Nyquist plot of $F(s)$
- ✦ $n_{p+}(F(s))$ the number of unstable poles of $F(s)$

the closed loop function $T(s)$ is asymptotically stable if and only if

$$\overleftarrow{\mathcal{N}} = n_{p+}(F(s)).$$

Moreover, if $\overleftarrow{\mathcal{N}} \neq n_{p+}$, the number of unstable poles of the closed loop function $T(s)$ is equal to $n_{p+}(F(s)) - \overleftarrow{\mathcal{N}}$.



Nyquist Stability Criterion: proof

- ✦ Let us define the so-called *Difference Function*

$$D(s) = 1 + F(s) = \frac{N_F(s) + D_F(s)}{D_F(s)}$$

- ✦ It is straightforward to notice that :

- ✦ The poles of $D(s)$ are the open loop control system poles , i.e. $D_D(s) = D_F(s)$

- ✦ The zeros of $D(s)$ are the closed loop control system poles, i.e. $N_D(s) = D_T(s)$



Nyquist Stability Criterion: proof

♣ Said

♣ $n_p(\mathbf{D}(s))$ the number of poles of $D(s)$

♣ $n_{p+}(\mathbf{D}(s))$ the number of poles with positive real part of $D(s)$

♣ $n_z(\mathbf{D}(s))$ the number of zeros of $D(s)$

♣ $n_{z+}(\mathbf{D}(s))$ the number of zeros with positive real part of $D(s)$

The phase variation of the difference function is

$$\Delta \angle D(j\omega)_{-\infty \omega \infty} = \pi(n_z(\mathbf{D}(s)) - n_p(\mathbf{D}(s))) - 2\pi(n_{z+}(\mathbf{D}(s)) - n_{p+}(\mathbf{D}(s)))$$



Nyquist Stability Criterion: proof

✦ However,

1. $F(s)$ strictly proper $\rightarrow D(s) = 1 + F(s) = \frac{N_F(s) + D_F(s)}{D_F(s)}$ proper and than

$$n_p(D(s)) = n_z(D(s))$$

2. Taking into account that $N_D(s) = D_T(s)$ and it is required the closed loop stability of the system, than

$$n_{z^+}(D(s)) = 0$$

✦ Hence the phase variation of the difference function is

$$\Delta \angle D(j\omega) = 2\pi \cdot n_{p^+}(D(s)) = 2\pi \cdot n_{p^+}(F(s))$$

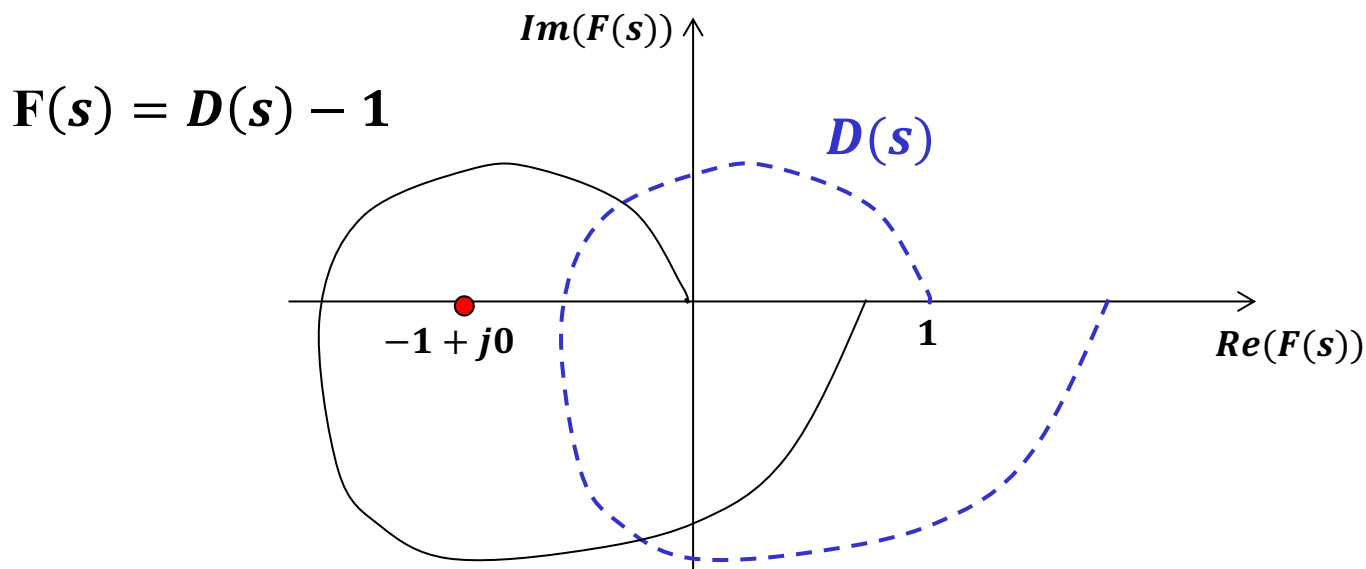


Nyquist Stability Criterion: proof

- ✦ The function $D(s)$ will encircle counter-clockwise the origin of the Nyquist plane a number of times given by

$$n_{p+}(F(s))$$

- ✦ The proof is concluded taking into account that the encirclements of the origin of the $D(s)$ Nyquist plot correspond to the encirclements of the critical point $-1 + j0$ of the $F(s)$ Nyquist plot

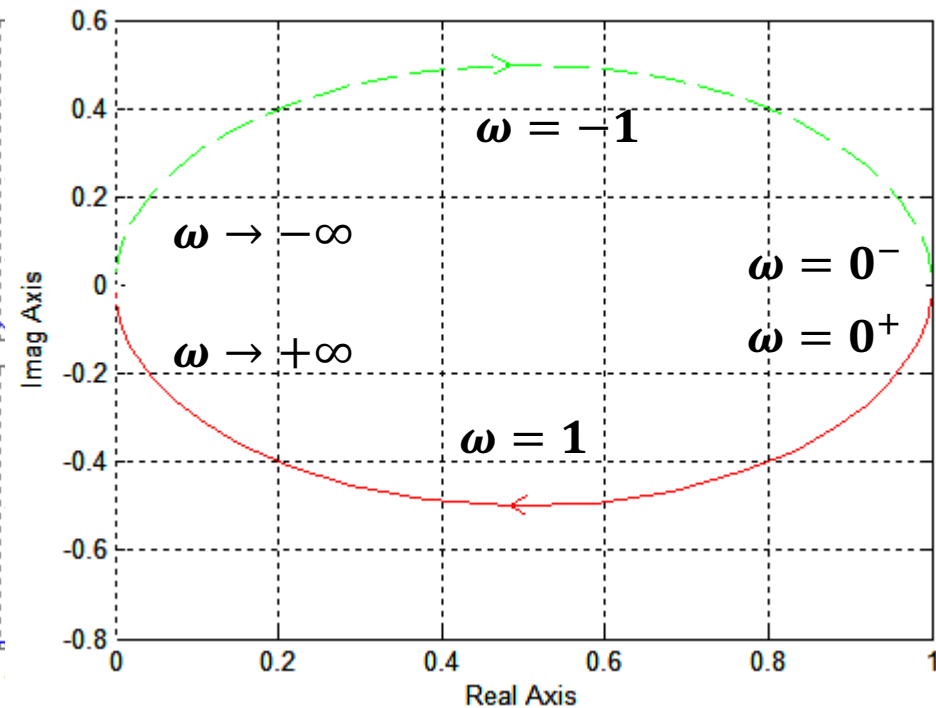
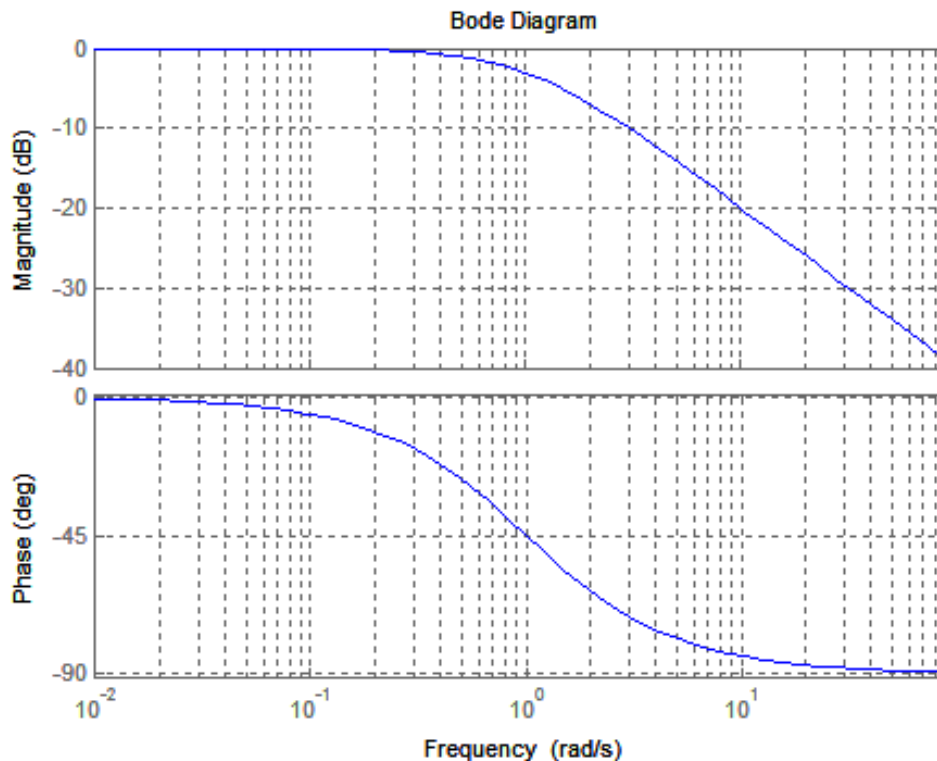




Nyquist stability criterion: example 1

Let us consider again the frequency response

$$F(s) = \frac{1}{1+s}$$

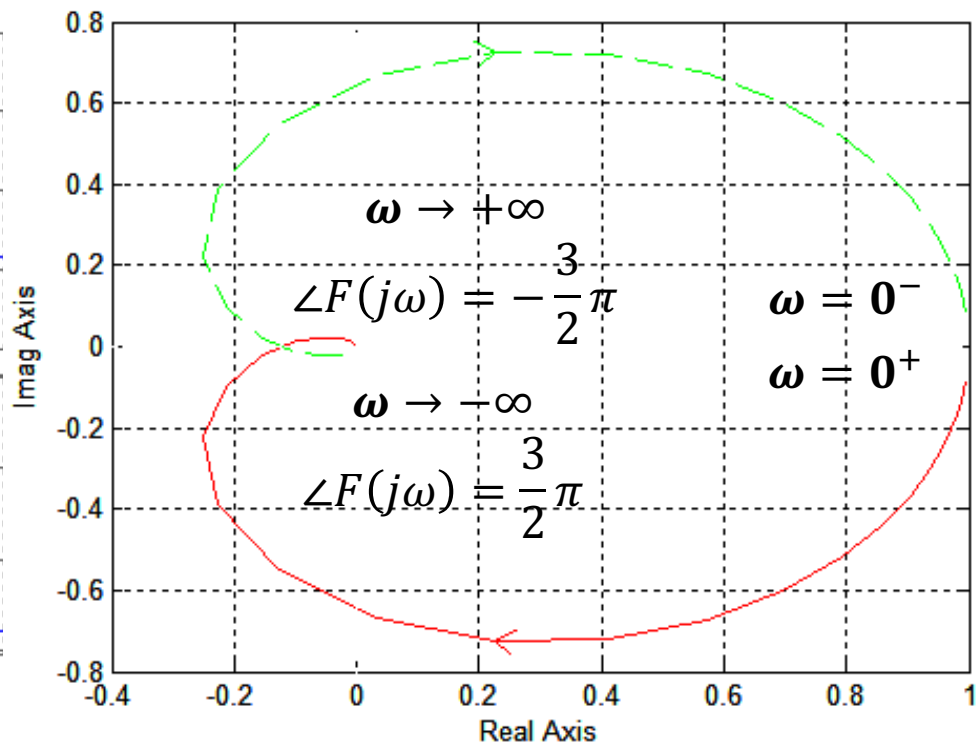
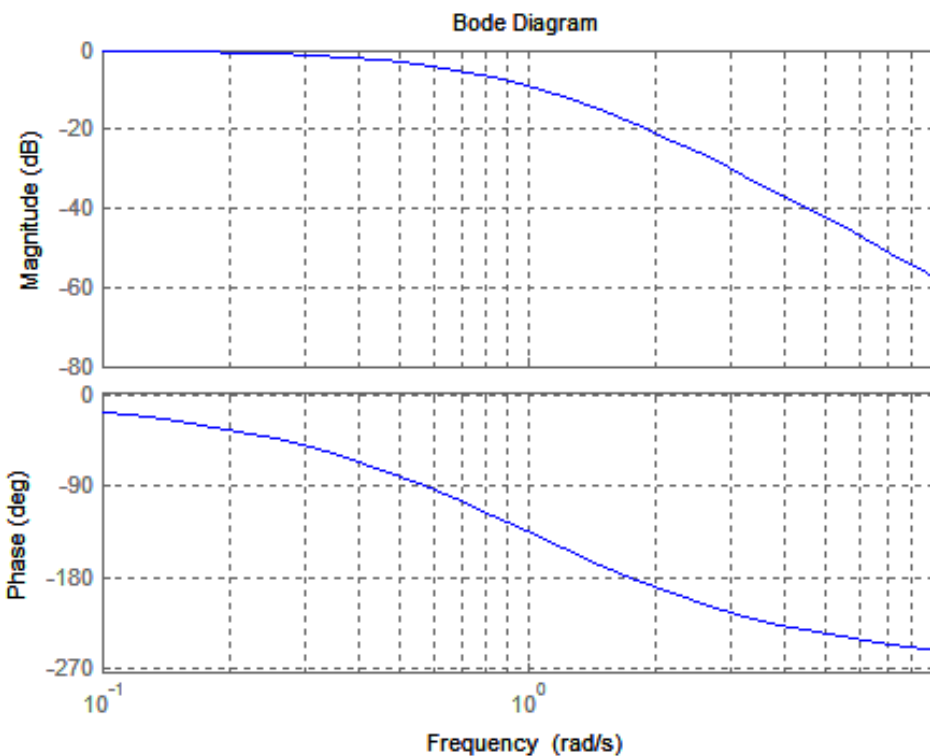


$\overline{N} = n_{p+}(F(s)) = 0 \rightarrow$ asymptotically stable closed loop function

Nyquist stability criterion: example 2

Let us consider the frequency response

$$F(s) = \frac{1}{(1+s)^3}$$



$\overline{N} = n_{p+}(F(s)) = 0 \rightarrow$ asymptotically stable closed loop function

However the two examples have an important difference in terms of robust stability