

# **Campi Elettromagnetici**

**Corso di Laurea in Ingegneria Informatica,  
Biomedica e delle Telecomunicazioni**

**a.a. 2020-2021 - Laurea "Triennale" – Secondo semestre - Secondo anno**

**Università degli Studi di Napoli "Parthenope"**

**Stefano Perna**

# Color legend

New formulas, important considerations,  
important formulas, important concepts

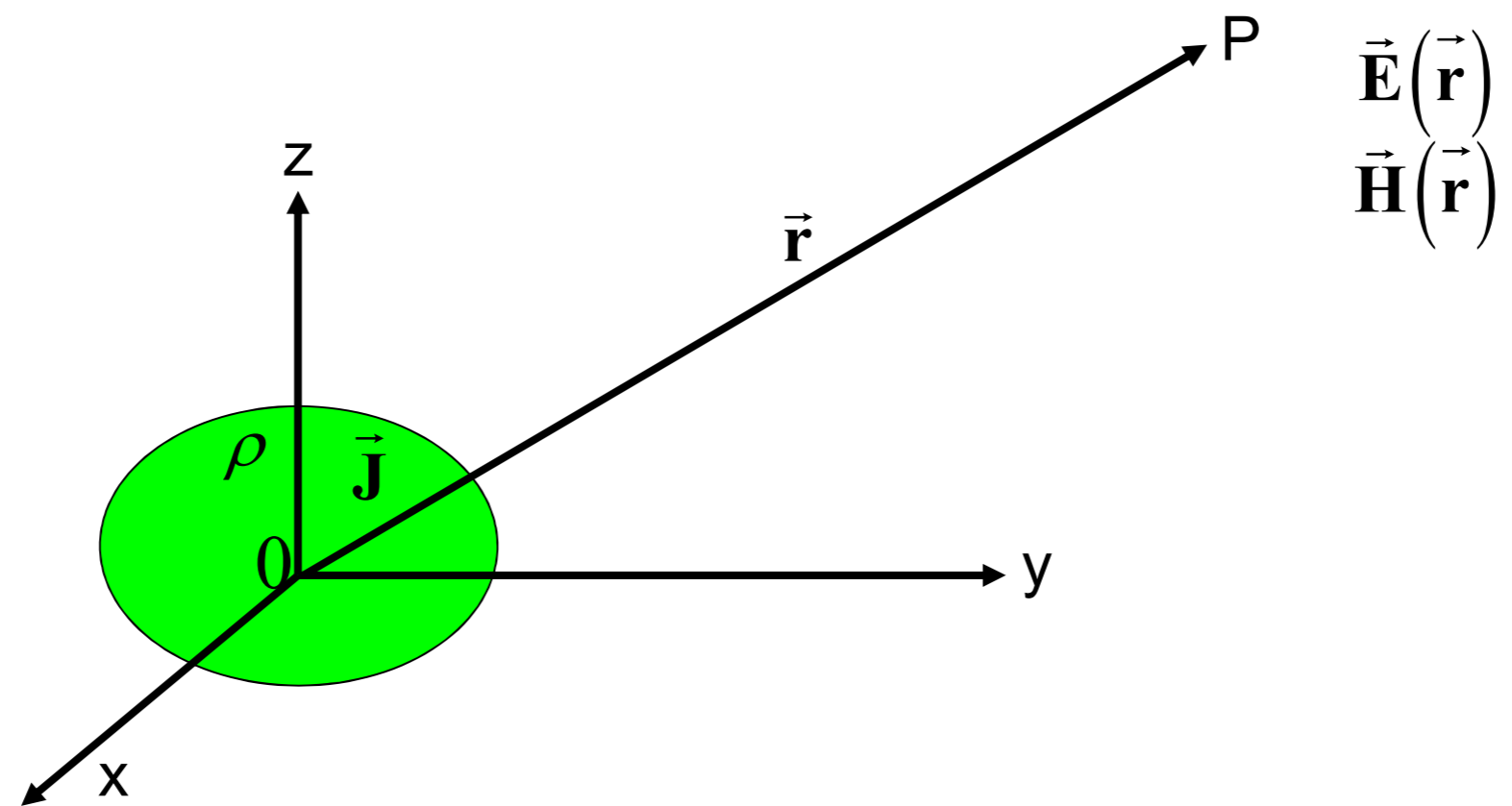
Very important for the discussion

Memo

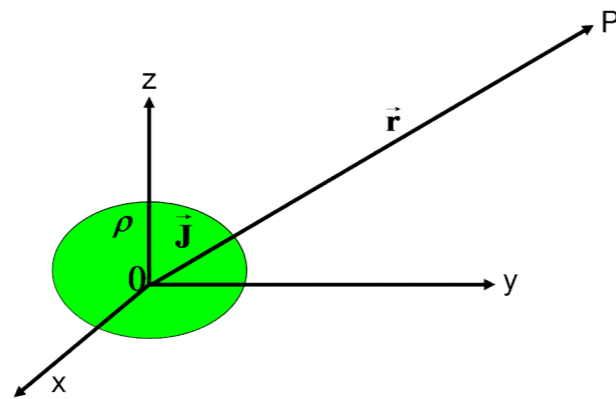
Mathematical tools to be exploited

Mathematics

# Radiation problem



An antenna (or aerial) is an electrical device which converts electric power into radio waves, and vice versa.



# Potentials

↓  $\mathbf{J}(\mathbf{r})$

$$\nabla^2 \mathbf{A} + k^2 \mathbf{A} = -\mu \mathbf{J} \quad \text{Vector Helmholtz equation}$$

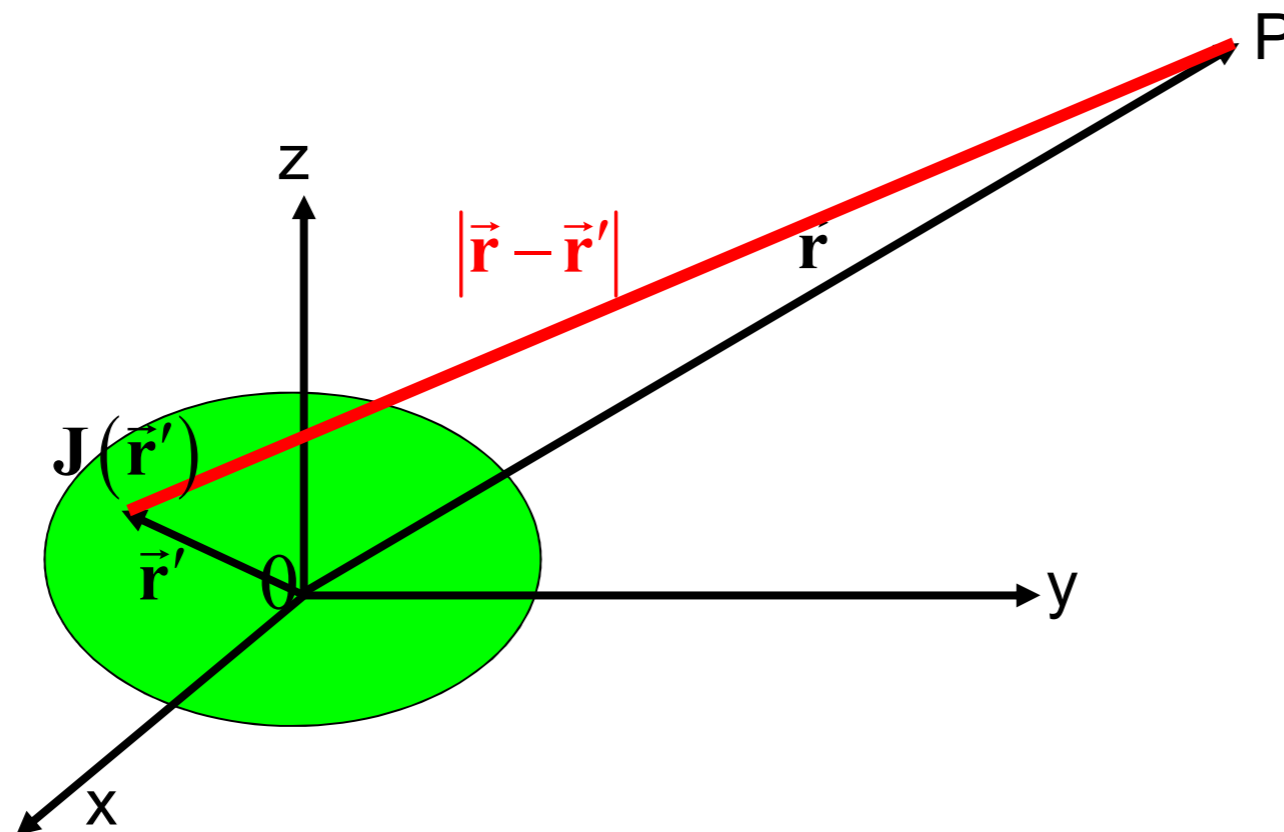
$$\downarrow \mathbf{A}(\mathbf{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\mathbf{r}') \frac{e^{-jk|\mathbf{r}-\mathbf{r}'|}}{|\mathbf{r}-\mathbf{r}'|} d\mathbf{r}'$$

$$\mathbf{H} = \frac{1}{\mu} \nabla \times \mathbf{A} \quad \text{and} \quad \mathbf{E} = -j\omega \mathbf{A} + \frac{\nabla \nabla \cdot \mathbf{A}}{j\omega \epsilon \mu}$$

↓  $\mathbf{E}(\mathbf{r})$   
 $\mathbf{H}(\mathbf{r})$

# Potentials

$$\mathbf{A}(\vec{\mathbf{r}}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{\mathbf{r}}') \frac{e^{-jk|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|}}{|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|} d\vec{\mathbf{r}}'$$

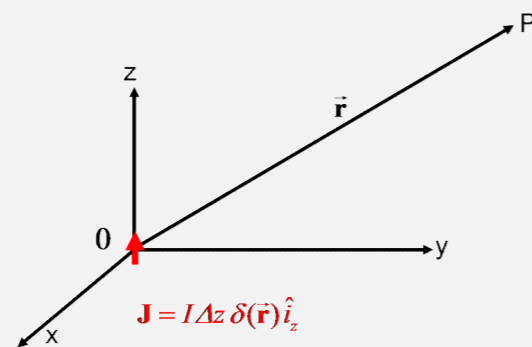


# Elementary electrical dipole vs. small loop antenna

## Elementary electrical dipole

$$\mathbf{J} = I \Delta z \delta(\vec{r}) \hat{i}_z$$

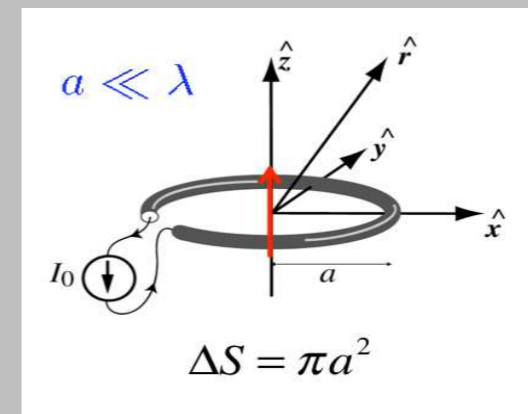
$$\begin{cases} E_r = \zeta \frac{I \Delta z}{2\pi} \left( \frac{1}{r^2} + \frac{1}{j\beta r^3} \right) \cos \vartheta \exp(-j\beta r) \\ E_\vartheta = \zeta \frac{I \Delta z}{4\pi} \left( \frac{j\beta}{r} + \frac{1}{r^2} + \frac{1}{j\beta r^3} \right) \sin \vartheta \exp(-j\beta r) \\ H_\varphi = \frac{I \Delta z}{4\pi} \left( \frac{j\beta}{r} + \frac{1}{r^2} \right) \sin \vartheta \exp(-j\beta r) \end{cases}$$



## Small loop antenna

$$\mathbf{J} = I \delta(z) \delta(r - a) \hat{i}_\varphi$$

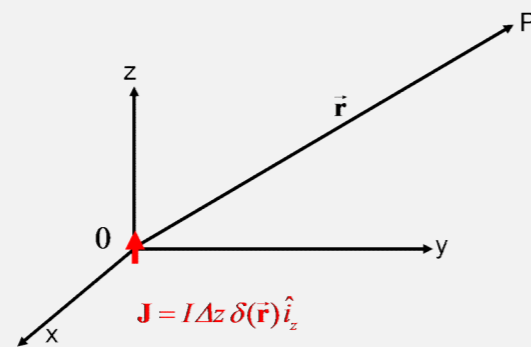
$$\begin{cases} H_r = \frac{I \Delta S}{2\pi} j\beta \left( \frac{1}{r^2} + \frac{1}{j\beta r^3} \right) \cos \vartheta \exp(-j\beta r) \\ H_\vartheta = \frac{I \Delta S}{4\pi} j\beta \left( \frac{j\beta}{r} + \frac{1}{r^2} + \frac{1}{j\beta r^3} \right) \sin \vartheta \exp(-j\beta r) \\ E_\varphi = -\frac{\zeta I \Delta S}{4\pi} j\beta \left( \frac{j\beta}{r} + \frac{1}{r^2} \right) \sin \vartheta \exp(-j\beta r) \end{cases}$$



# Elementary electrical and magnetic dipoles

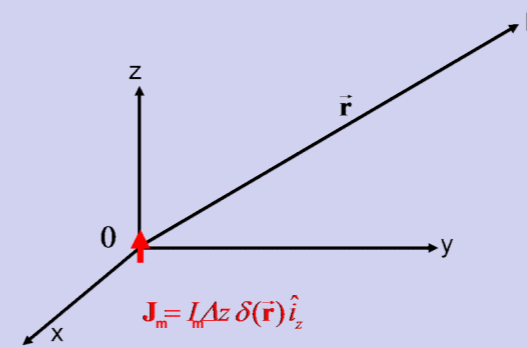
## Elementary electrical dipole

$$\mathbf{J} = I \Delta z \delta(\vec{\mathbf{r}}) \hat{i}_z$$



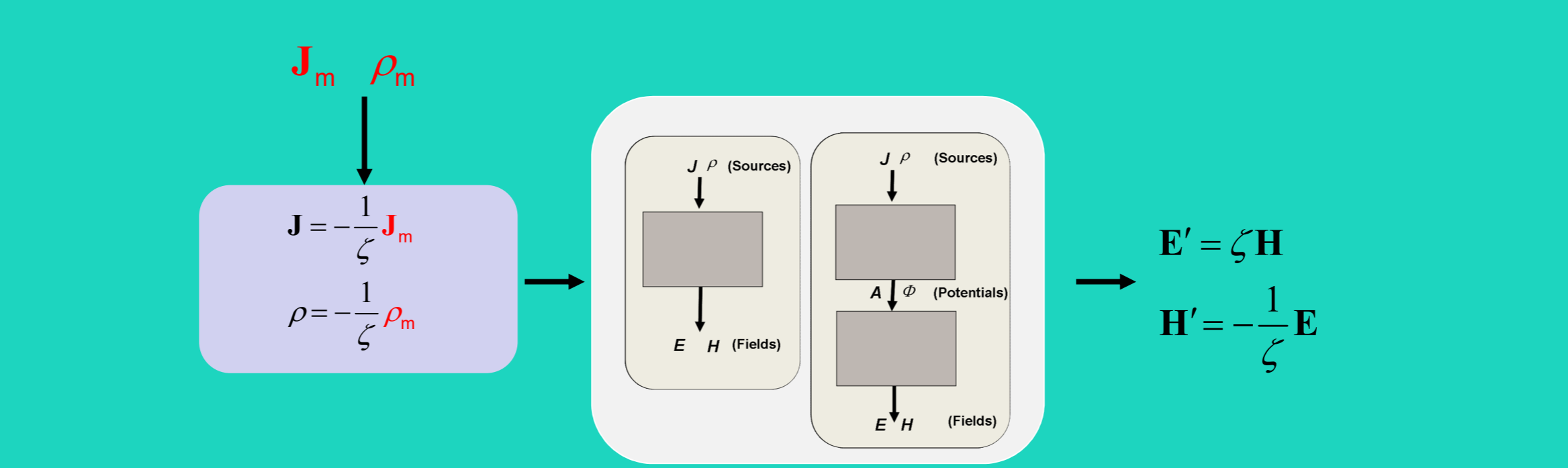
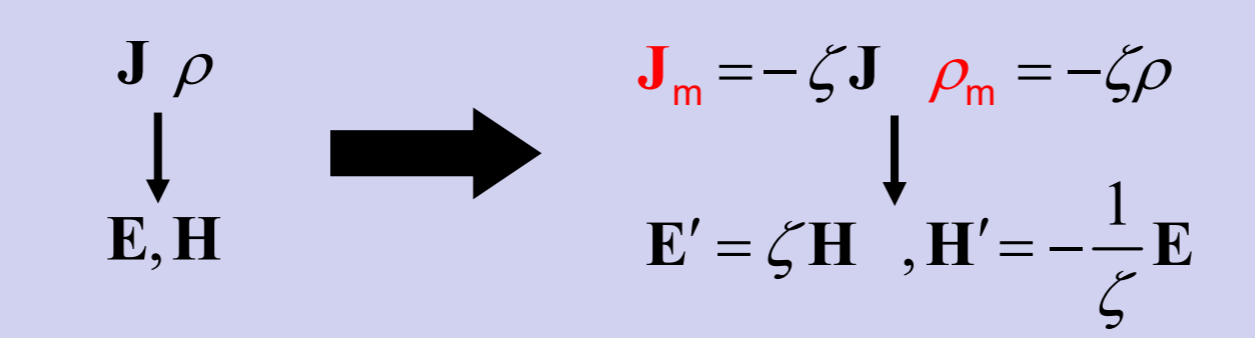
## Elementary magnetic dipole

$$\mathbf{J}_m = I_m \Delta z \delta(\vec{\mathbf{r}}) \hat{i}_z$$

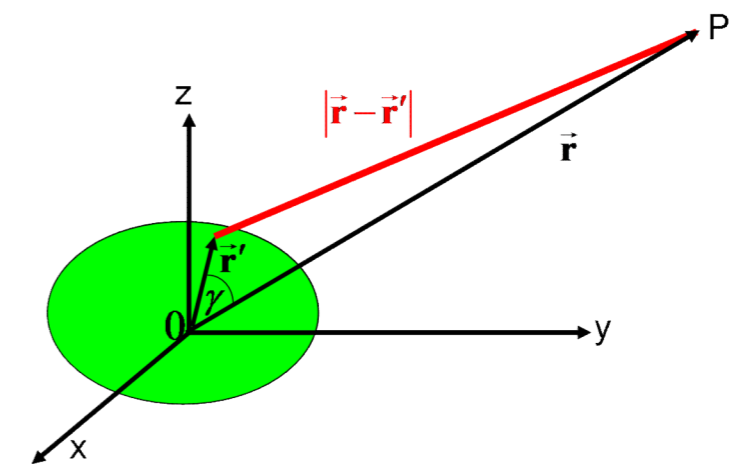
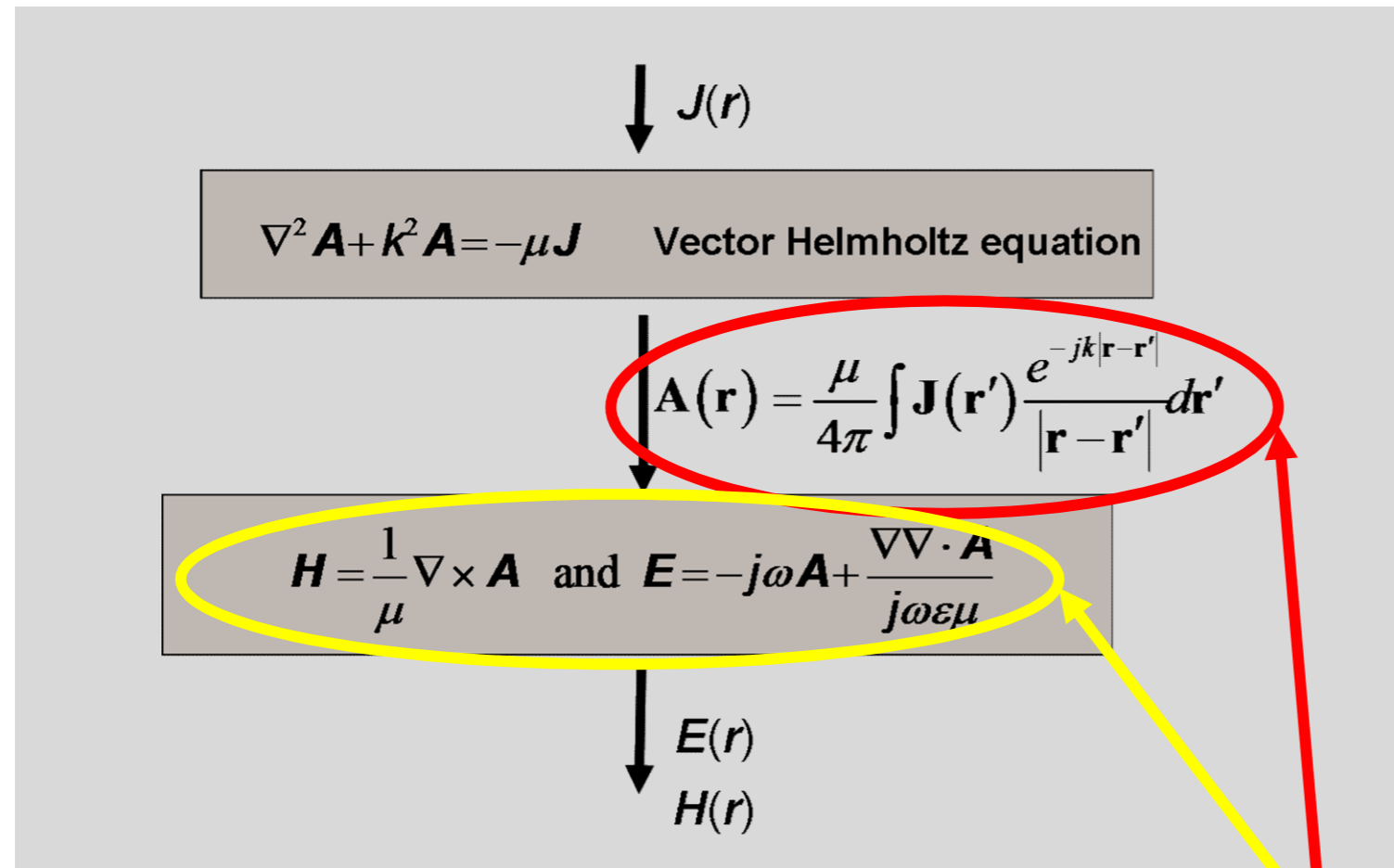




# Duality Theorem



# Extended antennas

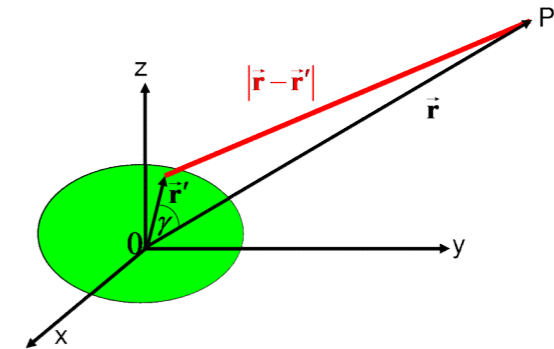


Is it possible to simplify the expressions of the fields, possibly via proper approximation of the vector potential  $\mathbf{A}$ ?

# Extended antennas

$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

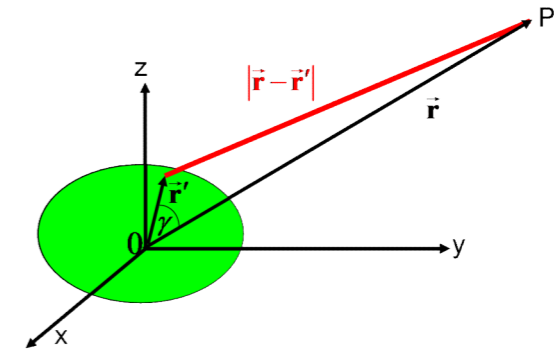
$$|\vec{r}-\vec{r}'| = r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$



# Extended antennas

$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

$$|\vec{r}-\vec{r}'| = r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$



$$\frac{1}{|\vec{r}-\vec{r}'|} = \frac{1}{r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots} \approx \frac{1}{r} \quad \text{if } r' \ll r$$

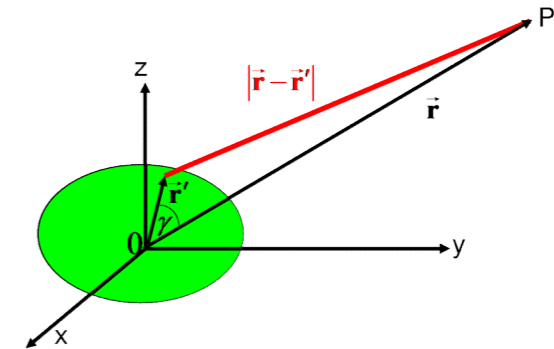
$$e^{-j\beta|\vec{r}-\vec{r}'|} = e^{-j\beta r} e^{j\beta r' \cos \gamma} e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \dots \approx e^{-j\beta r} \quad \text{if } r' \ll \lambda$$

$$e^{j\beta r' \cos \gamma} \approx 1 \quad \rightarrow \quad \frac{2\pi}{\lambda} r' \ll 2\pi \quad \rightarrow \quad r' \ll \lambda$$

# Extended antennas

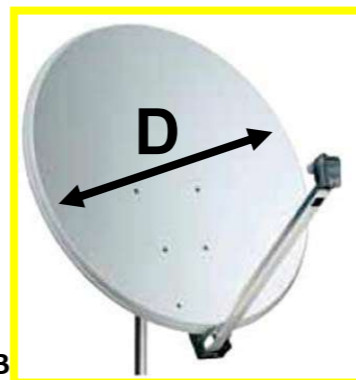
$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

$$|\vec{r}-\vec{r}'| = r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$



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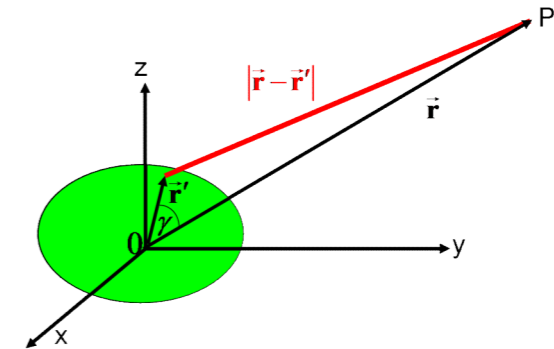
$$e^{-j\beta|\vec{r}-\vec{r}'|} = e^{-j\beta r} e^{j\beta r' \cos \gamma} e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \dots \approx e^{-j\beta r} \quad \text{if } D \ll \lambda$$



# Extended antennas

$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

$$|\vec{r}-\vec{r}'| = r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$



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$$e^{-j\beta|\vec{r}-\vec{r}'|} = e^{-j\beta r} e^{j\beta r' \cos \gamma} e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \dots \approx e^{-j\beta r} \quad \text{if } D \ll \lambda$$

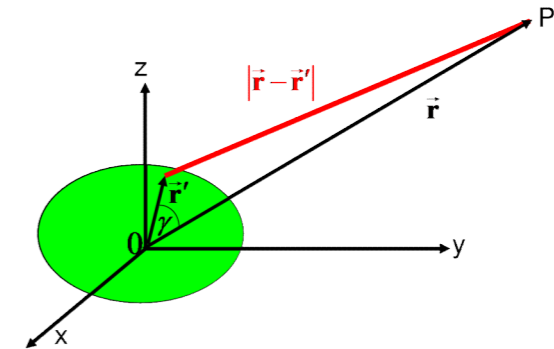
When the antennas are **small** with respect to the wavelength **and** to the distance from the observation point

$$\frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} \approx \frac{e^{-j\beta r}}{r} \quad \longrightarrow \quad \mathbf{A}(\vec{r}) \approx \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \int \mathbf{J}(\vec{r}') d\vec{r}'$$

# Extended antennas

$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

$$|\vec{r}-\vec{r}'| = r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$



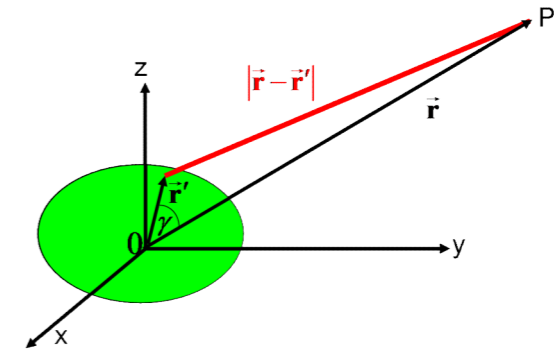
$$\frac{1}{|\vec{r}-\vec{r}'|} = \frac{1}{r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots} \approx \frac{1}{r} \quad \text{if } D \ll r$$

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# Extended antennas

$$\mathbf{A}(\vec{\mathbf{r}}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{\mathbf{r}}') \frac{e^{-j\beta|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|}}{|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|} d\vec{\mathbf{r}}'$$

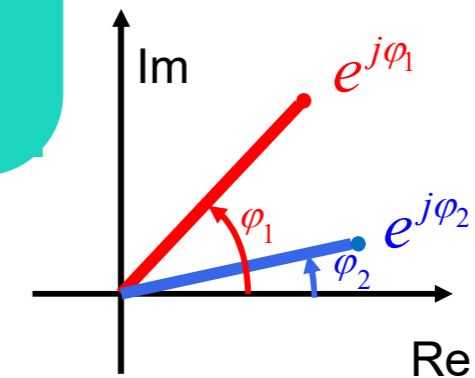
$$|\vec{\mathbf{r}}-\vec{\mathbf{r}}'| = r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$



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$$e^{-j\beta|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|} = e^{-j\beta r} e^{j\beta r' \cos \gamma} e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \dots \approx e^{-j\beta r} e^{j\beta r' \cos \gamma}$$

$$e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \approx 1 \quad \Rightarrow \quad \beta \frac{(r')^2}{2r} \sin^2 \gamma \ll 2\pi \quad \Rightarrow \quad \frac{2\pi}{\lambda} \left(\frac{D}{2}\right)^2 \frac{1}{2r} < \frac{\pi}{8} \quad \Rightarrow \quad r > \frac{2D^2}{\lambda}$$

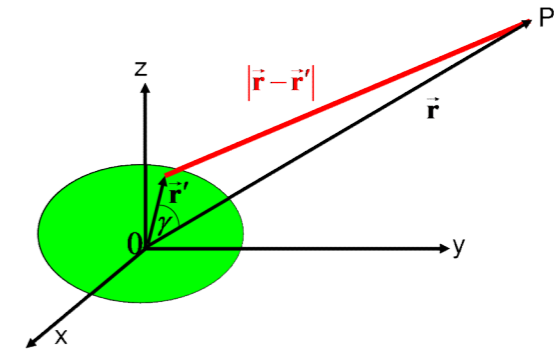




# Extended antennas

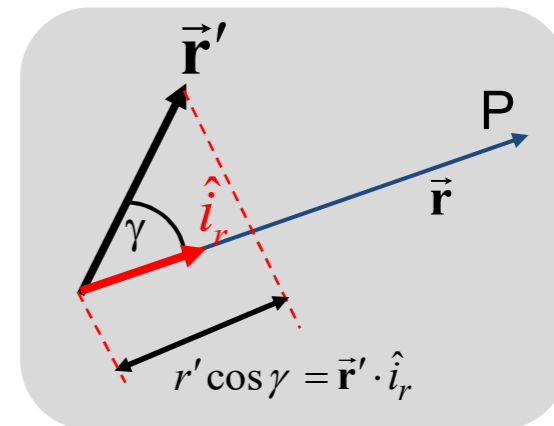
$$\mathbf{A}(\vec{\mathbf{r}}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{\mathbf{r}}') \frac{e^{-j\beta|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|}}{|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|} d\vec{\mathbf{r}}'$$

$$|\vec{\mathbf{r}}-\vec{\mathbf{r}}'| = r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$



$$\frac{1}{|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|} = \frac{1}{r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots} \approx \frac{1}{r} \quad \text{if } D \ll r$$

$$e^{-j\beta|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|} = e^{-j\beta r} e^{j\beta r' \cos \gamma} e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \dots \approx e^{-j\beta r} e^{j\beta r' \cos \gamma} \quad \text{if } r > \frac{2D^2}{\lambda}$$

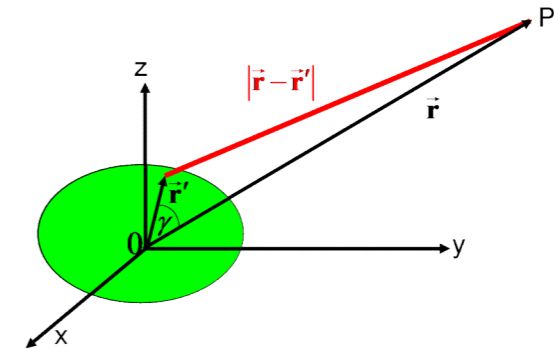


$$e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \approx 1 \quad \Rightarrow \quad \beta \frac{(r')^2}{2r} \sin^2 \gamma \ll 2\pi \quad \Rightarrow \quad \frac{2\pi}{\lambda} \left(\frac{D}{2}\right)^2 \frac{1}{2r} < \frac{\pi}{8} \quad \Rightarrow \quad r > \frac{2D^2}{\lambda}$$

# Extended antennas

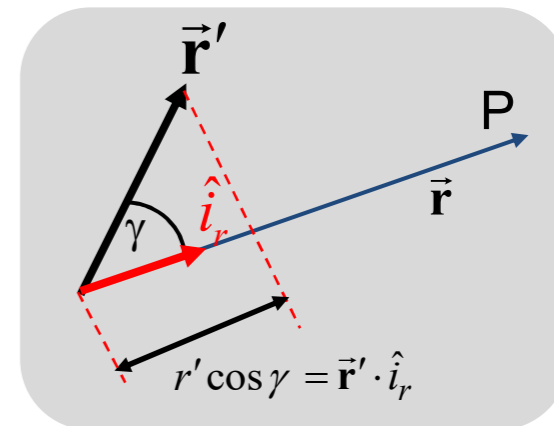
$$\mathbf{A}(\vec{\mathbf{r}}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{\mathbf{r}}') \frac{e^{-j\beta|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|}}{|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|} d\vec{\mathbf{r}}'$$

$$|\vec{\mathbf{r}}-\vec{\mathbf{r}}'| = r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$



$$\frac{1}{|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|} = \frac{1}{r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots} \approx \frac{1}{r} \quad \text{if } D \ll r$$

$$e^{-j\beta|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|} = e^{-j\beta r} e^{j\beta r' \cos \gamma} e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \dots \approx e^{-j\beta r} e^{j\beta \vec{\mathbf{r}} \cdot \hat{\mathbf{i}}_r} \quad \text{if } r > \frac{2D^2}{\lambda}$$

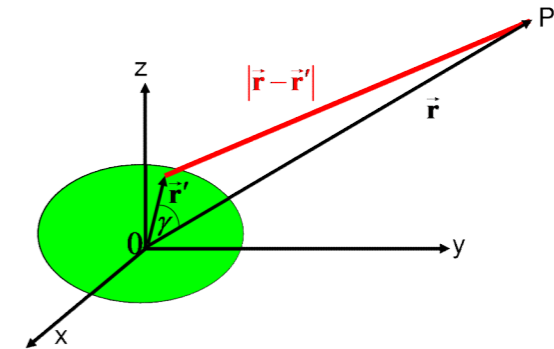


$$e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \approx 1 \quad \Rightarrow \quad \beta \frac{(r')^2}{2r} \sin^2 \gamma \ll 2\pi \quad \Rightarrow \quad \frac{2\pi}{\lambda} \left(\frac{D}{2}\right)^2 \frac{1}{2r} < \frac{\pi}{8} \quad \Rightarrow \quad r > \frac{2D^2}{\lambda}$$

# Extended antennas

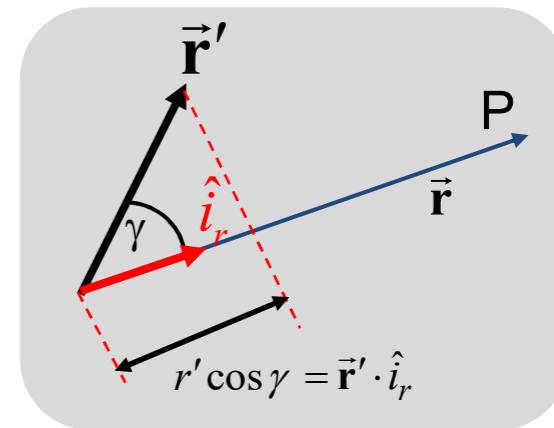
$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

$$|\vec{r}-\vec{r}'| = r - \vec{r}' \cdot \hat{i}_r + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$



$$\frac{1}{|\vec{r}-\vec{r}'|} = \frac{1}{r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots} \approx \frac{1}{r} \quad \text{if } D \ll r$$

$$e^{-j\beta|\vec{r}-\vec{r}'|} = e^{-j\beta r} e^{j\beta r' \cos \gamma} e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \dots \approx e^{-j\beta r} e^{j\beta \vec{r}' \cdot \hat{i}_r} \quad \text{if } r > \frac{2D^2}{\lambda}$$

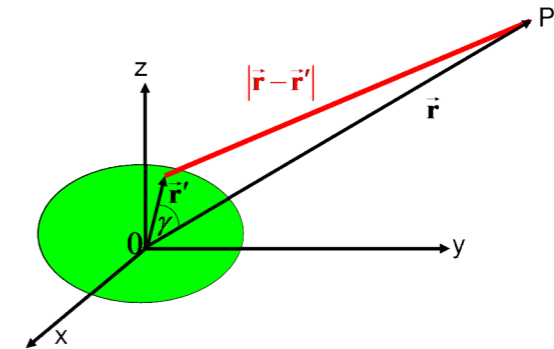


$$e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \approx 1 \quad \Rightarrow \quad \beta \frac{(r')^2}{2r} \sin^2 \gamma \ll 2\pi \quad \Rightarrow \quad \frac{2\pi}{\lambda} \left(\frac{D}{2}\right)^2 \frac{1}{2r} < \frac{\pi}{8} \quad \Rightarrow \quad r > \frac{2D^2}{\lambda}$$

# Extended antennas

$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

$$|\vec{r}-\vec{r}'| = r - \vec{r}' \cdot \hat{i}_r + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$



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$$e^{-j\beta|\vec{r}-\vec{r}'|} = e^{-j\beta r} e^{j\beta r' \cos \gamma} e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \dots \approx e^{-j\beta r} e^{j\beta \vec{r}' \cdot \hat{i}_r} \quad \text{if } r > \frac{2D^2}{\lambda}$$

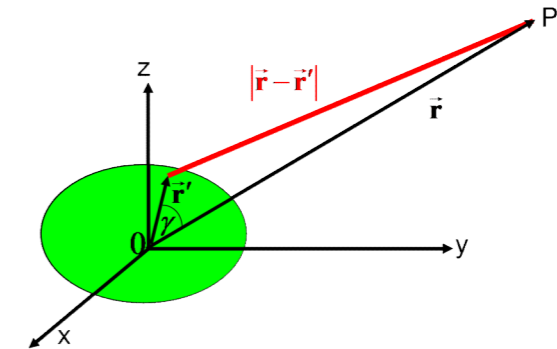
For **all** the antennas, if the distance from the observation point is **sufficiently large**

$$\frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} \approx \frac{e^{-j\beta r} e^{j\beta \vec{r}' \cdot \hat{i}_r}}{r} \quad \longrightarrow \quad \mathbf{A}(\vec{r}) \approx \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \int \mathbf{J}(\vec{r}') e^{j\beta \vec{r}' \cdot \hat{i}_r} d\vec{r}'$$

# Extended antennas

$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

~~$$|\vec{r}-\vec{r}'| = r - \vec{r}' \cdot \hat{i}_r + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$~~



$$\frac{1}{|\vec{r}-\vec{r}'|} = \frac{1}{r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots} \approx \frac{1}{r} \quad \text{if } D \ll r$$

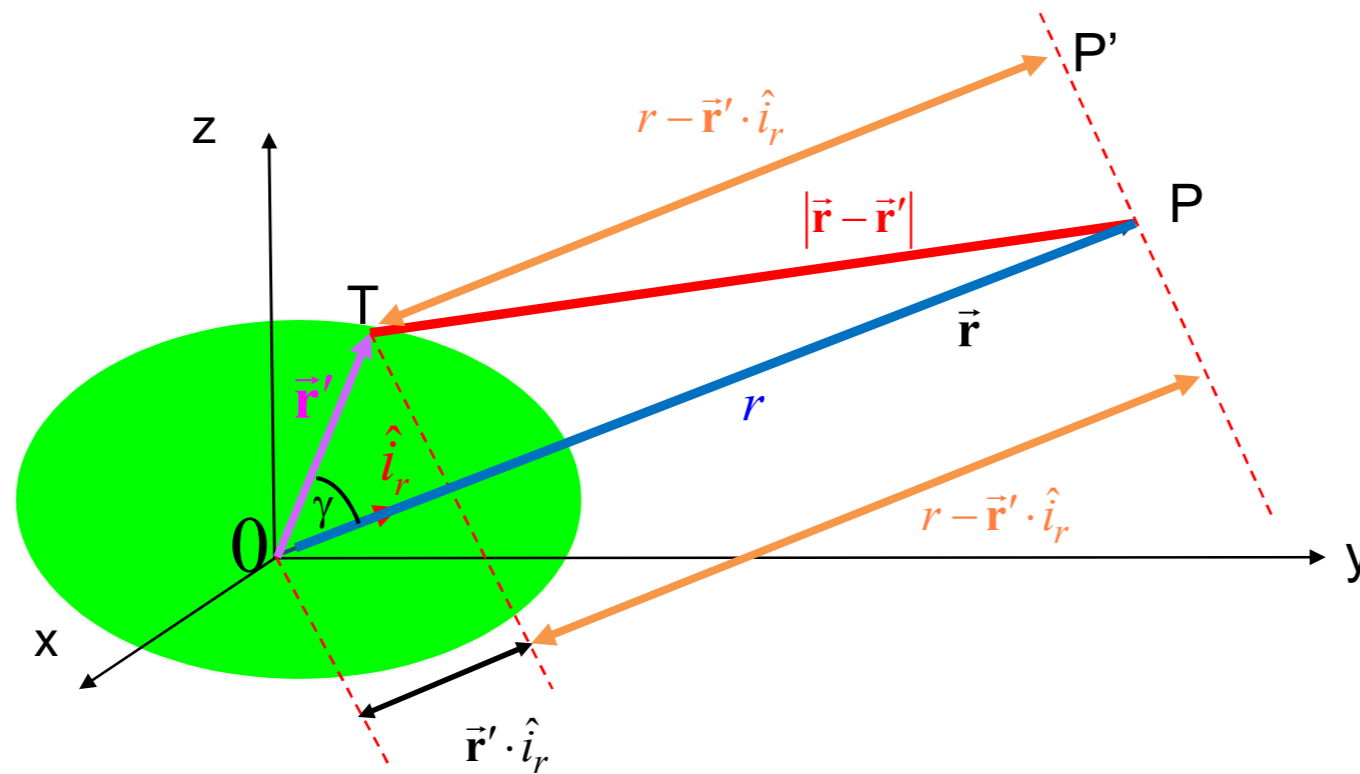
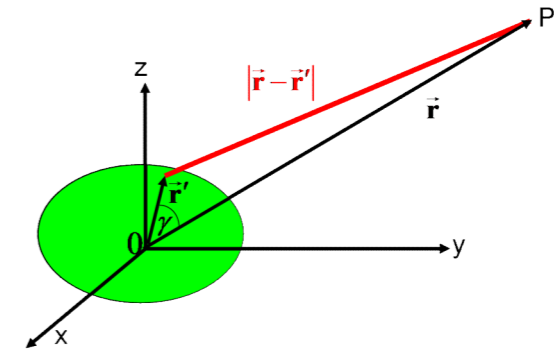
~~$$e^{-j\beta|\vec{r}-\vec{r}'|} = e^{-j\beta r} e^{j\beta r' \cos \gamma} e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \dots$$~~

$$\approx e^{-j\beta r} e^{j\beta \vec{r}' \cdot \hat{i}_r} \quad \text{if } r > \frac{2D^2}{\lambda}$$

# Extended antennas

$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

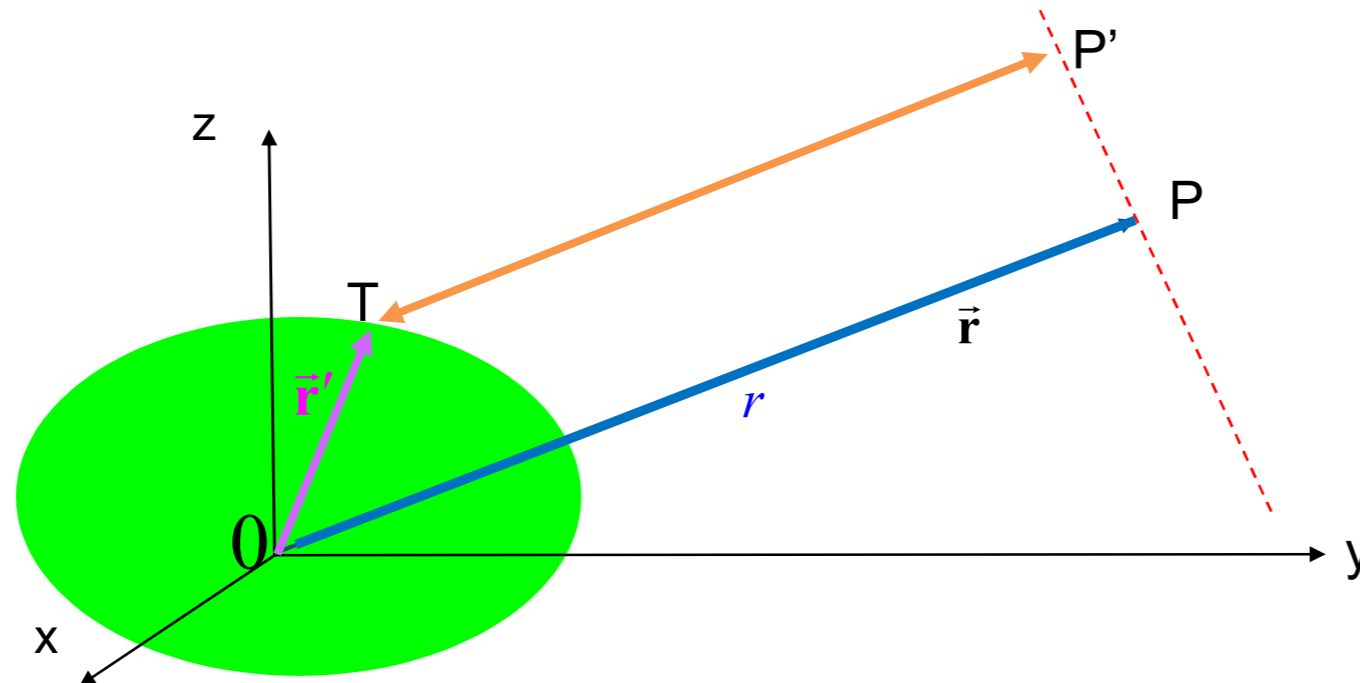
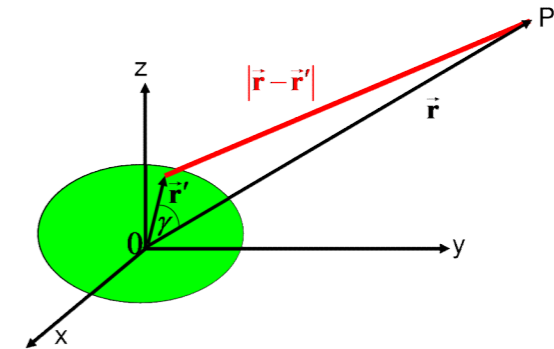
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# Extended antennas

$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

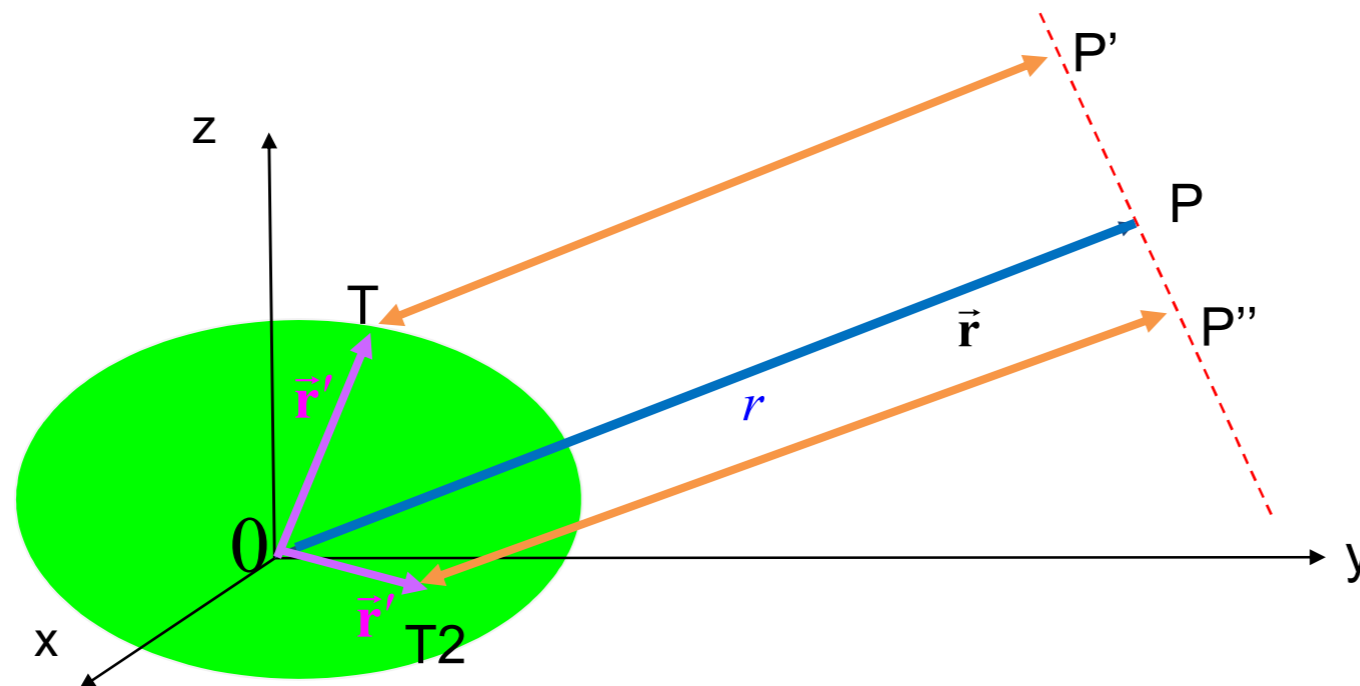
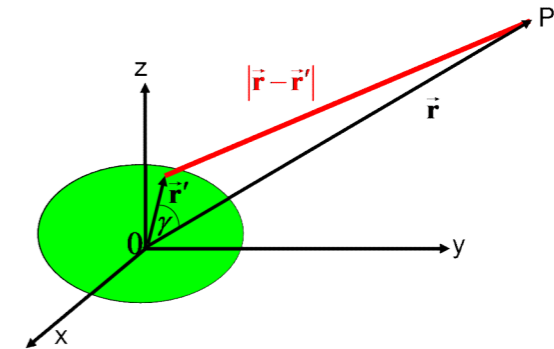
~~$$|\vec{r}-\vec{r}'| = r - \vec{r}' \cdot \hat{i}_r + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$~~



# Extended antennas

$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

~~$$|\vec{r}-\vec{r}'| = r - \vec{r}' \cdot \hat{i}_r + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$~~

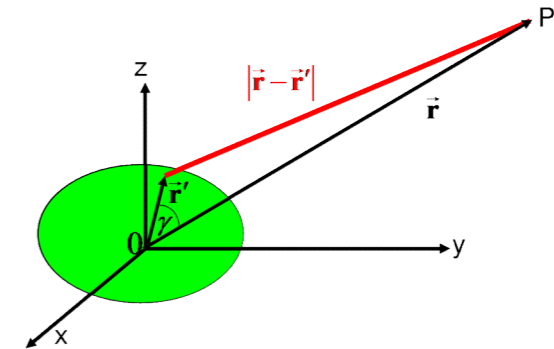




# Extended antennas

$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

$$|\vec{r}-\vec{r}'| = r - \vec{r}' \cdot \hat{i}_r + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$



$$\frac{1}{|\vec{r}-\vec{r}'|} = \frac{1}{r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots} \approx \frac{1}{r} \quad \text{if } D \ll r$$

$$e^{-j\beta|\vec{r}-\vec{r}'|} = e^{-j\beta r} e^{j\beta r' \cos \gamma} e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \dots \approx e^{-j\beta r} e^{j\beta \vec{r}' \cdot \hat{i}_r} \quad \text{if } r > \frac{2D^2}{\lambda}$$

For **all** the antennas, if the distance from the observation point is **sufficiently large**

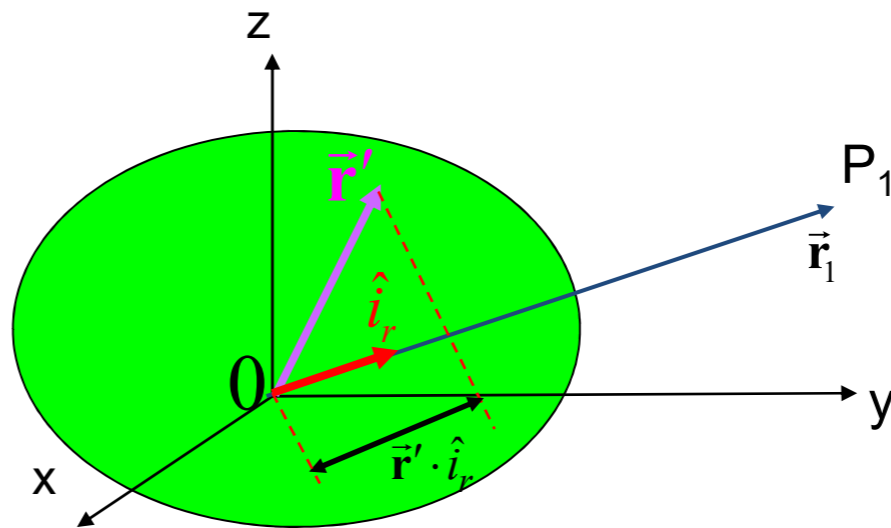
$$\frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} \approx \frac{e^{-j\beta r} e^{j\beta \vec{r}' \cdot \hat{i}_r}}{r} \quad \longrightarrow \quad \mathbf{A}(\vec{r}) \approx \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \int \mathbf{J}(\vec{r}') e^{j\beta \vec{r}' \cdot \hat{i}_r} d\vec{r}'$$

# Extended antennas

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

$$\int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{\mathbf{i}}_r} d\vec{\mathbf{r}}' =$$



For **all** the antennas, if the distance from the observation point is **sufficiently large**

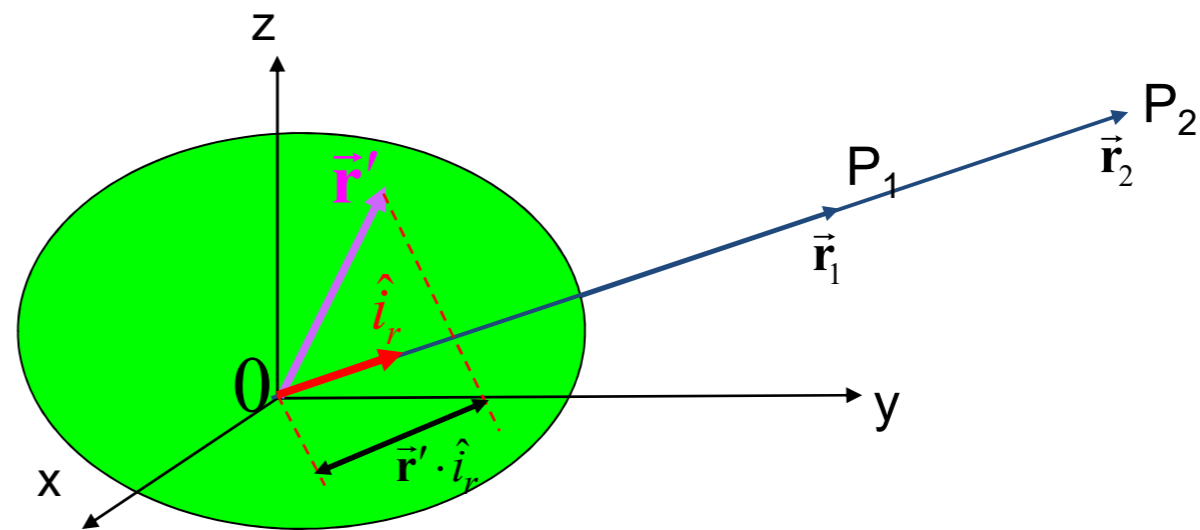
$$\frac{e^{-j\beta|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|}}{|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|} \approx \frac{e^{-j\beta r} e^{j\beta \vec{\mathbf{r}}' \cdot \hat{\mathbf{i}}_r}}{r} \quad \longrightarrow \quad \mathbf{A}(\vec{\mathbf{r}}) \approx \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{\mathbf{i}}_r} d\vec{\mathbf{r}}'$$

# Extended antennas

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

$$\int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{\mathbf{i}}_r} d\vec{\mathbf{r}}' =$$



For **all** the antennas, if the distance from the observation point is sufficiently large

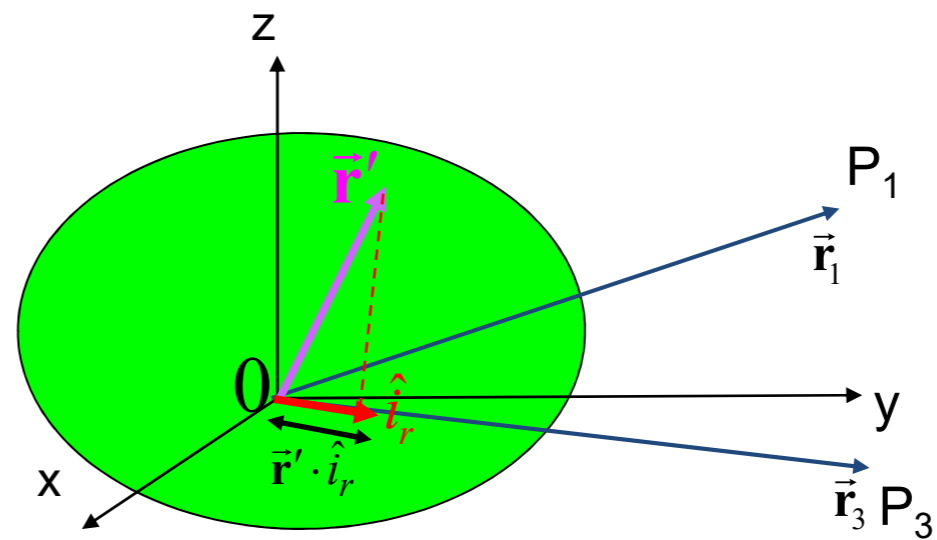
$$\frac{e^{-j\beta|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|}}{|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|} \approx \frac{e^{-j\beta r} e^{j\beta \vec{\mathbf{r}}' \cdot \hat{\mathbf{i}}_r}}{r} \quad \longrightarrow \quad \mathbf{A}(\vec{\mathbf{r}}) \approx \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{\mathbf{i}}_r} d\vec{\mathbf{r}}'$$

# Extended antennas

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

$$\int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{\mathbf{i}}_r} d\vec{\mathbf{r}}' = \mathbf{M}(\vartheta, \varphi) \quad \longrightarrow \quad \mathbf{A}(\vec{\mathbf{r}}) \approx \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \mathbf{M}(\vartheta, \varphi)$$



For **all** the antennas, if the distance from the observation point is sufficiently large

$$\frac{e^{-j\beta|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|}}{|\vec{\mathbf{r}}-\vec{\mathbf{r}}'|} \approx \frac{e^{-j\beta r} e^{j\beta \vec{\mathbf{r}}' \cdot \hat{\mathbf{i}}_r}}{r} \quad \longrightarrow \quad \mathbf{A}(\vec{\mathbf{r}}) \approx \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{\mathbf{i}}_r} d\vec{\mathbf{r}}'$$

# Extended antennas

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

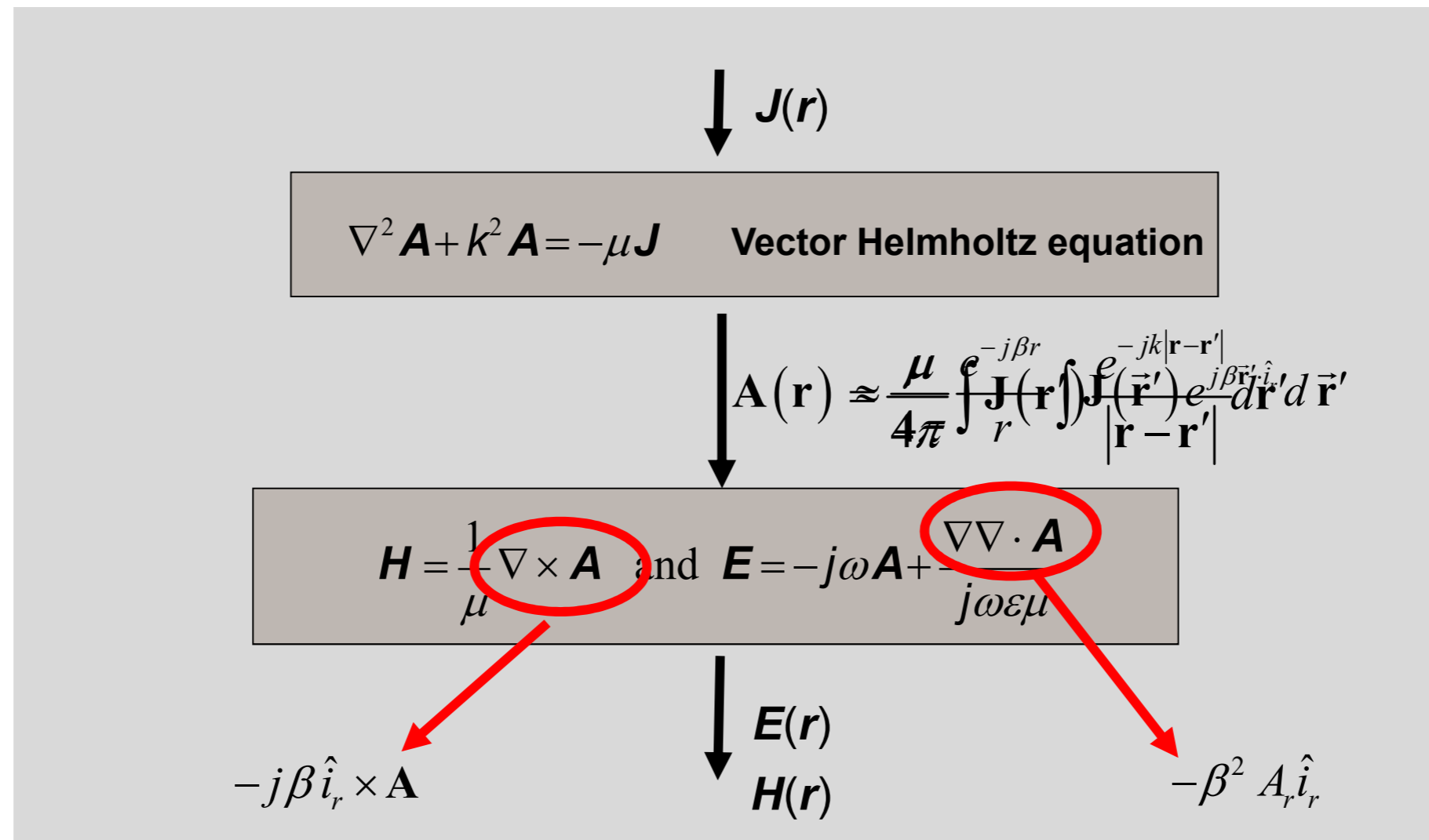
$$\int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{\mathbf{i}}_r} d\vec{\mathbf{r}}' = \mathbf{M}(\vartheta, \varphi) \quad \longrightarrow \quad \mathbf{A}(\vec{\mathbf{r}}) \approx \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \mathbf{M}(\vartheta, \varphi)$$

# Fraunhofer region

$$\begin{array}{l} r \gg D \\ r > \frac{2D^2}{\lambda} \\ r \gg \lambda \end{array} \longrightarrow \int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{\mathbf{i}}_r} d\vec{\mathbf{r}}' = \mathbf{M}(\vartheta, \varphi) \longrightarrow \mathbf{A}(\vec{\mathbf{r}}) \approx \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \mathbf{M}(\vartheta, \varphi)$$
$$\left\{ \begin{array}{l} \nabla \nabla \cdot \mathbf{A}(\vec{\mathbf{r}}) \approx -\beta^2 A_r(\vec{\mathbf{r}}) \hat{\mathbf{i}}_r \\ \nabla \times \mathbf{A}(\vec{\mathbf{r}}) \approx -j\beta \hat{\mathbf{i}}_r \times \mathbf{A}(\vec{\mathbf{r}}) \end{array} \right.$$

Fraunhofer region

# Radiation problem for extended antennas



$$\begin{array}{l}
 r \gg D \\
 r > \frac{2D^2}{\lambda} \\
 r \gg \lambda
 \end{array}
 \longrightarrow
 \int \mathbf{J}(\vec{r}') e^{j\beta \vec{r}' \cdot \hat{i}_r} d\vec{r}' = \mathbf{M}(\vartheta, \varphi) \longrightarrow \mathbf{A}(\vec{r}) \approx \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \mathbf{M}(\vartheta, \varphi)$$

$$\begin{array}{l}
 \nabla \nabla \cdot \mathbf{A}(\vec{r}) \approx -\beta^2 A_r(\vec{r}) \hat{i}_r \\
 \nabla \times \mathbf{A}(\vec{r}) \approx -j\beta \hat{i}_r \times \mathbf{A}(\vec{r})
 \end{array}$$

$$\nabla \nabla \cdot \mathbf{A}(\vec{r}) \approx -\beta^2 A_r(\vec{r}) \hat{i}_r$$



$$\mathbf{E} = -j\omega \mathbf{A} + \frac{\nabla \nabla \cdot \mathbf{A}}{j\omega \epsilon \mu} \approx -j\omega \mathbf{A} - \frac{\beta^2 A_r \hat{i}_r}{j\omega \epsilon \mu} = -j\omega \mathbf{A} + j\omega A_r \hat{i}_r = -j\omega [\mathbf{A} - A_r \hat{i}_r]$$

$$\mathbf{H} = \frac{1}{\mu} \nabla \times \mathbf{A} \quad \text{and} \quad \mathbf{E} = -j\omega \mathbf{A} + \frac{\nabla \nabla \cdot \mathbf{A}}{j\omega \epsilon \mu}$$

$$\nabla \nabla \cdot \mathbf{A}(\vec{r}) \approx -\beta^2 A_r(\vec{r}) \hat{i}_r$$

$$\nabla \times \mathbf{A}(\vec{r}) \approx -j\beta \hat{i}_r \times \mathbf{A}(\vec{r})$$



$$\mathbf{E}(\vec{\mathbf{r}}) = -j\omega \left[ \mathbf{A}(\vec{\mathbf{r}}) - A_r(\vec{\mathbf{r}}) \hat{i}_r \right]$$

$$\begin{aligned} r &\gg D \\ r &> \frac{2D^2}{\lambda} \\ r &\gg \lambda \end{aligned}$$

$$\int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{i}_r} d\vec{\mathbf{r}}' = \mathbf{M}(\vartheta, \varphi) \longrightarrow \mathbf{A}(\vec{\mathbf{r}}) \approx \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \mathbf{M}(\vartheta, \varphi)$$

$$\nabla \nabla \cdot \mathbf{A}(\vec{\mathbf{r}}) \approx -\beta^2 A_r(\vec{\mathbf{r}}) \hat{i}_r$$

$$\nabla \times \mathbf{A}(\vec{\mathbf{r}}) \approx -j\beta \hat{i}_r \times \mathbf{A}(\vec{\mathbf{r}})$$

$$\nabla \nabla \cdot \mathbf{A}(\vec{\mathbf{r}}) \approx -\beta^2 A_r(\vec{\mathbf{r}}) \hat{i}_r$$



$$\mathbf{E} = -j\omega \mathbf{A} + \frac{\nabla \nabla \cdot \mathbf{A}}{j\omega \epsilon \mu} \approx -j\omega \mathbf{A} - \frac{\beta^2 A_r \hat{i}_r}{j\omega \epsilon \mu} = -j\omega \mathbf{A} + j\omega A_r \hat{i}_r = -j\omega \left[ \mathbf{A} - A_r \hat{i}_r \right]$$

$$\mathbf{H} = \frac{1}{\mu} \nabla \times \mathbf{A} \quad \text{and} \quad \mathbf{E} = -j\omega \mathbf{A} + \frac{\nabla \nabla \cdot \mathbf{A}}{j\omega \epsilon \mu}$$

$$\nabla \nabla \cdot \mathbf{A}(\vec{\mathbf{r}}) \approx -\beta^2 A_r(\vec{\mathbf{r}}) \hat{i}_r$$

$$\nabla \times \mathbf{A}(\vec{\mathbf{r}}) \approx -j\beta \hat{i}_r \times \mathbf{A}(\vec{\mathbf{r}})$$

$$\mathbf{E}(\vec{\mathbf{r}}) = -j\omega \left[ \mathbf{A}(\vec{\mathbf{r}}) - A_r(\vec{\mathbf{r}}) \hat{i}_r \right]$$

$$\begin{aligned} r &\gg D \\ r &> \frac{2D^2}{\lambda} \\ r &\gg \lambda \end{aligned}$$

$$\int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{i}_r} d\vec{\mathbf{r}}' = \mathbf{M}(\vartheta, \varphi) \longrightarrow \mathbf{A}(\vec{\mathbf{r}}) \approx \frac{\mu}{4\pi r} e^{-j\beta r} \mathbf{M}(\vartheta, \varphi)$$

$$\nabla \nabla \cdot \mathbf{A}(\vec{\mathbf{r}}) \approx -\beta^2 A_r(\vec{\mathbf{r}}) \hat{i}_r$$

$$\nabla \times \mathbf{A}(\vec{\mathbf{r}}) \approx -j\beta \hat{i}_r \times \mathbf{A}(\vec{\mathbf{r}})$$

$$\nabla \times \mathbf{A}(\vec{\mathbf{r}}) \approx -j\beta \hat{i}_r \times \mathbf{A}(\vec{\mathbf{r}})$$



$$\zeta \mathbf{H} = \frac{\zeta}{\mu} \nabla \times \mathbf{A} \approx \frac{\zeta}{\mu} (-j\beta) \hat{i}_r \times \mathbf{A} = \frac{\zeta \sqrt{\mu} \sqrt{\varepsilon}}{\mu} \hat{i}_r \times [-j\omega \mathbf{A}] = \hat{i}_r \times [-j\omega \mathbf{A}] = \hat{i}_r \times \left[ -j\omega (\mathbf{A} - A_r \hat{i}_r) \right] = \hat{i}_r \times \mathbf{E}$$

$$\mathbf{H} = \frac{1}{\mu} \nabla \times \mathbf{A} \quad \text{and} \quad \mathbf{E} = -j\omega \mathbf{A} + \frac{\nabla \nabla \cdot \mathbf{A}}{j\omega \varepsilon \mu}$$

$$\nabla \nabla \cdot \mathbf{A}(\vec{\mathbf{r}}) \approx -\beta^2 A_r(\vec{\mathbf{r}}) \hat{i}_r$$

$$\nabla \times \mathbf{A}(\vec{\mathbf{r}}) \approx -j\beta \hat{i}_r \times \mathbf{A}(\vec{\mathbf{r}})$$

$$\mathbf{E}(\vec{r}) = -j\omega \left[ \mathbf{A}(\vec{r}) - A_r(\vec{r}) \hat{i}_r \right]$$

$$\zeta \mathbf{H}(\vec{r}) = \hat{i}_r \times \mathbf{E}(\vec{r})$$

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

$$r \gg \lambda$$

$$\int \mathbf{J}(\vec{r}') e^{j\beta \vec{r}' \cdot \hat{i}_r} d\vec{r}' = \mathbf{M}(\vartheta, \varphi) \longrightarrow \mathbf{A}(\vec{r}) \approx \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \mathbf{M}(\vartheta, \varphi)$$

$$\nabla \nabla \cdot \mathbf{A}(\vec{r}) \approx -\beta^2 A_r(\vec{r}) \hat{i}_r$$

$$\nabla \times \mathbf{A}(\vec{r}) \approx -j\beta \hat{i}_r \times \mathbf{A}(\vec{r})$$

$$\nabla \times \mathbf{A}(\vec{r}) \approx -j\beta \hat{i}_r \times \mathbf{A}(\vec{r})$$



$$\zeta \mathbf{H} = \frac{\zeta}{\mu} \nabla \times \mathbf{A} \approx \frac{\zeta}{\mu} (-j\beta) \hat{i}_r \times \mathbf{A} = \frac{\zeta \sqrt{\mu} \sqrt{\varepsilon}}{\mu} \hat{i}_r \times [-j\omega \mathbf{A}] = \hat{i}_r \times [-j\omega \mathbf{A}] = \hat{i}_r \times \left[ -j\omega (\mathbf{A} - A_r \hat{i}_r) \right] = \hat{i}_r \times \mathbf{E}$$

$$\mathbf{H} = \frac{1}{\mu} \nabla \times \mathbf{A} \quad \text{and} \quad \mathbf{E} = -j\omega \mathbf{A} + \frac{\nabla \nabla \cdot \mathbf{A}}{j\omega \varepsilon \mu}$$

$$\nabla \nabla \cdot \mathbf{A}(\vec{r}) \approx -\beta^2 A_r(\vec{r}) \hat{i}_r$$

$$\nabla \times \mathbf{A}(\vec{r}) \approx -j\beta \hat{i}_r \times \mathbf{A}(\vec{r})$$

# Fraunhofer region

$$\begin{array}{l}
 r \gg D \\
 r > \frac{2D^2}{\lambda} \\
 r \gg \lambda
 \end{array}
 \longrightarrow
 \int \mathbf{J}(\vec{r}') e^{j\beta \vec{r}' \cdot \hat{i}_r} d\vec{r}' = \mathbf{M}(\vartheta, \varphi) \longrightarrow \mathbf{A}(\vec{r}) \approx \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \mathbf{M}(\vartheta, \varphi)$$

$$\left\{ \begin{array}{l}
 \nabla \nabla \cdot \mathbf{A}(\vec{r}) \approx -\beta^2 A_r(\vec{r}) \hat{i}_r \\
 \nabla \times \mathbf{A}(\vec{r}) \approx -j\beta \hat{i}_r \times \mathbf{A}(\vec{r})
 \end{array} \right.$$

Fraunhofer region



$$\mathbf{A}(\vec{r}) = \cancel{A_r(\vec{r})} \hat{i}_r + A_\vartheta(\vec{r}) \hat{i}_\vartheta + A_\varphi(\vec{r}) \hat{i}_\varphi$$

$$\mathbf{E}(\vec{r}) = -j\omega \left[ \mathbf{A}(\vec{r}) - A_r(\vec{r}) \hat{i}_r \right] = -j\omega \left[ \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} \mathbf{M}(\vartheta, \varphi) - \frac{\mu}{4\pi} \frac{e^{-j\beta r}}{r} M_r(\vartheta, \varphi) \hat{i}_r \right] = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} \left[ \mathbf{M}(\vartheta, \varphi) - M_r(\vartheta, \varphi) \hat{i}_r \right]$$

$$\zeta \mathbf{H}(\vec{r}) = \hat{i}_r \times \mathbf{E}(\vec{r})$$

# Fraunhofer region

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

$$r \gg \lambda$$

$$\int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{\mathbf{i}}_r} d\vec{\mathbf{r}}' = \mathbf{M}(\vartheta, \varphi)$$

Fraunhofer region

$$\mathbf{E}(\vec{\mathbf{r}}) =$$

$$= -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} [\mathbf{M}(\vartheta, \varphi) - M_r(\vartheta, \varphi) \hat{\mathbf{i}}_r]$$

$$\zeta \mathbf{H}(\vec{\mathbf{r}}) = \hat{\mathbf{i}}_r \times \mathbf{E}(\vec{\mathbf{r}})$$

# Fraunhofer region

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

$$r \gg \lambda$$

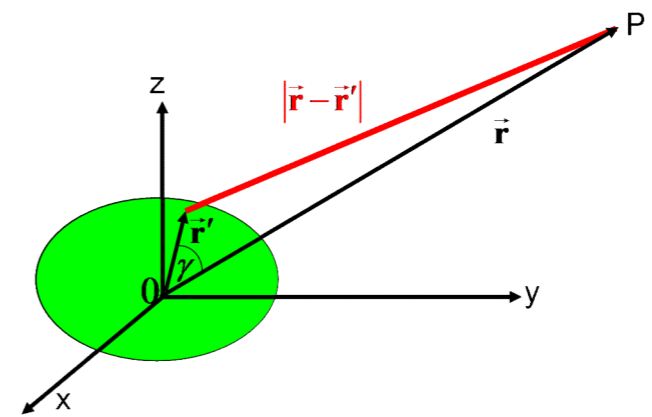
Fraunhofer region

$$\mathbf{E}(\vec{r}) = \mathbf{E}(r, \vartheta, \varphi) = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} \left[ \mathbf{M}(\vartheta, \varphi) - M_r(\vartheta, \varphi) \hat{i}_r \right]$$

$$\zeta \mathbf{H}(\vec{r}) = \hat{i}_r \times \mathbf{E}(\vec{r})$$

- the e.m. field propagates along  $\hat{i}_r$
- the e.m. field lies on the plane orthogonal to the propagation direction
- $|E|$  and  $|H|$  exhibit the decaying factor  $1/r$
- $|E|$  and  $|H|$  are proportional through  $\zeta$

$$\mathbf{M}(\vartheta, \varphi) = \int \mathbf{J}(\vec{r}') e^{j\beta \vec{r}' \cdot \hat{i}_r} d\vec{r}'$$



# Field regions

*Far-field (Fraunhofer) region* is defined as “that region of the field of an antenna where the angular field distribution is essentially independent of the distance from the antenna. If the antenna has a maximum overall dimension  $D$  ( $D > \lambda$ ), the far-field region is commonly taken to exist at distances greater than  $2D^2/\lambda$  from the antenna,  $\lambda$  being the wavelength”.

In this region, the field components are essentially transverse

# Fraunhofer region

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

$$r \gg \lambda$$

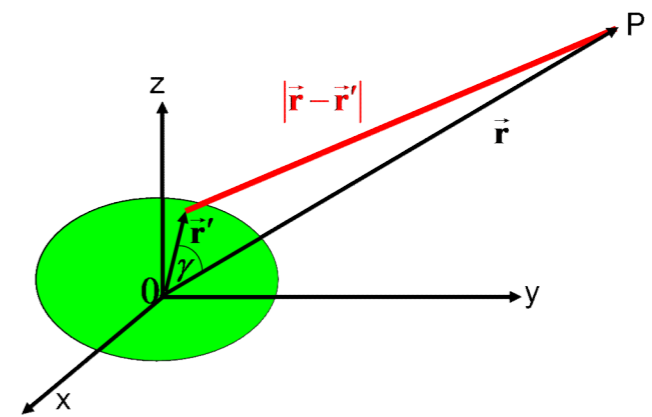
Fraunhofer region

$$\mathbf{E}(\vec{\mathbf{r}}) = \mathbf{E}(r, \vartheta, \varphi) = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} \left[ \mathbf{M}(\vartheta, \varphi) - M_r(\vartheta, \varphi) \hat{i}_r \right]$$

$$\zeta \mathbf{H}(\vec{\mathbf{r}}) = \hat{i}_r \times \mathbf{E}(\vec{\mathbf{r}})$$

- the e.m. field propagates along  $\hat{i}_r$
- the e.m. field lies on the plane orthogonal to the propagation direction
- $|E|$  and  $|H|$  exhibit the decaying factor  $1/r$
- $|E|$  and  $|H|$  are proportional through  $\zeta$

$$\mathbf{M}(\vartheta, \varphi) = \int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{i}_r} d\vec{\mathbf{r}}'$$





# The radiation condition

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

$$r \gg \lambda$$

Fraunhofer region

$$\mathbf{E}(\vec{r}) = \mathbf{E}(r, \vartheta, \varphi) = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} \left[ \mathbf{M}(\vartheta, \varphi) - M_r(\vartheta, \varphi) \hat{i}_r \right]$$

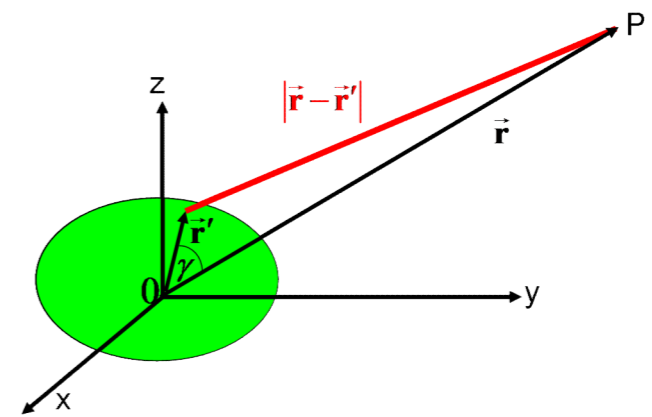
$$\zeta \mathbf{H}(\vec{r}) = \hat{i}_r \times \mathbf{E}(\vec{r})$$

$$\mathbf{E} \sim O\left(\frac{1}{r}\right)$$

$$\mathbf{H} \sim O\left(\frac{1}{r}\right) \quad \text{as } r \rightarrow \infty$$

$$\zeta \mathbf{H} - \hat{i}_r \times \mathbf{E} \sim o\left(\frac{1}{r}\right)$$

$$\mathbf{M}(\vartheta, \varphi) = \int \mathbf{J}(\vec{r}') e^{j\beta \vec{r}' \cdot \hat{i}_r} d\vec{r}'$$



# Fraunhofer region

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

$$r \gg \lambda$$

$$\mathbf{E}(\vec{r}) = \mathbf{E}(r, \vartheta, \varphi) = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} [\mathbf{M}(\vartheta, \varphi) - M_r(\vartheta, \varphi) \hat{i}_r] = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} [M_\vartheta(\vartheta, \varphi) \hat{i}_\vartheta + M_\varphi(\vartheta, \varphi) \hat{i}_\varphi]$$

$$\zeta \mathbf{H}(\vec{r}) = \hat{i}_r \times \mathbf{E}(\vec{r})$$

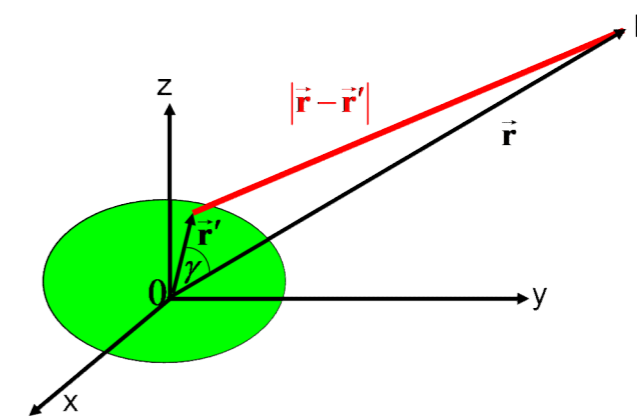
$$\mathbf{M}(\vartheta, \varphi) = \int \mathbf{J}(\vec{r}') e^{j\beta \vec{r}' \cdot \hat{i}_r} d\vec{r}'$$

## Fraunhofer region

$$\vec{\mathbf{S}} = \frac{1}{2} \vec{\mathbf{E}} \times \vec{\mathbf{H}}^* = \frac{1}{2\zeta} \vec{\mathbf{E}} \times (\hat{i}_r \times \vec{\mathbf{E}})^* = \frac{1}{2\zeta} \vec{\mathbf{E}} \times (\hat{i}_r \times \vec{\mathbf{E}}^*) = \frac{1}{2\zeta} \left[ |\vec{\mathbf{E}}|^2 \hat{i}_r - (\hat{i}_r \times \vec{\mathbf{E}}) \vec{\mathbf{E}}^* \right] = \frac{1}{2\zeta} |\vec{\mathbf{E}}|^2 \hat{i}_r$$

$$\vec{\mathbf{S}} = \frac{1}{2} \vec{\mathbf{E}} \times \vec{\mathbf{H}}^* = \frac{\zeta}{2} |\vec{\mathbf{H}}|^2 \hat{i}_r$$

$$\vec{\mathbf{A}} \times (\vec{\mathbf{B}} \times \vec{\mathbf{C}}) = \vec{\mathbf{B}}(\vec{\mathbf{A}} \cdot \vec{\mathbf{C}}) - \vec{\mathbf{C}}(\vec{\mathbf{A}} \cdot \vec{\mathbf{B}})$$



# Fraunhofer region

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

$$r \gg \lambda$$

Fraunhofer region

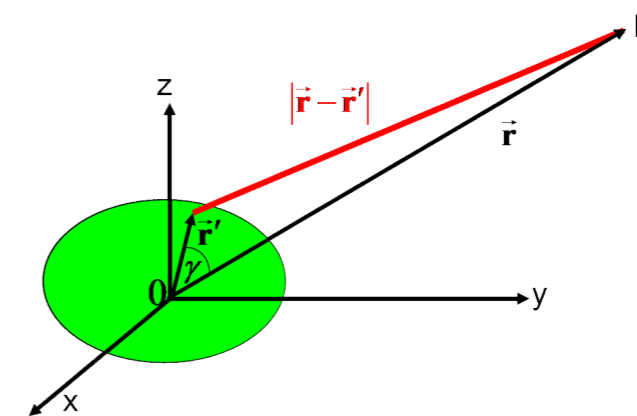
$$\mathbf{E}(\vec{\mathbf{r}}) = \mathbf{E}(r, \vartheta, \varphi) = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} \left[ \mathbf{M}(\vartheta, \varphi) - M_r(\vartheta, \varphi) \hat{i}_r \right] = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} \left[ M_\vartheta(\vartheta, \varphi) \hat{i}_\vartheta + M_\varphi(\vartheta, \varphi) \hat{i}_\varphi \right]$$

$$\zeta \mathbf{H}(\vec{\mathbf{r}}) = \hat{i}_r \times \mathbf{E}(\vec{\mathbf{r}})$$

- the e.m. field propagates along  $\hat{i}_r$
- the e.m. field lies on the plane orthogonal to the propagation direction
- $|E|$  and  $|H|$  exhibit the decaying factor  $1/r$
- $|E|$  and  $|H|$  are proportional through  $\zeta$

$$\vec{\mathbf{S}} = \frac{1}{2\zeta} |\vec{\mathbf{E}}|^2 \hat{i}_r = \frac{\zeta}{2} |\vec{\mathbf{H}}|^2 \hat{i}_r$$

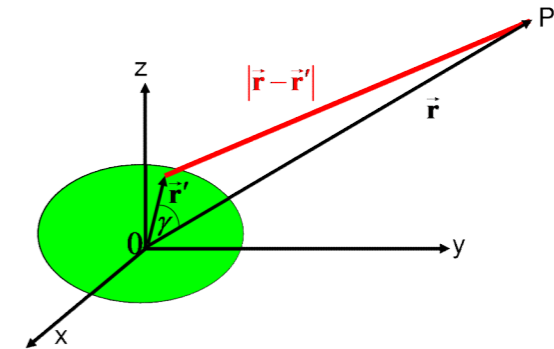
$$\mathbf{M}(\vartheta, \varphi) = \int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{i}_r} d\vec{\mathbf{r}}'$$



# Fraunhofer region

$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

$$|\vec{r}-\vec{r}'| = r - \vec{r}' \cdot \hat{i}_r + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$



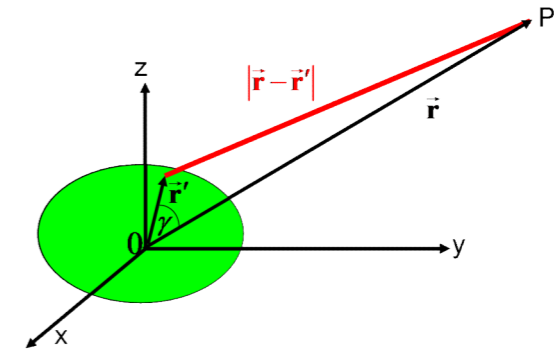
$$\frac{1}{|\vec{r}-\vec{r}'|} = \frac{1}{r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots} \approx \frac{1}{r} \quad \text{if } D \ll r$$

$$e^{-j\beta|\vec{r}-\vec{r}'|} = e^{-j\beta r} e^{j\beta \vec{r}' \cdot \hat{i}_r} e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \dots \approx e^{-j\beta r} e^{j\beta \vec{r}' \cdot \hat{i}_r} \quad \text{if } r > \frac{2D^2}{\lambda}$$

# Fresnel region

$$\mathbf{A}(\vec{r}) = \frac{\mu}{4\pi} \int \mathbf{J}(\vec{r}') \frac{e^{-j\beta|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} d\vec{r}'$$

$$|\vec{r}-\vec{r}'| = r - \vec{r}' \cdot \hat{i}_r + \frac{(r')^2}{2r} \sin^2 \gamma + \dots$$



$$\frac{1}{|\vec{r}-\vec{r}'|} = \frac{1}{r - r' \cos \gamma + \frac{(r')^2}{2r} \sin^2 \gamma + \dots} \approx \frac{1}{r} \quad \text{if } D \ll r$$

$$e^{-j\beta|\vec{r}-\vec{r}'|} \approx e^{-j\beta r} e^{j\beta \vec{r}' \cdot \hat{i}_r} e^{-j\beta \frac{(r')^2}{2r} \sin^2 \gamma} \dots \quad \text{if } r > 0.62 \sqrt{\frac{D^3}{\lambda}}$$

## Fresnel region

# Fraunhofer region

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

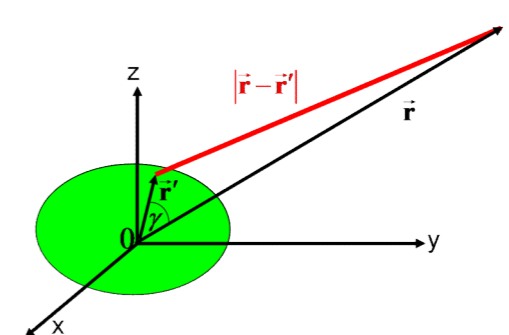
$$r \gg \lambda$$

$$\mathbf{E}(\vec{r}) = \mathbf{E}(r, \vartheta, \varphi) = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} [\mathbf{M}(\vartheta, \varphi) - M_r(\vartheta, \varphi) \hat{i}_r] = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} [M_\vartheta(\vartheta, \varphi) \hat{i}_\vartheta + M_\varphi(\vartheta, \varphi) \hat{i}_\varphi]$$

$$\zeta \mathbf{H}(\vec{r}) = \hat{i}_r \times \mathbf{E}(\vec{r})$$

$$\vec{S} = \frac{1}{2\zeta} |\vec{E}|^2 \hat{i}_r = \frac{\zeta}{2} |\vec{H}|^2 \hat{i}_r$$

$$\mathbf{M}(\vartheta, \varphi) = \int \mathbf{J}(\vec{r}') e^{j\beta \vec{r}' \cdot \hat{i}_r} d\vec{r}'$$



Fraunhofer region

$$\frac{j\omega\mu}{4\pi} = \frac{j\omega\sqrt{\mu}\sqrt{\mu}}{4\pi} \frac{\sqrt{\varepsilon}}{\sqrt{\varepsilon}} = \frac{j\zeta}{4\pi} \beta = \frac{j\zeta}{4\pi} \frac{2\pi}{\lambda} = \frac{j\zeta}{2\lambda}$$

$$\omega\sqrt{\mu\varepsilon} = \beta = \frac{2\pi}{\lambda}$$

$$\zeta = \sqrt{\frac{\mu}{\varepsilon}}$$

# Fraunhofer region

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

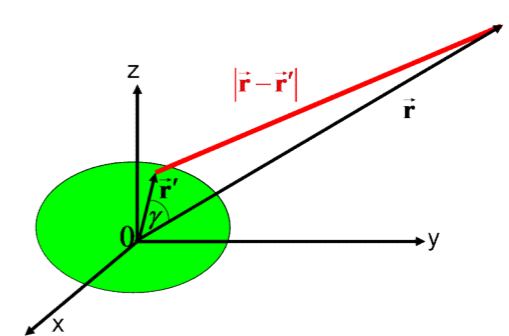
$$r \gg \lambda$$

$$\mathbf{E}(\vec{r}) = \mathbf{E}(r, \vartheta, \varphi) = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} [\mathbf{M}(\vartheta, \varphi) - M_r(\vartheta, \varphi)\hat{i}_r] = \frac{j\zeta}{2\lambda} \frac{e^{-j\beta r}}{r} [M_\vartheta(\vartheta, \varphi)\hat{i}_\vartheta + M_\varphi(\vartheta, \varphi)\hat{i}_\varphi]$$

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Fraunhofer region

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# Fraunhofer region

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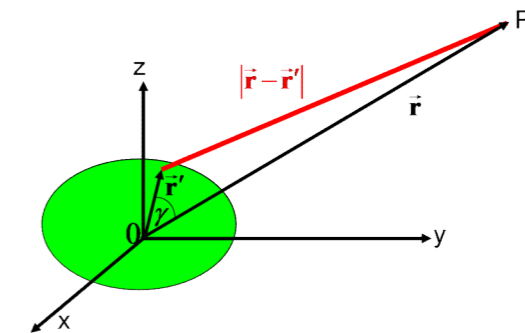
Fraunhofer region

$$\mathbf{E}(\vec{\mathbf{r}}) = \mathbf{E}(r, \vartheta, \varphi) = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} [\mathbf{M}(\vartheta, \varphi) - M_r(\vartheta, \varphi)\hat{i}_r] = -\frac{j\zeta}{2\lambda} \frac{e^{-j\beta r}}{r} [M_\vartheta(\vartheta, \varphi)\hat{i}_\vartheta + M_\varphi(\vartheta, \varphi)\hat{i}_\varphi]$$

$$\zeta \mathbf{H}(\vec{\mathbf{r}}) = \hat{i}_r \times \mathbf{E}(\vec{\mathbf{r}})$$

$$\vec{\mathbf{S}} = \frac{1}{2\zeta} |\vec{\mathbf{E}}|^2 \hat{i}_r = \frac{\zeta}{2} |\vec{\mathbf{H}}|^2 \hat{i}_r$$

$$\mathbf{M}(\vartheta, \varphi) = \int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{i}_r} d\vec{\mathbf{r}}'$$





# Fraunhofer region

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

$$r \gg \lambda$$

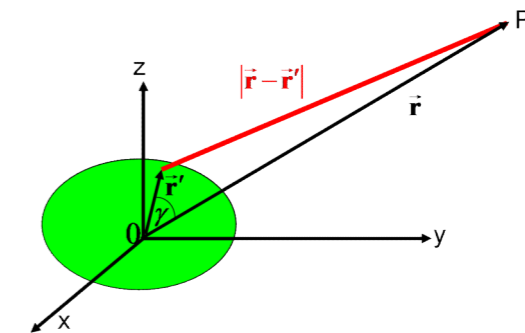
Fraunhofer region

$$\mathbf{E}(\vec{\mathbf{r}}) = \mathbf{E}(r, \vartheta, \varphi) = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} [\mathbf{M}(\vartheta, \varphi) - M_r(\vartheta, \varphi)\hat{i}_r] = \frac{j\zeta}{2\lambda} \frac{e^{-j\beta r}}{r} [-M_\vartheta(\vartheta, \varphi)\hat{i}_\vartheta - M_\varphi(\vartheta, \varphi)\hat{i}_\varphi]$$

$$\zeta \mathbf{H}(\vec{\mathbf{r}}) = \hat{i}_r \times \mathbf{E}(\vec{\mathbf{r}})$$

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# Fraunhofer region

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Fraunhofer region

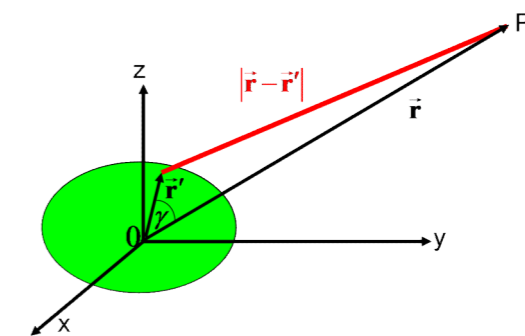
$$\mathbf{E}(\vec{\mathbf{r}}) = \mathbf{E}(r, \vartheta, \varphi) = -\frac{j\omega\mu}{4\pi} \frac{e^{-j\beta r}}{r} \left[ \mathbf{M}(\vartheta, \varphi) - M_r(\vartheta, \varphi) \hat{i}_r \right] = \frac{j\zeta}{2\lambda} \frac{e^{-j\beta r}}{r} I \left[ \frac{-M_\vartheta(\vartheta, \varphi) \hat{i}_\vartheta - M_\varphi(\vartheta, \varphi) \hat{i}_\varphi}{I} \right]$$

$$\zeta \mathbf{H}(\vec{\mathbf{r}}) = \hat{i}_r \times \mathbf{E}(\vec{\mathbf{r}})$$

$I$  is the phasor associated to the input current at the antenna input terminals

$$\vec{\mathbf{S}} = \frac{1}{2\zeta} |\vec{\mathbf{E}}|^2 \hat{i}_r = \frac{\zeta}{2} |\vec{\mathbf{H}}|^2 \hat{i}_r$$

$$\mathbf{M}(\vartheta, \varphi) = \int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{i}_r} d\vec{\mathbf{r}}'$$



# Fraunhofer region

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

$$r \gg \lambda$$

Fraunhofer region

$$\mathbf{E}(\vec{\mathbf{r}}) = \mathbf{E}(r, \vartheta, \varphi) =$$

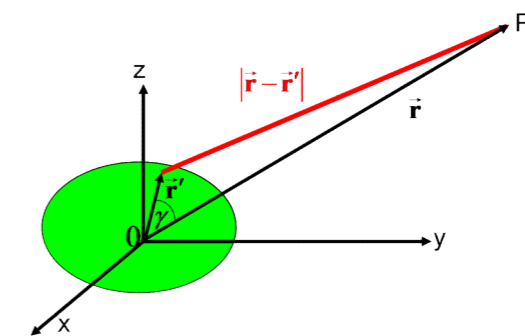
$$\zeta \mathbf{H}(\vec{\mathbf{r}}) = \hat{i}_r \times \mathbf{E}(\vec{\mathbf{r}})$$

$$= \frac{j\zeta}{2\lambda} \frac{e^{-j\beta r}}{r} I \left[ \frac{-M_\vartheta(\vartheta, \varphi) \hat{i}_\vartheta - M_\varphi(\vartheta, \varphi) \hat{i}_\varphi}{I} \right]$$

$I$  is the phasor associated to the input current at the antenna input terminals

$$\vec{\mathbf{S}} = \frac{1}{2\zeta} |\vec{\mathbf{E}}|^2 \hat{i}_r = \frac{\zeta}{2} |\vec{\mathbf{H}}|^2 \hat{i}_r$$

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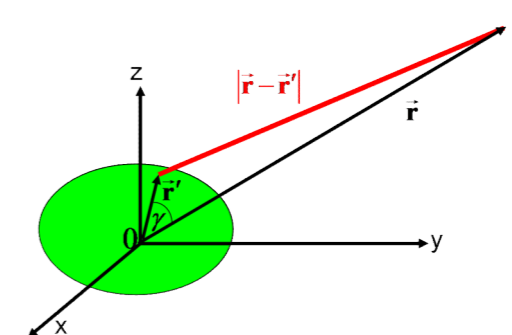
$$\mathbf{E}(\vec{\mathbf{r}}) = \mathbf{E}(r, \vartheta, \varphi) = \frac{j\zeta I e^{-j\beta r}}{2\lambda r} \frac{-M_\vartheta(\vartheta, \varphi)\hat{i}_\vartheta - M_\varphi(\vartheta, \varphi)\hat{i}_\varphi}{I} = \frac{j\zeta I e^{-j\beta r}}{2\lambda r} \mathbf{l}(\vartheta, \varphi)$$

$$\zeta \mathbf{H}(\vec{\mathbf{r}}) = \hat{i}_r \times \mathbf{E}(\vec{\mathbf{r}})$$

$I$  is the phasor associated to the input current at the antenna input terminals

$$\bar{\mathbf{S}} = \frac{1}{2\zeta} |\bar{\mathbf{E}}|^2 \hat{i}_r = \frac{\zeta}{2} |\bar{\mathbf{H}}|^2 \hat{i}_r$$

$$\mathbf{M}(\vartheta, \varphi) = \int \mathbf{J}(\vec{\mathbf{r}}') e^{j\beta \vec{\mathbf{r}}' \cdot \hat{i}_r} d\vec{\mathbf{r}}'$$



Fraunhofer region

$$[\mathbf{M}] = \frac{\text{Ampere}}{m^2} \times m^3 = \text{Ampere} \times m$$

$$\frac{-M_\vartheta(\vartheta, \varphi)\hat{i}_\vartheta - M_\varphi(\vartheta, \varphi)\hat{i}_\varphi}{I} = \mathbf{l}(\vartheta, \varphi) = l_\vartheta(\vartheta, \varphi)\hat{i}_\vartheta + l_\varphi(\vartheta, \varphi)\hat{i}_\varphi$$

is said **effective length** of the antenna

# Fraunhofer region

$$r \gg D$$

$$r > \frac{2D^2}{\lambda}$$

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$$\mathbf{E}(\vec{\mathbf{r}}) = \mathbf{E}(r, \vartheta, \varphi) = \frac{j\zeta I e^{-j\beta r}}{2\lambda r} \mathbf{l}(\vartheta, \varphi)$$

$$\zeta \mathbf{H}(\vec{\mathbf{r}}) = \hat{i}_r \times \mathbf{E}(\vec{\mathbf{r}})$$

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$$\vec{\mathbf{S}} = \frac{1}{2\zeta} |\vec{\mathbf{E}}|^2 \hat{i}_r = \frac{\zeta}{2} |\vec{\mathbf{H}}|^2 \hat{i}_r$$