

# **Campi Elettromagnetici**

**Corso di Laurea in Ingegneria Informatica,  
Biomedica e delle Telecomunicazioni**

**a.a. 2019-2020 - Laurea “Triennale” – Secondo semestre - Secondo anno**

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# Color legend

New formulas, important considerations,  
important formulas, important concepts

Very important for the discussion

Memo

Mathematical tools to be exploited

Mathematics

# Mathematical tools that we will exploit today

$$\nabla \cdot (\nabla \times \vec{\mathbf{A}}(\vec{\mathbf{r}})) = 0$$

# Maxwell equations

Time domain - Differential form

$$\begin{cases} \nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t) = -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \\ \nabla \times \vec{\mathbf{h}}(\vec{\mathbf{r}}, t) = \frac{\partial \vec{\mathbf{d}}(\vec{\mathbf{r}}, t)}{\partial t} + \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{d}}(\vec{\mathbf{r}}, t) = \rho(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t) = 0 \end{cases}$$



**James Clerk Maxwell 1831-1879**

# Maxwell equations

Time domain - Differential form

$$\begin{cases} \nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t) = -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \\ \nabla \times \vec{\mathbf{h}}(\vec{\mathbf{r}}, t) = \frac{\partial \vec{\mathbf{d}}(\vec{\mathbf{r}}, t)}{\partial t} + \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) + \vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{d}}(\vec{\mathbf{r}}, t) = \rho(\vec{\mathbf{r}}, t) + \rho_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t) = 0 \end{cases}$$



**James Clerk Maxwell 1831-1879**

# Maxwell equations

$$\begin{cases} \nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t) = -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \\ \nabla \times \vec{\mathbf{h}}(\vec{\mathbf{r}}, t) = \frac{\partial \vec{\mathbf{d}}(\vec{\mathbf{r}}, t)}{\partial t} + \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) + \vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{d}}(\vec{\mathbf{r}}, t) = \rho(\vec{\mathbf{r}}, t) + \rho_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t) = 0 \end{cases}$$

$$\begin{cases} \vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t) \\ \rho_0(\vec{\mathbf{r}}, t) \end{cases} \quad \text{Prescribed sources}$$

$$\begin{cases} \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) \\ \rho(\vec{\mathbf{r}}, t) \end{cases} \quad \text{Induced sources}$$

## Complex scenario



# The independence of the Maxwell equations

Time domain - Differential form

$$\left\{ \begin{array}{l} \nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t) = -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \\ \nabla \times \vec{\mathbf{h}}(\vec{\mathbf{r}}, t) = \frac{\partial \vec{\mathbf{d}}(\vec{\mathbf{r}}, t)}{\partial t} + \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) + \vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{d}}(\vec{\mathbf{r}}, t) = \rho(\vec{\mathbf{r}}, t) + \rho_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t) = 0 \end{array} \right.$$

$$\nabla \cdot (\nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t)) = \nabla \cdot \left( -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \right)$$

↓    ↓

$$0 = -\frac{\partial (\nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t))}{\partial t}$$

$$\nabla \cdot (\nabla \times \vec{\mathbf{A}}(\vec{\mathbf{r}})) = 0$$

# The independence of the Maxwell equations

Time domain - Differential form

$$\left\{ \begin{array}{l} \nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t) = -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \\ \nabla \times \vec{\mathbf{h}}(\vec{\mathbf{r}}, t) = \frac{\partial \vec{\mathbf{d}}(\vec{\mathbf{r}}, t)}{\partial t} + \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) + \vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{d}}(\vec{\mathbf{r}}, t) = \rho(\vec{\mathbf{r}}, t) + \rho_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t) = 0 \end{array} \right.$$

$$\nabla \cdot (\nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t)) = \nabla \cdot \left( -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \right)$$

↓    ↓

$$0 = -\frac{\partial (\nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t))}{\partial t}$$

$\nabla \cdot \vec{\mathbf{b}}$  is independent of time. If the fields are equal to zero before a given time, then  $\nabla \cdot \vec{\mathbf{b}} = 0$  for all times, thus recovering the last Maxwell equation

# The independence of the Maxwell equations

Time domain - Differential form

$$\left\{ \begin{array}{l} \nabla \times \vec{e}(\vec{r}, t) = -\frac{\partial \vec{b}(\vec{r}, t)}{\partial t} \\ \nabla \times \vec{h}(\vec{r}, t) = \frac{\partial \vec{d}(\vec{r}, t)}{\partial t} + \vec{j}(\vec{r}, t) + \vec{j}_0(\vec{r}, t) \\ \nabla \cdot \vec{d}(\vec{r}, t) = \rho(\vec{r}, t) + \rho_0(\vec{r}, t) \\ \nabla \cdot \vec{b}(\vec{r}, t) = 0 \end{array} \right.$$

Number of independent scalar equations:  
 $3+3+1=7$

Let us assume knowledge of the impressed sources  $\vec{j}_0(\vec{r}, t); \rho_0(\vec{r}, t)$

Number of unknown scalar quantities:  
 $\vec{e}(\vec{r}, t); \vec{d}(\vec{r}, t); \vec{h}(\vec{r}, t); \vec{b}(\vec{r}, t); \vec{j}(\vec{r}, t); \rho(\vec{r}, t)$

3	+3	+3	+3	+3	+1
				16	

# The independence of the Maxwell equations

Time domain - Differential form

$$\begin{cases} \nabla \times \vec{e}(\vec{r}, t) = -\frac{\partial \vec{b}(\vec{r}, t)}{\partial t} \\ \nabla \times \vec{h}(\vec{r}, t) = \frac{\partial \vec{d}(\vec{r}, t)}{\partial t} + \vec{j}(\vec{r}, t) + \vec{j}_0(\vec{r}, t) \\ \nabla \cdot \vec{d}(\vec{r}, t) = \rho(\vec{r}, t) + \rho_0(\vec{r}, t) \\ \nabla \cdot \vec{b}(\vec{r}, t) = 0 \end{cases}$$

Even assuming knowledge of the impressed sources  $\vec{j}_0(\vec{r}, t); \rho_0(\vec{r}, t)$

Number of independent scalar equations: 7

Number of unknown scalar quantities: 16

Maxwell equations involve a number of unknowns larger than the number of equations!



# Constitutive relationships

## Time domain - Differential form

$$\begin{cases} \nabla \times \vec{\mathbf{e}}(\vec{\mathbf{r}}, t) = -\frac{\partial \vec{\mathbf{b}}(\vec{\mathbf{r}}, t)}{\partial t} \\ \nabla \times \vec{\mathbf{h}}(\vec{\mathbf{r}}, t) = \frac{\partial \vec{\mathbf{d}}(\vec{\mathbf{r}}, t)}{\partial t} + \vec{\mathbf{j}}(\vec{\mathbf{r}}, t) + \vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{d}}(\vec{\mathbf{r}}, t) = \rho(\vec{\mathbf{r}}, t) + \rho_0(\vec{\mathbf{r}}, t) \\ \nabla \cdot \vec{\mathbf{b}}(\vec{\mathbf{r}}, t) = 0 \end{cases}$$

Even assuming knowledge of the impressed sources  $\vec{\mathbf{j}}_0(\vec{\mathbf{r}}, t); \rho_0(\vec{\mathbf{r}}, t)$

Number of independent scalar equations: 7

Number of unknown scalar quantities: 16

Maxwell equations involve a number of unknowns larger than the number of equations!

The additional missing equations are provided by the **constitutive relationships**, which describe interaction of fields and matter from a macroscopic point of view

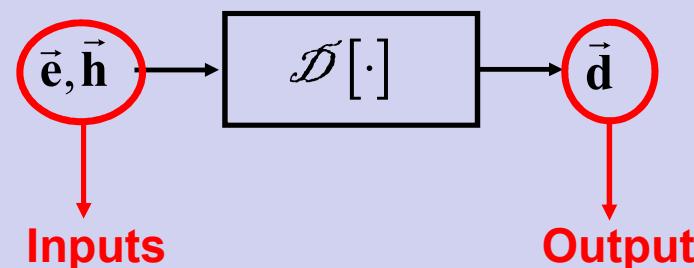
# Constitutive relationships

Inductions and currents must be represented in terms of fields

$$\vec{d} = \mathcal{D}[\vec{e}, \vec{h}]$$

$$\vec{b} = \mathcal{B}[\vec{e}, \vec{h}]$$

$$\vec{j} = \mathcal{J}[\vec{e}, \vec{h}]$$



$\mathcal{D}[\cdot]$ ,  $\mathcal{B}[\cdot]$  and  $\mathcal{J}[\cdot]$  are functionals which depend upon the medium in which the electromagnetic field is considered and upon the fields themselves.

# Constitutive relationships

## Linear media

$$\vec{d}_1 = \mathcal{D}[\vec{e}_1, \vec{h}_1] \quad \rightarrow \quad \vec{d}_1 + \vec{d}_2 = \mathcal{D}[\vec{e}_1 + \vec{e}_2, \vec{h}_1 + \vec{h}_2]$$
$$\vec{d}_2 = \mathcal{D}[\vec{e}_2, \vec{h}_2]$$

# Constitutive relationships

In the following we will consider linear media

# Constitutive relationships

## Linear media

### Example 1

$$\vec{d}(\vec{r}, t) = \epsilon(\vec{r}, t) \cdot \vec{e}(\vec{r}, t) + \chi(\vec{r}, t) \cdot \vec{h}(\vec{r}, t)$$

where

$\vec{d}(\vec{r}, t)$ ;  $\vec{e}(\vec{r}, t)$ ;  $\vec{h}(\vec{r}, t)$ : 3x1 column vectors

$\epsilon(\vec{r}, t)$ ;  $\chi(\vec{r}, t)$ : 3x3 matrices

■ Local (non-dispersive) media

■ Bianisotropic media

# Constitutive relationships

## Linear media

### Example 2

$$\vec{d}(\vec{r}, t) = \epsilon(\vec{r}, t) \cdot \vec{e}(\vec{r}, t)$$

where

$\vec{d}(\vec{r}, t)$ ;  $\vec{e}(\vec{r}, t)$ ;  $\vec{h}(\vec{r}, t)$ : 3x1 column vectors

$\epsilon(\vec{r}, t)$ : 3x3 matrices

■ Local (non-dispersive) media

■ Anisotropic media

# Constitutive relationships

## Linear media

### Example 2

$$\vec{d}(\vec{r}, t) = \epsilon(\vec{r}, t) \cdot \vec{e}(\vec{r}, t)$$

$$\vec{b}(\vec{r}, t) = \mu(\vec{r}, t) \cdot \vec{h}(\vec{r}, t)$$

$$\vec{j}(\vec{r}, t) = \sigma(\vec{r}, t) \cdot \vec{e}(\vec{r}, t)$$

where

$\vec{d}(\vec{r}, t); \vec{b}(\vec{r}, t); \vec{j}(\vec{r}, t); \vec{e}(\vec{r}, t); \vec{h}(\vec{r}, t)$ : 3x1 column vectors

$\epsilon(\vec{r}, t); \mu(\vec{r}, t); \sigma(\vec{r}, t)$ : 3x3 matrices

■ Local (non-dispersive) media

■ Anisotropic media

$\epsilon(\vec{r}, t)$ : permittivity [Farad/m]

$\mu(\vec{r}, t)$ : permeability [Henry/m]

$\sigma(\vec{r}, t)$ : conductivity [Siemens/m]

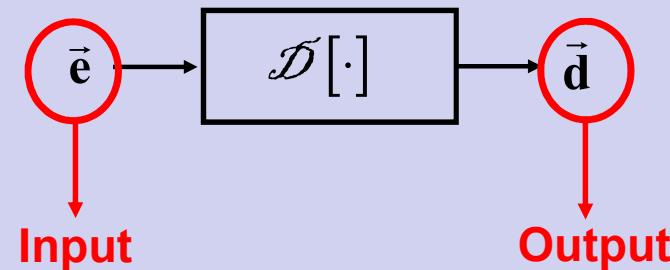
# Constitutive relationships

## Linear & Anisotropic media

$$\vec{d} = \mathcal{D}[\vec{e}]$$

$$\vec{b} = \mathcal{B}[\vec{h}]$$

$$\vec{j} = \mathcal{J}[\vec{e}]$$



# Constitutive relationships

## Linear media

$$\vec{d}(\vec{r}, t) = \epsilon(\vec{r}, t) \cdot \vec{e}(\vec{r}, t)$$

$\epsilon(\vec{r}, t)$ : 3x3 matrix

- Local (non-dispersive) media
- Anisotropic media

### Class

Isotropic

### Property

A **rotation** of the input implies  
the **same rotation** of the  
output

### Effect on the I-O relation

$\epsilon(\vec{r}, t)$  becomes scalar. It is  
not a matrix anymore!

# Constitutive relationships

## Linear media

### ■ Local (non-dispersive) media

### ■ Anisotropic media

$$\vec{d}(\vec{r}, t) = \boldsymbol{\varepsilon}(\vec{r}, t) \cdot \vec{e}(\vec{r}, t)$$

$$\vec{b}(\vec{r}, t) = \boldsymbol{\mu}(\vec{r}, t) \cdot \vec{h}(\vec{r}, t)$$

$$\vec{j}(\vec{r}, t) = \boldsymbol{\sigma}(\vec{r}, t) \cdot \vec{e}(\vec{r}, t)$$

$\boldsymbol{\varepsilon}(\vec{r}, t); \boldsymbol{\mu}(\vec{r}, t); \boldsymbol{\sigma}(\vec{r}, t)$ : 3x3 matrices

### ■ Local (non-dispersive) media

### ■ Isotropic media

$$\vec{d}(\vec{r}, t) = \varepsilon(\vec{r}, t) \vec{e}(\vec{r}, t)$$

$$\vec{b}(\vec{r}, t) = \mu(\vec{r}, t) \vec{h}(\vec{r}, t)$$

$$\vec{j}(\vec{r}, t) = \sigma(\vec{r}, t) \vec{e}(\vec{r}, t)$$

$\varepsilon(\vec{r}, t); \mu(\vec{r}, t); \sigma(\vec{r}, t)$ : scalar functions

# Constitutive relationships

## Linear media

### Class

Time-dispersive

Time-nondispersive

Time-variant

Time-invariant

### Property

### Effect on the I-O relation

# Constitutive relationships

## Linear media

### Class

**Space-dispersive**

**Space-nondispersive**

**Space-variant**

**Space-invariant**

### Property

### Effect on the I-O relation

# Constitutive relationships

## Linear media

### Class

Time-dispersive

Time-nondispersive

Time-variant

Time-invariant

### Property

The output **at time  $t$  depends only** on the value of the input **at the same time  $t$**

### Effect on the I-O relation

# Constitutive relationships

## Linear media

### Class

Time-dispersive

Time-nondispersive

Time-variant

Time-invariant

### Property

The output at time  $t$  depends on the values of the input throughout a time-interval.

### Effect on the I-O relation

# Constitutive relationships

## Linear media

### Class

Time-dispersive

Time-nondispersive

Time-variant

Time-invariant

### Property

A time translation of the input implies the same translation of the output

### Effect on the I-O relation

# Constitutive relationships

## Linear media

### Class

Time-dispersive

Time-nondispersive

Time-variant

Time-invariant

### Property

A time translation of the input does not imply the same translation of the output

### Effect on the I-O relation

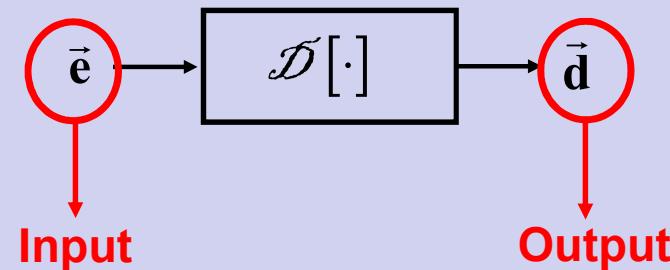
# Constitutive relationships

## Linear & Anisotropic media

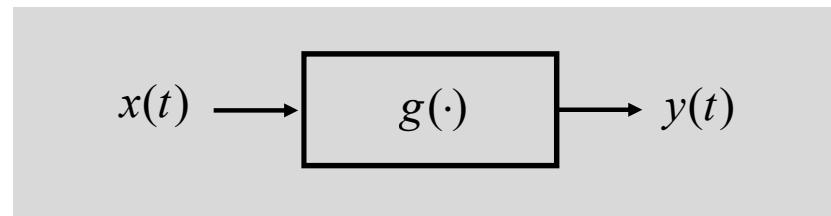
$$\vec{d} = \mathcal{D}[\vec{e}]$$

$$\vec{b} = \mathcal{B}[\vec{h}]$$

$$\vec{j} = \mathcal{J}[\vec{e}]$$



# Memo: time-dispersive (TD) linear systems



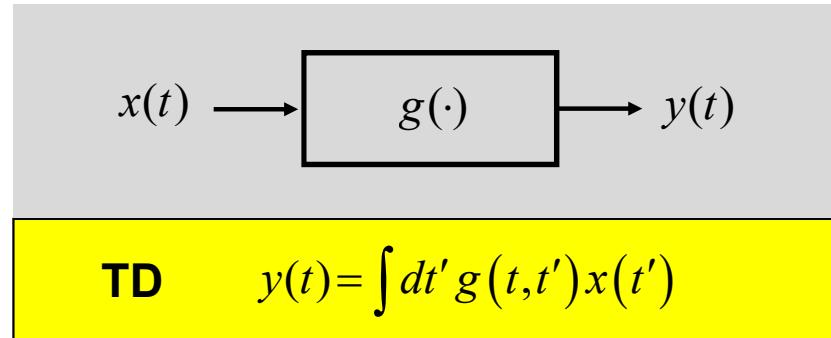
## Effect on the I-O relation

$$y(t) = \int dt' g(t, t') x(t')$$

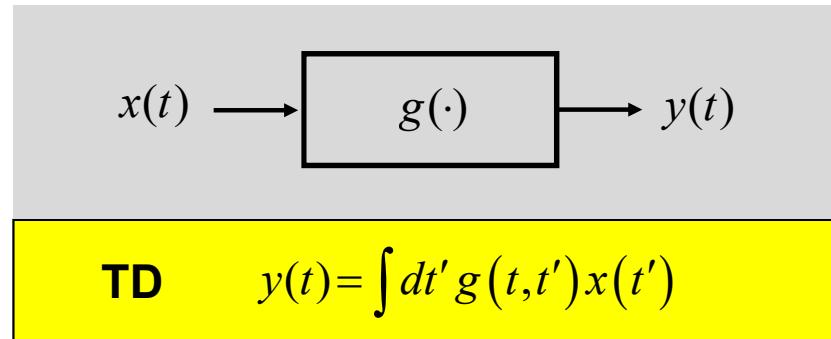
The output **at time  $t$  depends** on the values of the input **throughout a time-interval**.

In the most general case, these systems possess an heredity: they are called **dispersive**

# Memo: Time-nondispersive (TND) linear systems



# Memo: Time-nondispersive (TND) linear systems



## Property

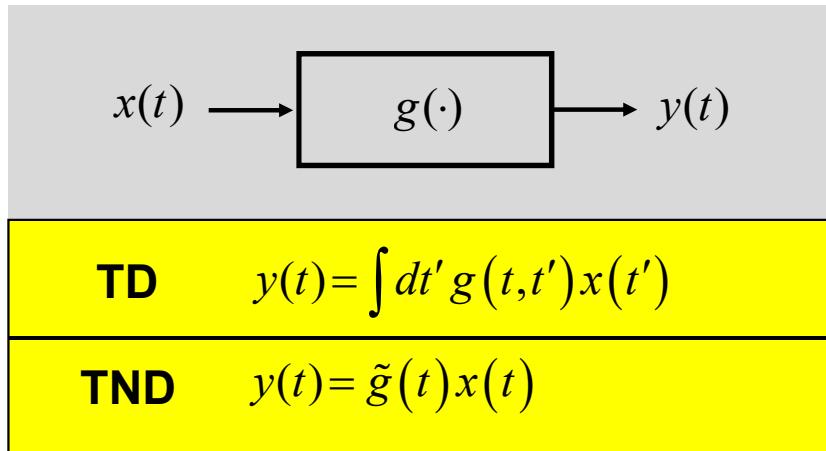
The output **at time  $t$  depends only** on the value of the input **at the same time  $t$**

## Effect on the I-O relation

$$g(t, t') \rightarrow \delta(t - t') \tilde{g}(t')$$

$$y(t) = \int dt' \delta(t - t') \tilde{g}(t') x(t') = \tilde{g}(t) x(t)$$

# Memo: Time-nondispersive (TND) linear systems



## Property

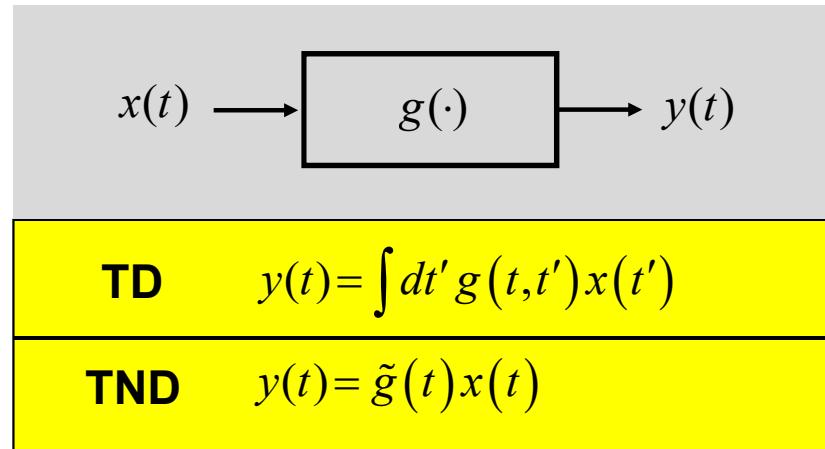
The output **at time  $t$  depends only** on the value of the input **at the same time  $t$**

## Effect on the I-O relation

$$g(t, t') \rightarrow \delta(t - t') \tilde{g}(t')$$

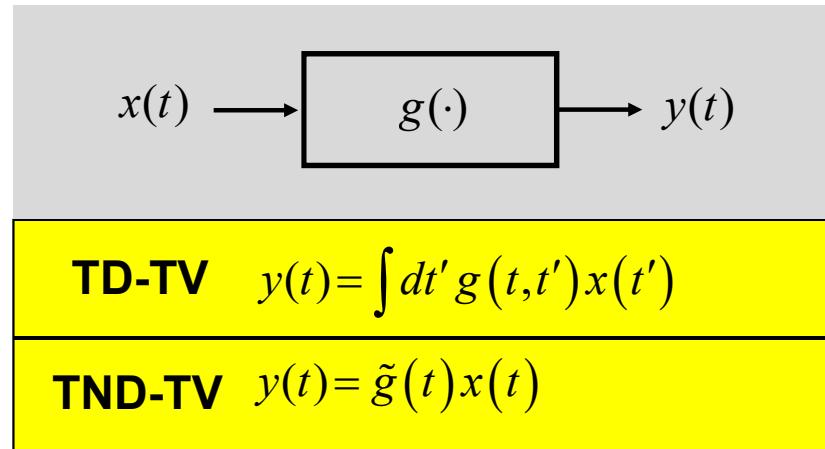
$$y(t) = \int dt' \delta(t - t') \tilde{g}(t') x(t') = \tilde{g}(t) x(t)$$

# Memo: time-invariant (TI) linear systems



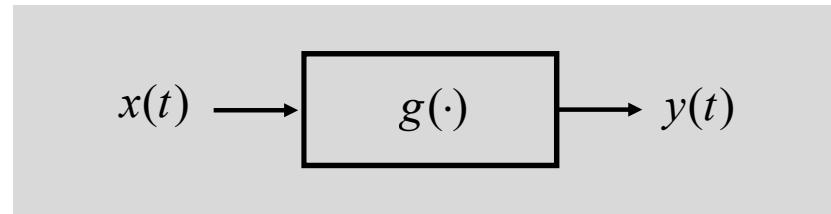
Property	Effect on the I-O relation TD case	Effect on the I-O relation TND case
$x(t) \rightarrow g(\cdot) \rightarrow y(t)$		
$x(t-T) \rightarrow g(\cdot) \rightarrow y(t-T)$		

# Memo: time-invariant (TI) linear systems



Property	Effect on the I-O relation TD case	Effect on the I-O relation TND case
$x(t) \rightarrow g(\cdot) \rightarrow y(t)$	$g(t, t') \rightarrow g(t - t')$ $y(t) = \int dt' g(t - t') x(t')$	$\tilde{g}(t) \rightarrow \tilde{g}$ $y(t) = \tilde{g} x(t)$

# Memo: linear systems



$x(t)$  and  $y(t)$  are scalar

	Time-dispersive (TD)	Time-nondispersive (TND)
Time-variant (TV)	$y(t) = \int dt' g(t, t') x(t')$	$y(t) = \tilde{g}(t) x(t)$
Time-invariant (TI)	$y(t) = \int dt' g(t - t') x(t')$	$y(t) = \tilde{g} x(t)$